

Experimental investigation of a control scheme for a zero-detuning resonant sideband extraction interferometer for next-generation gravitational-wave detectors

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Abstract

Some next-generation gravitational-wave detectors, such as the American Advanced LIGO project and the Japanese LCGT project, plan to use power recycled resonant sideband extraction (RSE) interferometers for their interferometer's optical configuration. A power recycled zero-detuning (PRZD) RSE interferometer, which is the default design for LCGT, has five main length degrees of freedom that need to be controlled in order to operate a gravitational-wave detector. This task is expected to be very challenging because of the complexity of optical configuration. A new control scheme for a PRZD RSE interferometer has been developed and tested with a prototype interferometer. The PRZD RSE interferometer was successfully locked with the control scheme. It is the first experimental demonstration of a PRZD RSE interferometer with suspended test masses. The result serves as an important step for the operation of LCGT.

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(Some figures in this article are in colour only in the electronic version)

1. Introduction

Presently several laser interferometer gravitational-wave detectors are in operation in the United States (LIGO [1]), in Europe (GEO600 [2] and VIRGO [3]), and in Japan (TAMA300 [4]). In addition to the present detectors, there are plans to upgrade them to next-generation interferometers. Amongst them are Advanced LIGO [5] and LCGT [6], which plan to use the power recycled resonant sideband extraction (RSE) technique to enhance detector sensitivities.

Despite the great advantage of being able to achieve better sensitivity by avoiding problems of thermal absorption by substrates, the RSE configuration poses a more difficult challenge in controlling the interferometer in order to use it as a gravitational-wave detector due to the increased number of degrees of freedom (DOF) that need to be controlled. Therefore, designing and demonstrating a control scheme as simple as possible are vital before the technique is adapted in large-scale interferometers such as LCGT. Recently, a control of a power recycled detuned RSE has been demonstrated with a prototype experiment on the 40 m interferometer at Caltech [7], and a zero-detuned signal recycling as well as a detuned signal recycling has been demonstrated on the GEO600 detector [8]. The difference between the RSE and the signal recycling is explained in detail in a paper such as [9]. A PRZD RSE operation was not demonstrated previously.

We have developed a novel control scheme for LCGT [9], and have carried out experimental work [10]. This experiment aims to control a PRZD RSE interferometer using the scheme to establish a lock acquisition sequence, to measure the sensing matrix, to compare it with modeling, and is focused on sensing-signal separation issues.

2. Control scheme

The control scheme consists mainly of two parts; the signal extraction scheme and the lock sequence.

2.1. Signal extraction scheme

The power recycled RSE interferometer has five DOFs to be controlled as shown in figure 1. They are the average length and the differential length of the two Fabry–Perot (FP) arm cavities, L_+ and L_- , respectively, as indicated by arrows with solid lines, the average length and the differential length of the power recycling cavity (PRC), l_+ and l_- , respectively, as indicated by arrows with short dashed lines, and the average length of the signal extraction cavity (SEC), l_s , as indicated by arrows with long dashed lines.

It is known from experience with present detectors, which use the Pound–Drever–Hall (PDH) method [11], that the arm cavities are relatively easy to control with clean control signals as the arm cavities have high finesse. On the other hand, it is expected to be quite challenging to obtain clean control signals of the central part of the RSE (i.e. l_+ , l_- and l_s), because the resonant conditions of the light fields inside the central part will be strongly affected by both the PRC and the SEC. Thus, the length-sensing scheme has to be designed in such a way that it manipulates the resonant conditions of the light inside by properly designed cavity lengths and sideband frequencies.

