

High-energy Neutrino follow-up search of Gravitational Wave Event GW150914 with ANTARES and IceCube

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We present the high-energy-neutrino follow-up observations of the first gravitational wave transient GW150914 observed by the Advanced LIGO detectors on Sept. 14th, 2015. We search for coincident neutrino candidates within the data recorded by the IceCube and ANTARES neutrino detectors. A possible joint detection could be used in targeted electromagnetic follow-up observations, given the significantly better angular resolution of neutrino events compared to gravitational waves.

We find no neutrino candidates in both temporal and spatial coincidence with the gravitational wave event. Within ± 500 s of the gravitational wave event, the number of neutrino candidates detected by IceCube and ANTARES were three and zero, respectively. This is consistent with the expected atmospheric background, and none of the neutrino candidates were directionally coincident with GW150914. We use this non-detection to constrain neutrino emission from the gravitational-wave event.

I. INTRODUCTION

Advanced LIGO’s first observation periods [1, 2] represent a major step in probing the dynamical origin of high-energy emission from cosmic transients [3]. The significant improvement in gravitational wave (GW) search sensitivity enables a comprehensive multimessenger observational effort involving partner electromagnetic observatories from radio to gamma-rays, as well as neutrino detectors. The goals of multimessenger observations are to gain a more complete understanding of cosmic processes through a combination of information from different probes, and to increase search sensitivity over an analysis using a single messenger [4, 5].

The merger of neutron stars and black holes, and potentially massive stellar core collapse with rapidly rotating cores, are expected to be significant sources of GWs [3]. These events can result in a black hole plus accretion disk system that drives a relativistic outflow [6, 7]. Energy dissipation in the outflow produces non-thermal, high-energy radiation that is observed as gamma-ray bursts (GRBs), and may have a $\gg$ GeV neutrino component at comparable luminosities.

Multiple detectors have been built that can search for this high-energy neutrino signature, including the IceCube Neutrino Observatory—a cubic-kilometer facility at the South Pole [8, 9, 10], and ANTARES [11, 12, 13] in the Mediterranean sea. The construction of the KM3NeT cubic-kilometer scale neutrino detector in the Mediterranean Sea has started in December 2015 with the successful deployment of the first detection string [14]. IceCube is planning a substantial increase in sensitivity with near-future upgrades [15, 16]. Another facility, the Baikal Neutrino Telescope is also planning an upgrade to cubic-kilometer volume [17]. An astrophysical high-energy neutrino flux has recently been discovered by IceCube [18, 19, 20, 21], demonstrating the production of non-thermal high-energy neutrinos. The specific origin of this neutrino flux is currently unknown. Multimessenger analyses constraining the common sources of high-energy neutrinos and GWs have been carried out in the past with both ANTARES and IceCube [22, 23, 24].

On Sept. 14th, 2015 at 09:50:45 UTC, a highly significant GW signal was recorded by the LIGO Hanford, WA and Livingston, LA detectors [25]. The event, labeled GW150914, was produced by a stellar-mass binary black

hole merger at redshift $z = 0.09^{+0.03}_{-0.04}$. The reconstructed mass of each black hole is $\sim 30 M_{\odot}$. Such a system may produce electromagnetic emission and emit neutrinos if the merger happens in a sufficiently baryon-dense environment, and a black hole plus accretion disk system is formed [26]. Current consensus is that such a scenario is unlikely, nevertheless, there are no significant observational constraints.

Here we report the results of a neutrino follow-up search of GW150914 using ANTARES and IceCube. After brief descriptions of the GW search (Section II) and the neutrino follow-up (Section III), we present the joint analysis, results of the search and source constraints, and conclusions (Section IV).

