

by R. Weiss
(received 03/14/89, E.J.F)

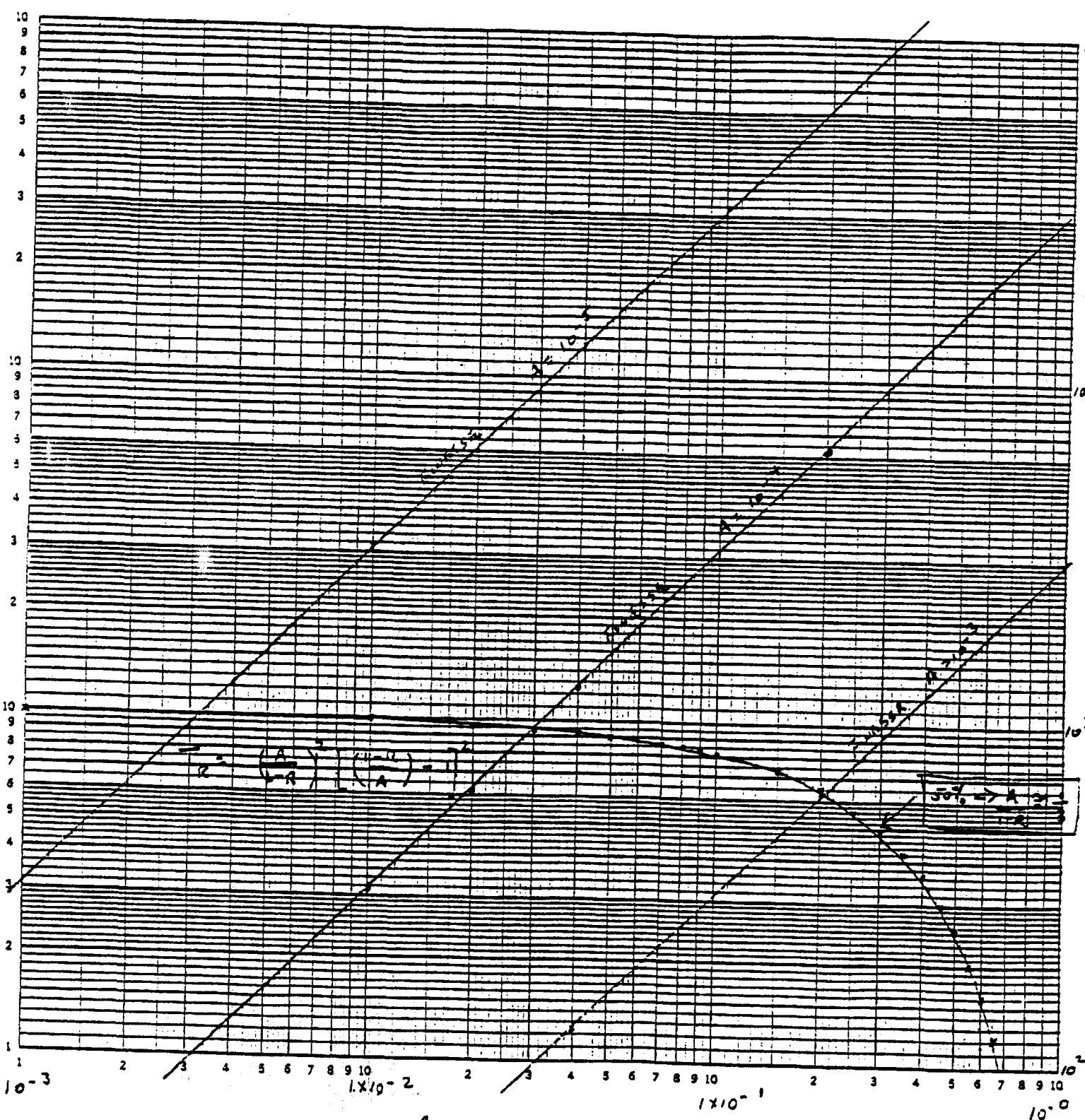
Finesse and transmission of a symmetric Fabry-Perot cavity