The outline of the length-sensing scheme is as follows. It is based on the PDH method. The FP arm cavity lengths are controlled with a single modulation–demodulation technique and the central part of the RSE is controlled with a double modulation–demodulation technique with amplitude modulation (AM) sidebands and phase modulation (PM) sidebands. By using the double modulation–demodulation technique, the control signals for the central part will

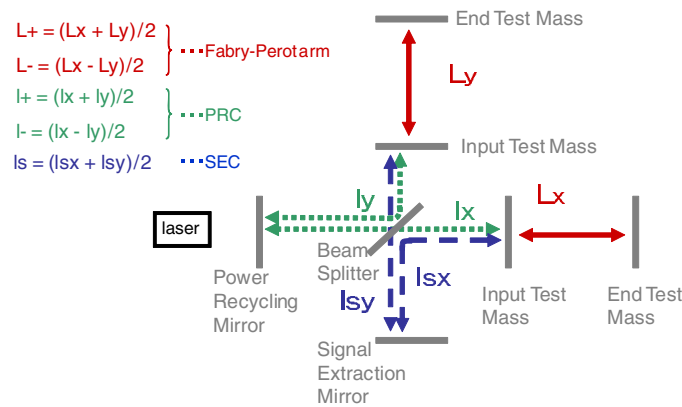


Figure 1. Degrees of freedom in the power recycled RSE interferometer.

be affected very little by the signals derived from the carrier which are dominated by the FP arm cavities. This is because both the AM and the PM modulation sideband frequencies are designed in such a way that they are not resonant in the FP arm cavities; thus, it decouples the FP arm cavities and the central part.

2.1.1. The central part. The central part is designed so that the AM and the PM sidebands behave in the following way. The lengths of two paths that compose the Michelson interferometer have a macroscopic asymmetry such that when the carrier interferes destructively at the dark port (DP), the AM sidebands interfere constructively at the bright port (BP) and destructively at the DP, while the PM sidebands interfere destructively at the BP and constructively at the DP. Thus, the AM sidebands ‘reflect completely’ from the Michelson part while the PM sidebands ‘transmit completely’ through the Michelson part. This condition is met when the round trip Michelson asymmetry length is designed to be equal to $(2m + 1)c/2f$ ($m = 0, 1, 2, \dots$) for the PM sidebands, where c is the speed of light and f is the modulation frequency, and integer multiple of c/f for the AM sidebands. In our design it is $3c/f$ for the AM sidebands and $c/2f$ for the PM sidebands. The two cavities’ macroscopic lengths are designed so that the AM sidebands resonate inside the PRC and the PM sidebands resonate inside the compound cavity made of the PRC and the SEC. This enables the PM sidebands to be sensitive to the length of the SEC while the AM sidebands are not affected by the SEC length, thus ensuring independent control signals for l_+ and l_s .

2.1.2. The whole RSE. Figure 2 shows how the control signals of the prototype RSE interferometer are obtained. The L_+ control signal is obtained at BP and is fed back to the end test masses (ETM), the L_- control signal is obtained at DP and is fed back to the ETMs. The l_- control signal is obtained at DP and is fed back to the beam splitter (BS), the l_+ control signal is obtained at the BP and is fed back to the power recycling mirror (PRM), and the l_s control signal is obtained at the pick-off port (PO) and is fed back to the signal extraction mirror (SEM).

Table 1 shows the theoretical length-sensing signal matrix. The signals are at dc. It is calculated with parameters used for the prototype RSE interferometer, of which there will be a detailed explanation in section 3. Demodulation phases are chosen so that each main signal (i.e., the signal that should be obtained at the corresponding detection port) is maximized. The first from the left column shows the detection ports; SD and DD stand for single/double

In reality, the uncontrolled interferometer is not near its operating point most of the time, therefore feedback signals for all the DOFs are almost never present simultaneously. In order to successfully lock the interferometer, a sequence of locking needs to be established. Here we show one possible order of locking the whole interferometer with simulation work results. FINESSE [12] is used for the calculation. The sequence is as follows:

- (i) Lock the central part
 - (a) Lock l_- DOF
 - (b) Lock l_+ DOF
 - (c) Lock l_s DOF
- (ii) Lock the FP arm cavities.