II. GRAVITATIONAL WAVE DATA ANALYSIS AND DISCOVERY

GW150914 was initially identified by low-latency searches for generic GW transients [27, 28, 29]. Subsequent analysis with three independent matched-filter analyses using models of compact binary coalescence waveforms [30, 31] confirmed that the event was produced by the merger of two black holes. The analyses established a false alarm rate of less than 1 per event per 203 000 years, equivalent to a significance $> 5.1\sigma$ [25]. Source parameters were reconstructed using the LALInference package [31, 32, 33], finding black-hole masses $36^{+5}_{-4} M_{\odot}$ and $29^{+4}_{-4} M_{\odot}$ and luminosity distance $D_{\text{gw}} = 410^{+160}_{-180}$ Mpc. The duration of the signal within LIGO’s sensitive band was 0.2 s.

The directional point spread function (sky map) of the GW event was computed through the full parameter estimation of the signal, carried out using the LALInference package [32, 33]. The LALInference results presented here account for calibration uncertainty in the GW strain signal. The sky map is shown in Fig. 1. At 90% (50%) credible level (CL), the sky map covers 590 deg^2 (140 deg^2).

III. HIGH-ENERGY NEUTRINO COINCIDENCE SEARCH

High-energy neutrino observatories are primarily sensitive to neutrinos with $\gg$ GeV energies. IceCube and ANTARES are both sensitive to through-going muons (called track events), produced by neutrinos near the detector, above ~ 100 GeV. In this analysis, ANTARES data include only up-going tracks for events originat-

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ing from the Southern hemisphere, while IceCube data include both up-going tracks (from the Northern hemisphere) as well as down-going tracks (from the Southern hemisphere). The energy threshold of neutrino candidates increases in the Southern hemisphere for IceCube, since downward-going atmospheric muons are not filtered by the Earth, greatly increasing the background at lower energies. Neutrino times of arrival are determined at μs precision.

Since neutrino telescopes continuously take data observing the whole sky, it is possible to look back and search for neutrino counterparts to an interesting GW signal at any time around the GW observation.

To search for neutrinos coincident with GW150914, we used a time window of $\pm 500\text{ s}$ around the GW transient. This search window, which was used in previous GW-neutrino searches, is a conservative, observation-based upper limit on the plausible emission of GWs and high-energy neutrinos in the case of GRBs, which are thought to be driven by a stellar-mass black hole—accretion disk system [34]. While the relative time of arrival of GWs and neutrinos can be informative [35, 36, 37], here we do not use detailed temporal information beyond the $\pm 500\text{ s}$ time window.

The search for high-energy neutrino candidates recorded by IceCube within $\pm 500\text{ s}$ of GW150914 used IceCube’s online event stream. The online event stream implements an event selection similar to the event selection used for neutrino point source searches [38], but optimized for real-time performance at the South Pole. This event selection consists primarily of cosmic-ray-induced background events, with an expectation per 1000 seconds of 2.2 events in the Northern sky (atmospheric neutrinos), and 2.2 events in the Southern sky (high-energy atmospheric muons). In the search window of $\pm 500\text{ s}$ centered on the GW alert time (see below), one event was found in the Southern sky and two in the Northern sky, which is consistent with the background expectation. The properties of these events are listed in Table I. The neutrino candidates’ directions are shown in Fig. 1.

The muon energy in Table I is reconstructed assuming a single muon is producing the event. While the event from the Southern hemisphere has a significantly greater reconstructed energy [39] than the other two events, 12.5% of the background events in the same declination range in the Southern hemisphere have energies in excess of the one observed. The intense flux of atmospheric muons and bundles of muons that constitute the background for IceCube in the Southern hemisphere gradually falls as the cosmic ray flux declines with energy [40]. The use of energy cuts to remove most of this background is the reason that IceCube’s sensitivity in the Southern sky is shifted to higher energies.

An additional search was performed using the high-energy starting event selection described in [18]. No events were found in coincidence with GW150914.