$R_1 = R_2 \quad A_1 = A_2$

$(R_1, A_1 \quad R_2, A_1)$

Finesse = $\frac{\nu_{fs}}{\Delta\nu(1/2\text{power})} = \frac{\pi R^{1/2}}{1-R}$

$\nu_{fs} = \frac{c}{2L} = 150\text{MHz/meter}$



Fabry-Perot properties

Assumptions:

- 1) A single radial mode is excited
- 2) All calculations made at input which is a wavefront with $R \rightarrow \infty$; for example a plane mirror of a hemispherical cavity.

Plan of the calculation

- 1) Determine Green's function (useful for transient analysis later)
- 2) Determine response to sinusoidal excitation
- 3) Show familiar intensity formula
- 4) Phase and amplitude of electric field reflection coefficient

<p>Amplitude coefficients $z = 0$</p> <p style="margin-left: 40px;">$\rightarrow r_1$</p> <p style="margin-left: 40px;">$\leftarrow -r_1$</p> <p style="margin-left: 40px;">$\rightarrow t_1$</p>	<p>$z = \ell$ $r = \frac{n' - n}{n' + n}$</p> <p style="margin-left: 40px;">$\leftarrow -r_2$</p> <p style="margin-left: 40px;">$\rightarrow t_2$</p> <p style="margin-left: 40px;">$t = \frac{2n}{n' + n}$</p>
---	--

Green's function: An impulse applied at input at time t'

$$\epsilon_{\text{inc}}(t) = \epsilon_0 \delta(t')$$

produces an exponentially decaying set of pulses at both $z = \ell$ (transmitted) and $z = 0$ (reflected) pulses

$$\begin{aligned}
 h_T(t) &= \frac{E_T(t)}{E_0} = t_1 t_2 \delta\left(t - \frac{\ell}{c}\right) + t_1 r_2 r_1 t_2 \delta\left(t - \frac{3\ell}{c}\right) + t_1 r_2 r_1 r_2 r_1 t_2 \delta\left(t - \frac{5\ell}{c}\right) + \dots \\
 &= t_1 t_2 \sum_{n=0}^{\infty} (r_2 r_1)^n \delta\left(t - \frac{(2n+1)\ell}{c}\right) \\
 h_R(t) &= \frac{E_R(t)}{E_0} = r_1 \delta(t) - t_1 t_1 r_2 \delta\left(t - \frac{2\ell}{c}\right) - t_1 r_2 r_1 r_2 t_1 \delta\left(t - \frac{4\ell}{c}\right) \\
 &\quad - t_1 r_2 r_1 r_2 r_1 r_2 t_1 \delta\left(t - \frac{6\ell}{c}\right) + \dots \\
 &= r_1 \delta(t) - t_1 t_1 r_2 \sum_{n=0}^{\infty} (r_1 r_2)^n \delta\left(t - \frac{(2n+2)\ell}{c}\right)
 \end{aligned}$$

The response of the system for an arbitrary incident electric field is the convolution of the Green's function with the incident field. The electric field at $z = 0$ (only the reflected part) will be the convolution of the Green's function with the incident electric field at $z = 0$

t is now, t' is the time associated with the incident field and τ is the delay time $t' = t - \tau$

$$\begin{aligned} \frac{E_R(t)}{\text{field now}} &= \int_{-\infty}^0 h_R(\tau) E_{\text{inc}}(t - \tau) d\tau && \text{reflected field at } z = 0, t \end{aligned}$$

The transmitted field at $z = \ell$ is then

$$\frac{E_T(t)}{\text{field now}} = \int_{-\infty}^0 h_T(\tau) E_{\text{inc}}(t - \tau) d\tau \quad \text{transmitted field at } z = \ell, t$$