Figure 3 shows the feedback signal of the l_- DOF when the other two DOFs in the central part are free of control. Here the FP arms are not assumed to be present for simplicity, i.e. the FP arms are replaced with the input test masses of the FP arms. The plots are contour plots of the error signal for the l_- DOF when the other two DOFs are uncontrolled, i.e. the microscopic position of the PRM expressed in the phase gained by the light (hereinafter called ϕ) is scanned from its operating point by $0, 30, 60$ and 90° as indicated by the vertical arrow, while the SEM's ϕ is scanned from its operating point by -90 to 90° in the direction of the y -axis. The x -axis shows the ϕ of the l_- DOF, with the $\phi = 0$ being the operating point. On the right of each plot there is a color bar that shows the signal amplitude.

The l_- DOF can be locked to where the amplitude of the error signal changes its sign at the so-called zero crossing. A clear vertical line of a zero crossing is present at the operating point in the top figure when the l_+ and the l_s DOFs happen to be at their operating points, (i.e. their ϕ 's are zero). The polarity of the zero crossing needs to be taken into account; in this case, the right polarity is the one that goes from plus to minus as the ϕ is moved from minus to plus. With the wrong polarity the error signal does not guide the DOF to be locked at the operating point. A zero-crossing line at the operating point having the wrong polarity is seen as the ϕ of the l_+ DOF is scanned between 0 and 90° . One such example is marked by the ellipse in the figure. According to the simulation results the chance of such a line being present is $\sim 30\%$ at the maximum. This indicates that most of the time (with a minimum chance of being $\sim 70\%$), when the other two DOFs are completely free of control (i.e. their detunings are scanned independently from each other), the l_- DOF can be locked at its operating point.

Then, with the assumption that the l_- is now locked at its operating point, the feedback signal of the l_+ DOF is shown in figure 4 when the l_s is still free of control so it could be anywhere in the y -axis. In the same manner, as in figure 3, it is clear that a zero crossing for the l_+ is present very close to its operating point regardless of the l_s 's ϕ detuning. Therefore, the l_+ DOF can be locked at its operating point. When this is the case, the zero-crossing line of the l_- DOF is always present as shown in the top plot of figure 3, so the lock of the l_- DOF becomes stable.

Next, along with the assumption that the l_- is locked, there is a new assumption that the l_+ is locked at its operating point; a feedback signal of the l_s DOF is shown in figure 5. The figure shows that a zero-crossing point is present at its operation point as marked by the arrow. In this way, the central part is locked at its operating point.

This is a good approximation because the chance that the carrier is resonant inside the FP arm cavities, and thus disturbs the central part length-sensing signals is very small because of the high finesse, and the chance that the sidebands are accidentally resonant inside the FP arm cavities is very small for the same reason.

Next, the FP arms are locked by acquiring the lock of the individual FP arm, whose control signal is always present and dominant at each detection port for each FP arm cavity