The IceCube detector also has sensitivity to outbursts of MeV neutrinos (as occur for example in core-collapse

#	ΔT [s]	RA [h]	Dec [°]	$\sigma_{\mu}^{\text{rec}}$ [°]	E_{μ}^{rec} [TeV]	fraction
1	+37.2	8.84	−16.6	0.35	175	12.5%
2	+163.2	11.13	12.0	1.95	1.22	26.5%
3	+311.4	−7.23	8.4	0.47	0.33	98.4%

TABLE I. Parameters of neutrino candidates identified by IceCube within the $\pm 500\text{ s}$ time window around GW150914. ΔT is the time of arrival of the neutrino candidates relative to that of GW150914. E_{μ}^{rec} is the reconstructed muon energy; energies at the source would be $\sim 10\%$ higher. $\sigma_{\mu}^{\text{rec}}$ is the angular uncertainty of the reconstructed track direction [42]. The last column shows the fraction of background neutrino candidates with higher reconstructed energy at the same declination ($\pm 5^\circ$).

supernovae) via a sudden increase in the photomultiplier rates [41]. The global photomultiplier noise rate is monitored continuously, and deviations sufficient to trigger the lowest-level of alert occur roughly once per hour. No alert was triggered during the ± 500 second time-window around the GW candidate event.

The search for coincident neutrinos for ANTARES within $\pm 500\text{ s}$ of GW150914 used ANTARES’s online reconstruction pipeline [43]. A fast and robust algorithm [44] selected up-going neutrino candidates with $\sim \text{mHz}$ rate, with atmospheric muon contamination less than 10%. In addition, to reduce the background of atmospheric neutrinos [45], a requirement of a minimum reconstructed energy reduced the online event rate to 1.2 events/day. Consequently, for ANTARES the expected number of neutrino candidates from the Southern sky in a 1000s window in the Southern sky is 0.014. We found no neutrino event from ANTARES that were temporally coincident with GW150914. This is consistent with the expected background event rate.

IV. RESULTS

A. Joint analysis

We carried out the joint GW and neutrino search following the analysis developed for previous GW and neutrino datasets using initial GW detectors [22, 24, 34, 46]. After identifying the GW event GW150914 with the cWB pipeline, we used reconstructed neutrino candidates to search for temporal and directional coincidences between GW150914 and neutrinos. We assumed that the *a priori* source directional distribution is uniform. For temporal coincidence, we searched within a $\pm 500\text{ s}$ time window around GW150914.

The relative difference in propagation time for $\gg \text{GeV}$ neutrinos and GWs (which travel at the speed of light in general relativity) traveling to Earth from the source is expected to be $\ll 1\text{ s}$. The relative propagation time between neutrinos and GWs may change in alternative gravity models [47, 48]. However, discrepancies from gen-

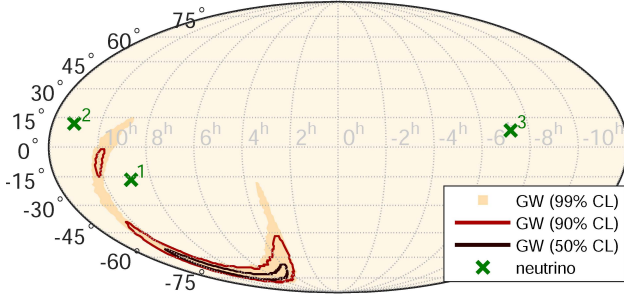


FIG. 1. GW skymap in equatorial coordinates, showing the reconstructed probability density contours of the GW event at 50%, 90% and 99% CL, and the reconstructed directions of high-energy neutrino candidates detected by IceCube (crosses) during a ± 500 s time window around the GW event. The neutrino directional uncertainties are $< 1^\circ$ and are not shown. GW shading indicates the reconstructed probability density of the GW event, darker regions corresponding to higher probability. Neutrino numbers refer to the first column of Table I.

eral relativity could in principle be probed with a joint GW-neutrino detection by comparing the arrival times against the expected time frame of emission.

Directionally, we searched for overlap between the GW sky map and the neutrino point spread functions, assumed to be Gaussian with standard deviation σ_μ^{rec} (see Table I).

The search identified no ANTARES neutrino candidate that were temporally coincident with GW150914.