The transfer function for-reflection or transmission is the response for a sinusoidal incident field

$$E_{\text{inc}}(t') = E_0 e^{i\omega t'}$$

$$\frac{E_r(t)}{E_0} = e^{i\omega t} \left[r_1 - t_1 t_1 r_2 e^{-i\frac{2\omega\ell}{c}} \sum_{n=0}^{\infty} (r_1 r_2)^n e^{-i\frac{2\omega n\ell}{c}} \right]$$

Since $|(r_1 r_2)| < 1$ the series can be summed

$$= e^{i\omega t} \left[r_1 - \frac{t_1 t_1 r_2 e^{-i\frac{2\omega\ell}{c}}}{(1 - r_1 r_2 e^{-i\frac{2\omega\ell}{c}})} \right] = e^{i\omega t} \left[\frac{r_1 - r_2 (r_1^2 + t_1^2) e^{-i\frac{2\omega\ell}{c}}}{(1 - r_1 r_2 e^{-i\frac{2\omega\ell}{c}})} \right]$$

Summing series for transmitted beam gives

$$\frac{E_T(t)}{E_0} = e^{i\omega t} \left[\frac{t_1 t_2}{1 - r_1 r_2 e^{-i\frac{2\omega\ell}{c}}} \right]$$

Check for familiar intensity formulae

power
coefficients
 $T_i = t_i t_i^*$
 $R_i = r_i r_i^*$

$$\begin{aligned} \frac{I_T}{I_{\text{inc}}} &= \frac{E_T^* E_T}{E_0^2} = \frac{T_1 T_2}{1 + R_1 R_2 - 2\sqrt{R_1 R_2} \cos \frac{2\omega\ell}{c}} \\ \frac{I_R}{I_{\text{inc}}} &= \frac{E_R^* E_R}{E_0^2} = \frac{R_1 + (R_1 + T_1)^2 R_2 - 2\sqrt{R_1 R_2} (R_1 + T_1) \cos \frac{2\omega\ell}{c}}{1 + R_1 R_2 - 2\sqrt{R_1 R_2} \cos \frac{2\omega\ell}{c}} \end{aligned}$$

If $\begin{matrix} R_1 + T_1 = 1 \\ R_1 + T_2 = 1 \end{matrix}$ $I_T + I_R = I_{\text{inc}}$

Amplitude coefficient for reflection cavity

$$\frac{E_R}{E_{\text{inc}}}(x) = \frac{r_1 - r_1(r_1^2 + t_1^2) e^{-ix}}{(1 - r_1 r_2 e^{-ix})} \quad x = 2\omega l/c$$

$$r_1 = |r_1| e^{i\varphi_1}$$

$$r_2 = |r_2| e^{i\varphi_2}$$

π Phase reversal has been taken into account already

Approximation

Let r_1 and r_2 be real and take account of the losses in the amplitude; T_1 = front mirror transmission, A_1 , and A_2 = mirror loss

Approximate

$$r_1 \cong (1 - T_1 - A_1)^{1/2} \quad r_2 = (1 - A_2)^{1/2}$$

Define $A^2 = 1 - A_1 - T_1$ $B^2 = (1 - A_2)(1 - A_1)^2$ $C^2 = (1 - A_1 - T_1)(1 - A_2)$

$$\frac{E_R}{E_{\text{inc}}}(x) = \left[\frac{A^2 - 2AB \cos x + B^2}{1 - 2C \cos x + C^2} \right] e^{i \underbrace{\left[\tan^{-1} \left(\frac{B \sin x}{A - B \cos x} \right) - \tan^{-1} \left(\frac{C \sin x}{1 - C \cos x} \right) \right]}_{\varphi}}$$

$$\frac{d\varphi}{dx} = \frac{B}{A^2 + B^2 - 2B \cos x} [A \cos x - B] - \frac{C}{1 + C^2 - 2C \cos x} [\cos x - C]$$