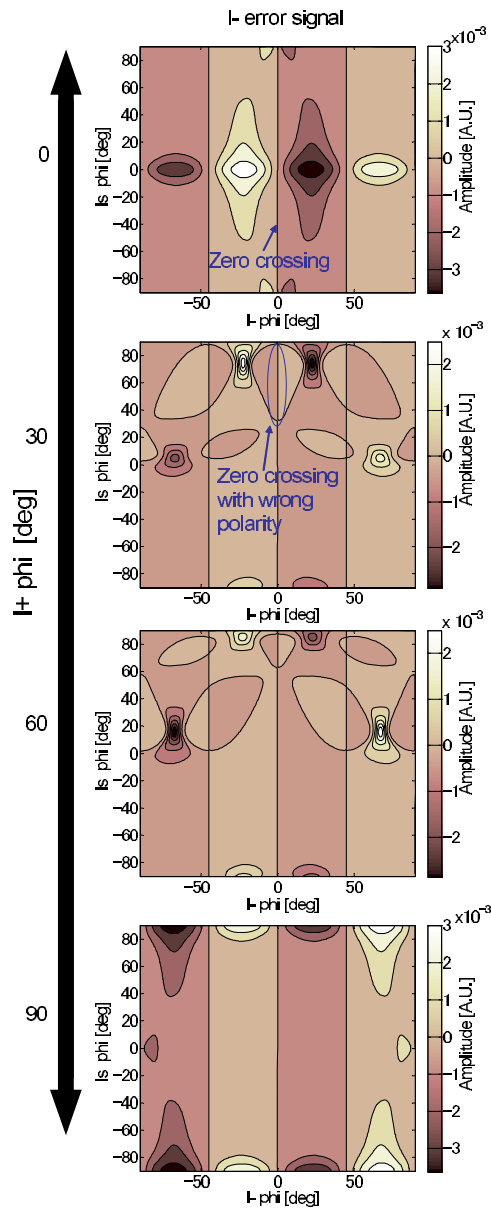


Figure 3. L_- error signal with SEM and PRM free.

which is detecting the reflected light from each cavity⁸. Once the two FP arm cavities are locked, the control servos can be switched to those of the L_- and L_+ DOF to change them into the controls of L_- and L_+ DOFs simultaneously. Therefore the FP arms are locked. Thus, the whole RSE interferometer can be locked in this order.

⁸ These detection ports are not shown in figure 2 for simplicity.

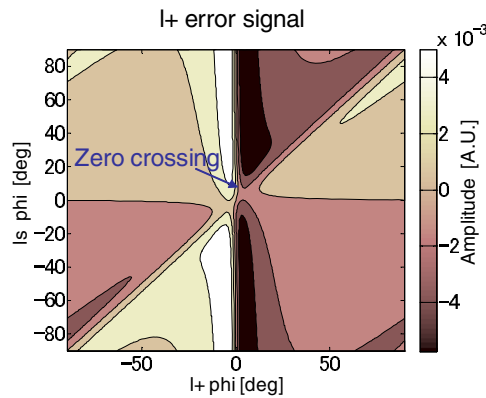


Figure 4. I_+ error signal with SEM free and I_- in lock.

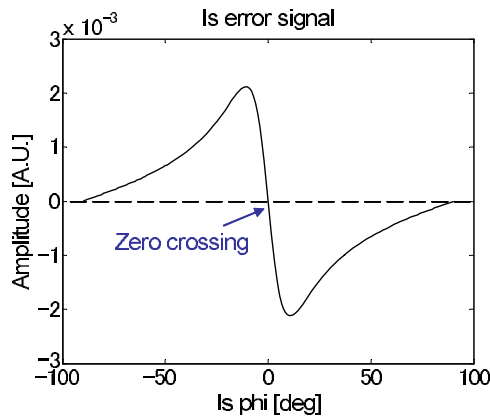


Figure 5. I_s error signal with I_+ and I_- in lock.

3. Experimental demonstration of a control of the PRZD RSE

3.1. 4 m prototype interferometer

Figure 6 shows the optical layout of the prototype interferometer. A light at $1.064 \mu\text{m}$ from a Nd:YAG laser enters a Mach-Zehnder interferometer (MZ) where two Pockels cells are placed in different paths of the MZ, and it is phase modulated at 17.25 MHz and amplitude modulated at 103.25 MHz. Light power at the output of the MZ is 90 mW with modulation depths of approximately 0.15 rad for both frequencies. The FP arm cavity length is 4.15 m, and is formed by a flat input mirror and a curved end mirror with a radius of curvature of 6 m. Two mode-matching lenses are used for the FP arm cavities, and an additional two are used to compensate for the mismatch due to the relatively large Michelson macroscopic length asymmetry (4.35 m). A coupling ratio between the incident beam and the fundamental mode of the FP arm cavity of more than 98% is achieved. The finesse of the cavity is approximately 120. Although this is about one order of magnitude smaller than the designed value of LCGT, it does not affect the lock acquisition strategy in terms of signal separation. All of the seven