For IceCube, none of the three neutrino candidates temporally coincident with GW150914 were compatible with the GW direction at 90% CL. Additionally, the reconstructed energy of the neutrino candidates with respect to the expected background does not make them significant. See Fig. 1 for the directional relation of GW150914 and the IceCube neutrino candidates detected within the ± 500 s window. This non-detection is consistent with our expectation from a binary black hole merger.

To better understand the probability that the detected neutrino candidates are being consistent with background, we briefly consider different aspects of the data separately. First, the number of detected neutrino candidates, i.e. 3 and 0 for IceCube and ANTARES, respectively, is fully consistent with the expected background rate of 4.4 and $\ll 1$ for the two detectors, with p-value $1 - F_{\text{pois}}(N_{\text{observed}} \leq 2, N_{\text{expected}} = 4.4) = 0.81$, where F_{pois} is the Poisson cumulative distribution function. Second, for the most significant reconstructed muon energy (Table I), 12.5% of background events will have greater muon energy. The probability that at least one neutrino candidate, out of 3 detected events, has an energy high enough to make it appear even less background-like, is $1 - (1 - 0.125)^3 \approx 0.33$. Third, with the GW sky area 90% CL of $\Omega_{\text{gw}} = 590 \text{ deg}^2$, the probability of a

background neutrino candidate being directionally coincident is $\Omega_{\text{gw}}/\Omega_{\text{all}} \approx 0.014$. We expect $3\Omega_{\text{gw}}/\Omega_{\text{all}}$ directionally coincident neutrinos, given 3 temporal coincidences. Therefore, the probability that at least one of the 3 neutrino candidates is directionally coincident with the 90% CL skymap of GW150914 is $1 - (1 - 0.014)^3 \approx 0.04$.

B. Constraints on the source

We used the non-detection of coincident neutrino candidates by ANTARES and IceCube to derive a standard frequentist neutrino spectral fluence upper limit for GW150914 at 90% CL. Considering no spatially and temporally coincident neutrino candidates, we calculated the source fluence that on average would produce 2.3 detected neutrino candidates. We carried out this analysis as a function of source direction, and independently for ANTARES and IceCube.

The obtained spectral fluence upper limits as a function of source direction are shown in Fig. 2. We consider a standard $dN/dE \propto E^{-2}$ source model, as well as a model with a spectral cutoff at high energies: $dN/dE \propto E^{-2} \exp[-\sqrt{(E/100\text{TeV})}]$. For each spectral model, the upper limit shown in each direction of the sky is the more stringent limit provided by one or the other detector. We see in Fig. 2 that the constraint strongly depends on the source direction, and is mostly within $E^2 dN/dE \sim 10^{-1} - 10 \text{ GeV}^{-1} \text{ cm}^{-2}$. Furthermore, the upper limits by ANTARES and IceCube constrain different energy ranges in the region of the sky close to the GW candidate. For an E^{-2} power-law source spectrum, 90% of ANTARES signal neutrinos are in the energy range from 3 TeV to 1 PeV, whereas for IceCube at this southern declination the corresponding energy range is 200 TeV to 100 PeV.

We now convert our fluence upper limit into a constraint on the total energy emitted in neutrinos by the source. To obtain this constraint, we integrate emission within $[100 \text{ GeV}, 100 \text{ PeV}]$ for the standard $dN/dE \propto E^{-2}$ source model, and within $[100 \text{ GeV}, 100 \text{ TeV}]$ assuming neutrino emission with a cutoff at 100 TeV. We find non-detection to correspond to the following upper limit on the total energy radiated in neutrinos:

$$E_{\nu, \text{tot}}^{\text{ul}} \sim 10^{52} - 10^{54} \left(\frac{D_{\text{gw}}}{410 \text{ Mpc}} \right)^2 \text{ erg} \quad (1)$$

The wide allowed range is primarily due to the large directional uncertainty of the GW event. For comparison, the total energy radiated in GWs from the source is $\sim 5 \times 10^{54} \text{ erg}$. Typical GRB isotropic-equivalent energies are $\sim 10^{51} \text{ erg}$ for long and $\sim 10^{49} \text{ erg}$ for short GRBs [49], which may be similar to the total energy radiated in neutrinos in GRBs [50, 51, 52, 53].