Special but important case

$$A_1 \ll 1 \quad A_2 \ll 1 \quad T_1 < 1 \quad T_2 = 0$$

near resonance

φ	x	
π	0	on resonance
$\pi \pm \pi/2$	$\pm \frac{T_1 + A_1 + A_2}{2}$	1/2 power pt
$\pi \pm 3\pi/4$	$\pm T_1$	
$2\pi \pm T_{1/2}$	$\pm \pi/2$	
$2\pi, 0$	$\pm \pi$	1/2 free spectral range

$$\left. \frac{\partial \varphi}{\partial x} \right|_{x=0} = \frac{B - CA}{(A - B)(1 - C)} \cong - \frac{4}{T_1 \left(1 - \frac{(A_1 + A_2)^2}{T_1} \right)}$$

$$\left. \frac{\partial \varphi}{\partial x} \right|_{x=\pi/2} \cong -T_{1/2}$$

$$\left| \frac{E_R}{E_{\text{inc}}} \right|^2(x=0) = 1 - \frac{4(A_1 + A_2)}{T_1}$$

$$\left| \frac{E_R}{E_{\text{inc}}} \right|^2(x=\pi/2) = 1 - A_1$$

1/2 Power points

$$\Delta x = \pm \frac{T_1 + A_1 + A_2}{2} \text{ radians}$$

$$\left| \frac{E_R}{E_{inc}} \right|^2 (\Delta x) = 1 - \frac{2(A_1 + A_2)}{T_1}$$

phase $\varphi(\Delta x) = \pi \pm \frac{\pi}{2}$

Frequency and position sensitivity of reflection cavity

at $x = 0$

$$\Delta x = \frac{2\Delta\omega\ell_0}{c} + \frac{2\omega_0\Delta\ell}{c}$$

Frequency sensitivity $\Delta\ell = 0$ change in optical phase reflected/change in frequency

$$\left. \frac{\Delta\varphi}{\Delta f} \right|_0 = \left. \frac{\partial\varphi}{\partial x} \right|_0 \frac{dx}{df} = -\frac{16\pi}{T_1 \left(1 - \left(\frac{A_1 + A_2}{T_1}\right)^2\right)} \frac{\ell_0}{c}$$

In terms of energy storage time defined by a source internal to the cavity

$$\text{Energy stored in cavity} = u\ell_0 A_c$$

 A_c = beam cross section ℓ_0 = length of cavity u = energy density

Power dissipated or emitted

$$\frac{dE}{dt} = -\frac{cuA_c}{2} (T_1 + A_1 + A_2)$$

$$E(t) = E(0) e^{-t/\tau_{\text{store en}}}$$

$$\tau_{\text{store en}} = \frac{2\ell_0}{c(T_1 + A_1 + A_2)} = \frac{2\tau_{\text{trans}}}{(T_1 + A_1 + A_2)}$$

Using energy storage time

$$\tau_{\text{store field}} = 2\tau_{\text{store en}} = \frac{4\tau_{\text{trans}}}{(T_1 + A_1 + A_2)}$$

$$\frac{1}{T_1 \left(1 - \left(\frac{A_1 + A_2}{T_1}\right)^2\right)} = \frac{1}{2} \frac{\tau_c}{\tau_t} \frac{\left(1 - \frac{1}{2} \frac{\tau_c}{\tau_t} (T_1 + A_1 + A_2)\right)}{\left(1 - \frac{\tau_c}{\tau_t} (A_1 + A_2)\right)}$$

another formulation of energy storage time

$$\frac{\Delta\varphi}{\Delta f} = -8\pi \tau_{\text{sen}} \frac{\left(1 - \frac{1}{2} \frac{\tau_c}{\tau_t} (A_1 + A_2)\right)}{\left(1 - \frac{\tau_c}{\tau_t} (A_1 + A_2)\right)}$$

$$\tau_{\text{sen}} = \frac{\tau_t}{(1 - |\tau_1 \tau_2|)}$$

$$A_1 + A_2 \ll T_1$$

$$\frac{\Delta\varphi}{\Delta f} \cong -8\pi \tau_{\text{store en}} = -4\pi \tau_{\text{store field}}$$