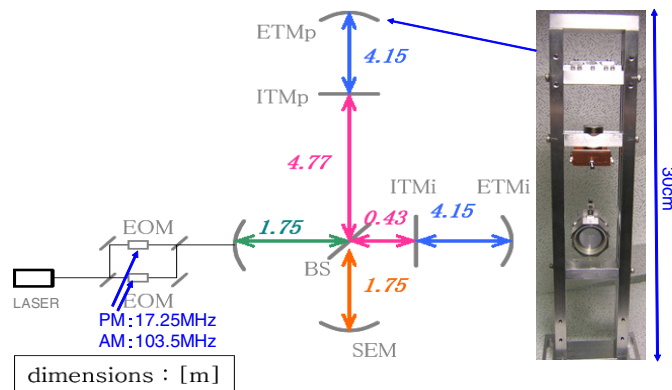


Figure 6. Optical layout of the prototype interferometer.

test masses (diameter 2.54 cm) are suspended as double pendulums (height 30 cm) to suppress the mirror motion at frequencies above the resonant frequency (at about 1.5 Hz). The test mass motion around its resonant frequency is damped by an eddy-current damping system. The length control is performed via coil-magnets actuators.

3.2. Lock sequence

The lock of the RSE has been successfully demonstrated. The sequence that has been determined with the simulation work is used; the central part first then the FP arm cavities. The direct lock of the FP arm cavities' two DOFs, the common, L_+ , and the differential, L_- , has been realized without an additional step (locking the individual arms with each control signal, then switch the control servos for those of the common and the differential control). This is because often one of the two FP arm cavities is locked with the control signal for the L_+ DOF first, and as soon as the other FP arm cavity's DOF is close to its operating point the L_- DOF is controlled, consequently switching the L_+ control signal to lock the L_+ DOF. Under quiet conditions the cavity mirrors' peak-to-peak motion is on the order of a little less than half wavelength, and hence the cavity goes through about one resonance per second. Typically, the time interval between the lock of one FP arm cavity and both cavities is less than a second, indicating that the movement of the second cavity during the locking event is enhanced by the L_+ feedback force.

We define several lock states, which are shown in figure 7. Each state is indicated as follows:

- State 0: the interferometer is uncontrolled.
- State 1: Michelson is locked to dark fringe at the DP.
- State 2: power-recycled Michelson is locked.
- State 3: central part is locked.
- State 4: tuned RSE is locked.

In state 2, the carrier is anti-resonant inside the PRC. In order for the carrier to be resonant inside the PRC, the phase needs to be shifted by π inside the PRC, which happens when the carrier is resonant inside the FP arm cavities as in state 4. Figure 8 shows the dc power detected at various ports (i.e., DP, BP, PO, transmitted port for the inline FP arm cavity and for the perpendicular FP arm cavity). Each lock state is separated by boxes with colors specified

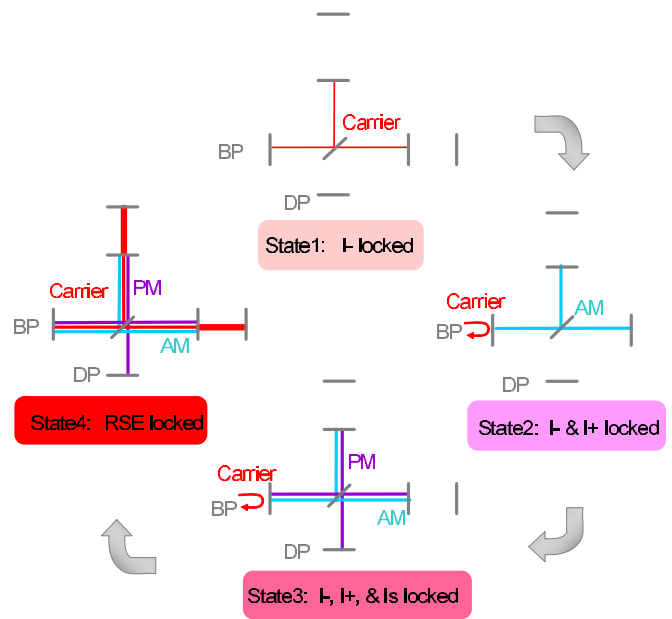


Figure 7. Lock states of the interferometer.