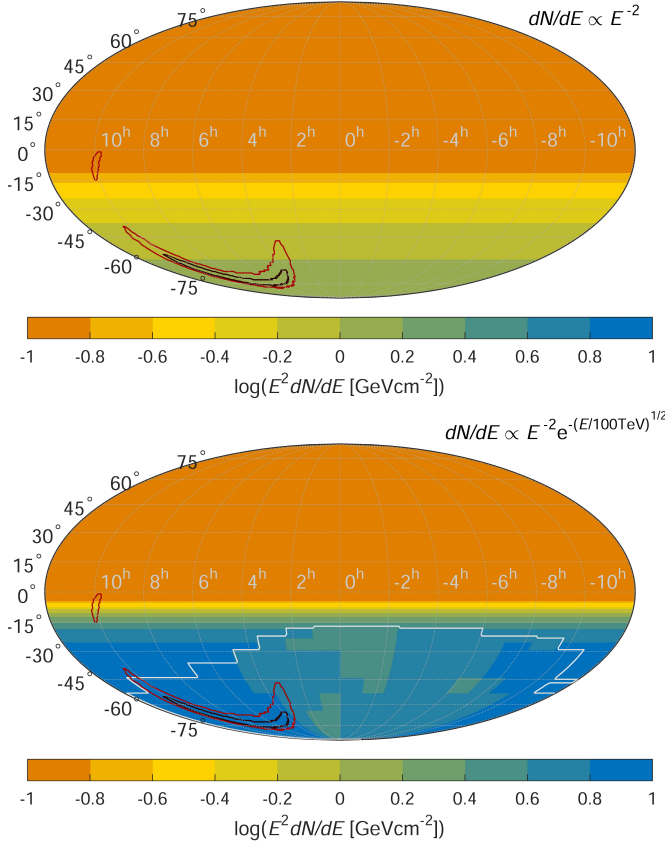


FIG. 2. Upper limit on the high-energy neutrino spectral fluence ($\nu_\mu + \bar{\nu}_\mu$) from GW150914 as a function of source direction, assuming $dN/dE \propto E^{-2}$ (top) and $dN/dE \propto E^{-2} \exp[-\sqrt{E/100\text{TeV}}]$ (bottom) neutrino spectra. The region surrounded by a white line shows the part of the sky in which ANTARES is more sensitive (close to nadir), while on the rest of the sky, IceCube is more sensitive. For comparison, the 50% CL and 90% CL contours of the GW sky map are also shown.

V. CONCLUSION

The results above represent the first concrete limit on neutrino emission from this GW source type, and the first neutrino follow-up of a significant GW event. With the continued increase of Advanced LIGO-Virgo sensitivities for the next observation periods, and the implied source rate of $2\text{--}400 \text{ Gpc}^{-3}\text{yr}^{-1}$ in the comoving frame based on this first detection [54], we can expect to detect a significant number of GW sources, allowing for stacked neutrino analyses and significantly improved constraints. Similar analyses for the upcoming observation periods of Advanced LIGO-Virgo will be important to provide constraints on or to detect other joint GW and neutrino sources.

Joint GW and neutrino searches will also be used to improve the efficiency of electromagnetic follow-up observations over GW-only triggers. Given the significantly

more accurate direction reconstruction of neutrinos ($\sim 1 \text{ deg}^2$ for track events in IceCube [39, 42] and $\sim 0.2 \text{ deg}^2$ in ANTARES [55]) compared to GWs ($\gtrsim 100 \text{ deg}^2$), a joint event candidate provides a greatly reduced sky area for follow-up observatories [56]. The delay induced by the event filtering and reconstruction after the recorded trigger time is typically 3–5 s for ANTARES [43], 20–30 s for IceCube [57], and $\mathcal{O}(1 \text{ min})$ for LIGO-Virgo, making data available for rapid analyses.

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