 $\frac{1}{2}$ power pt full width

$$\Delta f\left(\frac{1}{2}\right) = \frac{c}{2\ell_0} \frac{(T_1 + A_1 + A_2)}{2\pi} = \Delta\nu_{fs} \frac{(T_1 + A_1 + A_2)}{2\pi} \quad \Delta\nu_{fs} = \text{free spectral range}$$

$$= \frac{1}{2\pi \tau_{\text{store en}}}$$

Position sensitivity

$$\frac{\Delta\varphi}{\Delta l} \cong \frac{\partial\varphi}{\partial x} \frac{\partial x}{\partial l} = \frac{-16\pi f_0/c}{T_1(1 - (\frac{A_1+A_2}{T_1})^2)} = \frac{-16\pi}{\lambda_0 T_1(1 - (\frac{A_1+A_2}{T_1})^2)}$$

In terms of the energy storage time

$$\frac{\Delta\varphi}{\Delta l} \cong \frac{-8\pi\tau_{\text{store en}}}{\lambda_0\tau_{\text{transit}}} \frac{(1 - \frac{1}{2}\frac{\tau_s}{\tau_t}(A_1 + A_2))}{(1 - \tau_s/\tau_t(A_1 + A_2))}$$

$$A_1 + A_2 \ll T_1$$

$$\frac{\Delta\varphi}{\Delta l} \cong \frac{-8\pi}{\lambda_0} \frac{\tau_{\text{store en}}}{\tau_{\text{transit}}} = -\frac{4\pi}{\lambda_0} \frac{\tau_{\text{store field}}}{\tau_{\text{transit}}}$$

Note: Maximum bandwidth for frequency or length control loops is determined by phase (φ) response going to $\varphi = \pi \pm \pi$. This corresponds to a frequency width of $1/2 \Delta\nu_{fs}$.

$$\Delta BW_{\text{max}} = \frac{c}{4L_0}$$

For LIGO with 4km arms

$$\Delta BW_{\text{max}} = 18.8 \text{ KHz}$$

Further properties of cavities

DC transfer function of a Fabry-Perot cavity to a length displacement

Again assume input mirror has R_1 T_1 A_1 Back mirror has A_2 $R_2 = 1 - A_2$

$$E_{in} = E_0 e^{i\omega t} \rightarrow$$

$$E_{ref} = E_0 A(\Delta k) e^{i\delta(x)} (\leftarrow \ell \rightarrow)$$

Phase change in reflection due to a change of length at resonance

$$x = 2\omega\ell/c$$

$$\left. \frac{d\varphi}{dx} \right|_{x=n2\pi} = -\frac{4}{T_1 \left(1 - \left(\frac{A_1 + A_2}{T_1}\right)^2\right)}$$

$$A^2(o) = 1 - \frac{4(A_1 + A_2)}{T_1}$$

For a length change $\Delta\ell$

$$\Delta\varphi = \frac{d\varphi}{dx} \frac{dx}{d\ell} \Delta\ell = -\frac{8\omega}{cT_1 \left(1 - \left(\frac{A_1 + A_2}{T_1}\right)^2\right)} = -\frac{16\pi}{T_1 \left(1 - \left(\frac{A_1 + A_2}{T_1}\right)^2\right)} \frac{\Delta\ell}{\lambda}$$

 $h = 2\frac{\Delta\ell}{\ell}$ strain to gravity wave amplitude

$$\frac{\Delta\varphi}{h} = -\frac{8\pi}{T_1 \left(1 - \left(\frac{A_1 + A_2}{T_1}\right)^2\right)} \frac{\ell}{\lambda} \text{ for one cavity}$$