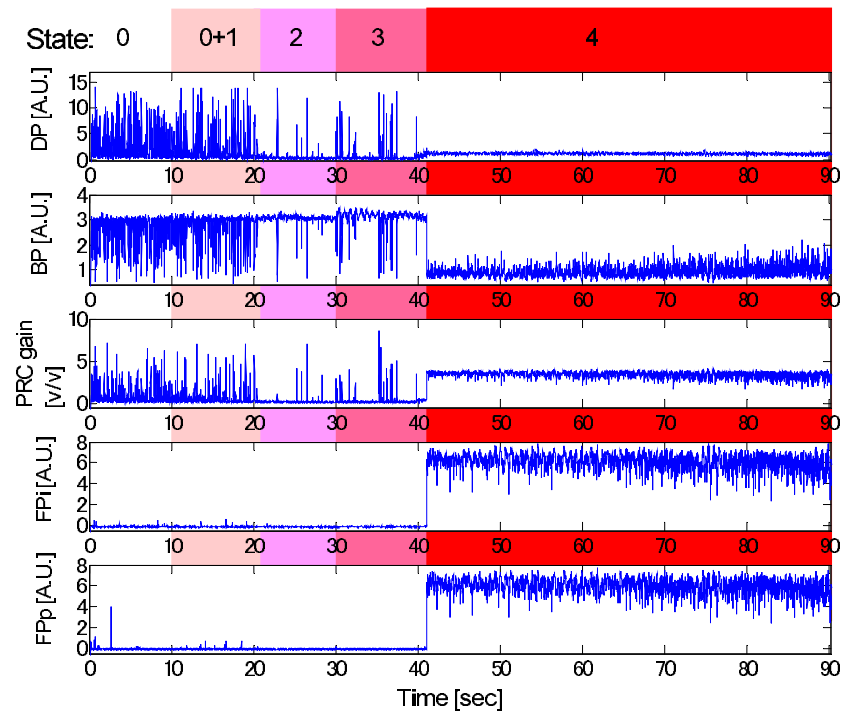


Figure 8. DC power at various ports.

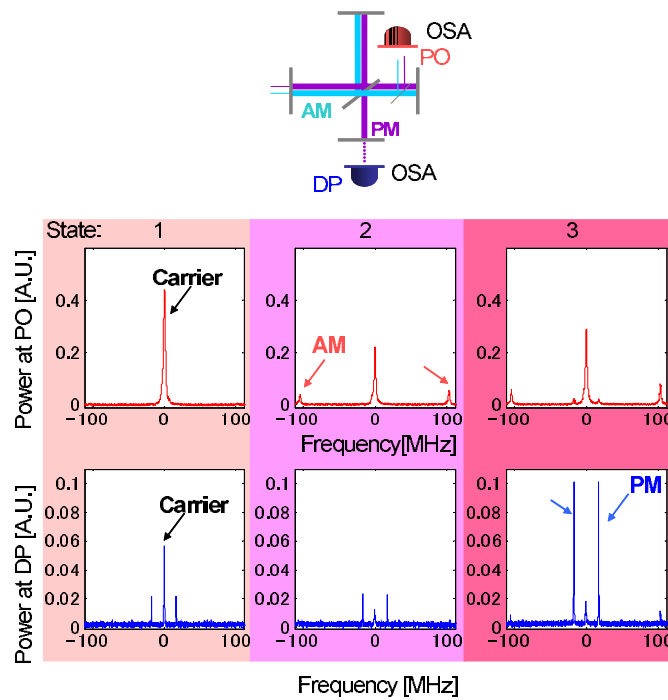


Figure 9. Sideband resonant peaks inside the PRC and the SEC.