Try to define storage time as when intensity drops by $1/e$

Begin with wave inside cavity

$$\begin{pmatrix} M_1 & M_2 \\ T_1 A_1 & A_2 \end{pmatrix}$$

$$\begin{aligned} \bar{S} &= \frac{c}{4\pi} (\bar{E} \times \bar{B}) & u &= \frac{1}{8\pi} (E^2 + B^2) \\ &= \frac{c}{4\pi} E^2 & &= \frac{1}{4\pi} E^2 \end{aligned}$$

The energy in the cavity is made up of a standing wave with $1/2$ the energy traveling $R \rightarrow L$ and the other $L \rightarrow R$

$$U_{total} = u_+ A\ell + u_- A\ell \quad u_+ = u = u_-$$

$$U_{total} = 2uA\ell$$

$$\text{flux out and into mirror 1 } S_+ A = v u A (T_1 + A_1)$$

$$\text{flux into mirror 2 } S_- A = c u A (A_2)$$

$$\left. \frac{dU}{dt} \right|_{\text{inside cavity}} = -(S_+ A + S_- A) = -c u A [T_1 + A_1 + A_2] = -\frac{c U_{total}}{2A\ell} A [T_1 + A_1 + A_2]$$

$$\frac{dU}{U_{\text{total}}} = -\frac{c}{2\ell}(T_1 + A_1 + A_2)dt$$

$$U(t) = U_0 e^{-\frac{c(T_1 + A_1 + A_2)}{2\ell}t} \quad \tau_{STU} = \frac{2\ell}{c(T_1 + A_1 + A_2)} = \frac{\tau_t}{(1 - |r_1 r_2|)}$$

$$E^2(t) = E_0^2 e^{-\frac{c(T_1 + A_1 + A_2)}{2\ell}t}$$

$$E(t) = E_0 e^{-\frac{c(T_1 + A_1 + A_2)}{4\ell}t} \quad \tau_{ST_{\text{field}}} = \frac{4\ell}{c(T_1 + A_1 + A_2)}$$

With this definition of storage time

$$\frac{\Delta\varphi}{h} = -\frac{4\pi r_s n c}{\lambda} \frac{\left(1 - \frac{1}{2} \frac{r_s}{r_t} (A_1 + A_2)\right)}{\left(1 - r_s/r_t (A_1 + A_2)\right)} \stackrel{\text{energy storage time}}{=} -\frac{2\pi r_s n c}{\lambda} \frac{\left(1 - \frac{1}{4} \frac{r_s n}{r_t} (A_1 + A_2)\right)}{\left(1 - \frac{1}{2} \frac{r_s n}{r_t} (A_1 + A_2)\right)} \stackrel{E \text{ field storage time}}{=}$$

one cavity
h → DC

OPTICAL INHOMOGENEITY

Optical loss when averaged over wavefront

$$A_{L_{w_0}} = (1 - \langle J_0(\Gamma_{\text{beam}}^{\text{rms}}) \rangle) = \frac{\Gamma_{\text{rms}}^2}{4} \text{ for small phase shift}$$

In terms of wavefront distortion in cm required to correct

φ_c contour of constant phase
average φ line

$$\Gamma_{\text{rms}} = \left(\frac{\int \delta x^2(r, \theta) A_{\text{mode dimension}} da}{\int A_{\text{mode dimension}} da} \right)^{1/2}$$

We do not know the distribution $\delta x^2(r, \theta)$. Optics industry gives average Δx over ill defined beam size.

Nevertheless using the average values gives some indication. Usually the specification is given in

$$\Delta x / \ell_{\text{path}} = Q$$

$$A_L = \pi^2 \left(\frac{Q \ell_{\text{path}}}{\lambda} \right)^2$$

Another way to specify is in terms of wave front distortion in units of the wavelength

$$\Delta x = f \lambda$$

$$A_L = \pi^2 (f^2)$$

A_L	f	$Q \text{ mm/cm } \ell = 10\text{cm } \lambda = 5145\text{\AA} \Delta n$
.1	1/10	5nm/cm 5×10^{-7}
.025	1/20	2.5
.011	1/30	1.7
.0062	1/40	1.25
.0040	1/50	1 1×10^{-7}