in figure 7, and on top of each box the state number is shown. The start time of each state is defined as the time when the switch of a servo loop of each DOF is turned on. Typically, each DOF is locked within a fraction of a second after the servo loop is switched on. Between the time 10 and 20 s, the l_- DOF is not locked the whole time. This is because the carrier is not yet anti-resonant inside the PRC and disturbs the control signal of the l_- DOF. When the l_+ servo loop is switched on at time 20 s, the carrier is anti-resonant inside the PRC and the l_- DOF is completely locked as well as the l_+ DOF. The difference between states 2 and 3 is not obvious from the dc signals, due to the relatively small size of the modulation sidebands compared to that of the carrier. The sharp peaks seen between time 20 and 40 s can be explained as follows: the carrier light is sometimes resonant in one of the two arms, consequently adding a relative phase shift of π between the two beams interfering at the BS. This will switch the DP into BP, and vice versa, so the DP is no longer dark. Note that the corresponding peaks in light transmitted from the FP arm cavities are too small to be seen in the plots. Due to the fact that the double modulation–demodulation scheme always yields the same sign for the error signal irrespective of the dark port condition together with the large stable range of the control servo, the lock of the central part easily survives the short reversal from DP to BP.

In order to verify the locking status, resonant conditions for the two sets of sidebands are monitored with optical spectrum analyzers (OSA) placed at the DP and the PO. Figure 9 shows the output of the optical spectrum analyzers. The upper three plots show the output power at the PO, and the bottom three show the output power at the DP. Colored boxes indicate the lock state as specified in figure 7. In state 1, neither the AM nor the PM sidebands are resonant. In state 2, the AM sidebands are resonant inside the PRC, thus there are resonant peaks of the AM sidebands detected at the PO. The PM sidebands are not yet resonant. In state 3, the PM sidebands are resonant inside the compound cavity made of the PRC and the SEC; thus, there

Table 2. Measured normalized matrix.

	L_+	L_-	l_+	l_-	l_s
BP(SD)	1	4.5×10^{-2}	4.8×10^{-2}	-1.9×10^{-2}	-1.3×10^{-2}
DP(SD)	-1.9×10^{-1}	1	-2.1×10^{-3}	-4.9×10^{-2}	-1.5×10^{-3}
BP(DD)	-1.2×10^{-1}	6.6×10^{-3}	1	1.8×10^{-2}	-5.0×10^{-2}
DP(DD)	-1.1×10^{-1}	-3.1×10^{-1}	1.1×10^{-1}	1	1.2×10^{-1}
PO(DD)	-1.1	2.2×10^{-1}	4.0×10^{-1}	5.1×10^{-1}	1

are resonant peaks of the PM sidebands detected at the DP. Each cavity length is controlled throughout the locking process.

Table 2 shows the measured length-sensing signal matrix, normalized by the main signals at each detection port. Demodulation phases are tuned to maximize the main signal at each detection port. A sinusoidal signal is applied to a feedback path of each DOF to move the test masses. The signal frequency (at 2.2 kHz) is well above the unity gain frequency of all the control servo so that loop gains do not have to be taken into account. The frequency is also well below the cut-off frequency of the optical response of the RSE (at about 1 MHz), so the values can be directly compared with the theoretical dc values.

The general pattern of the measured matrix is in good agreement with the theoretical one; the $L_{+/-}$ signals dominate at their detection ports, the l_+ and the l_s mix at BP and PO but the signals are linearly independent, and the $L_{+/-}$ signals that mix at DD systems are relatively small. In general, unwanted signals other than the one that should be obtained at the corresponding detection ports tend to be greater than the theoretical values due to present imperfections such as the unwanted carrier light at the DP. It is not essential for this experiment, where agreement in the overall feature of the matrix is of importance in order to verify the control scheme.

4. Results and conclusion

A new control scheme has been developed for LCGT, and a PRZD RSE interferometer has been successfully controlled with the scheme with the prototype interferometer. It is the first experimental demonstration of a PRZD RSE interferometer with suspended test masses. In parallel the control sequence has been established with a simulation. The measured signal matrix showed good agreement with modeling. This result has shown that the LCGT can use a PRZD RSE interferometer as its optical configuration and use our control scheme to lock the interferometer. This is one important step toward the successful operation of the LCGT detector.

Acknowledgments

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