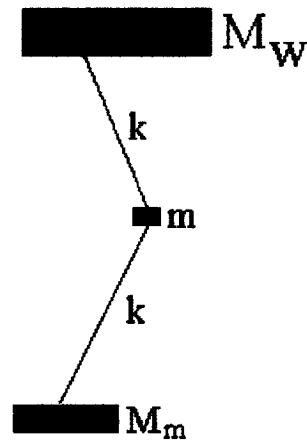


## Violin Mode Noise

Robert L. Coldwell and Bernard F. Whiting  
University of Florida  
Supported by NSF Grant PHY-9722114



### Introduction

A series of pulses in the driving force of a violin mode are in principle capable of reproducing any output data signal. This is a rather serious noise problem since an accidental series of such pulses can produce results that might be confused with a gravity wave. Fortunately the complete set needed to reproduce most signals is fairly large. An incomplete set has a signature in the form of spurious bumps in the data that do not correspond to the expected signal.

### The differential equation

So that the equation for the signal from a violin mode in the time domain is

$$\ddot{d}(t) + 4\pi w \dot{d}(t) + (2\pi f_0)^2 d(t) = \rho(t)$$

In the frequency domain this is

$$[-f^2 - 2jwf + f_0^2](2\pi)^2 D(f) = P(f)$$

In the case that  $\rho(t) = \delta(t - t_0)$ ,  $P(f) = \int_{-\infty}^{\infty} \delta(t - t_0) e^{j2\pi ft} dt = e^{j2\pi ft_0}$ , this yields

$$D(f) = \frac{-e^{j2\pi ft_0}}{(2\pi)^2 [f^2 + 2jwf - f_0^2]}$$

## Two pulses canceling the center of the natural peak

Suppose I have a pulse at  $t_0 = 0$  and a second at  $t_1$  sometime later. Then

$$D_2(f) = \frac{-(1 + c_1 e^{j2\pi f t_1})}{(2\pi)^2 [f^2 + 2jwf - f_0^2]} = \frac{-(1 + c_1 \cos(2\pi f t_1) + jc_1 \sin(2\pi f t_1))}{(2\pi)^2 [f^2 + 2jwf - f_0^2]}$$

The object is to make the central peak disappear. There are three parameters,  $\text{Re}(c_1)$ ,  $\text{Im}(c_1)$  and  $t_1$ . There are two conditions, that the real part and imaginary part of  $D_1(f_0)$  equal zero. If I choose  $c_1$  to be real, the imaginary condition is satisfied by

$$t_1 = \frac{m}{2f_0}, \quad m = \text{any integer} \quad . \quad \text{Then the cosine is given by}$$

$$\cos(\pi m) = (-1)^m \quad \text{so that the real part of } c_1 \text{ is given by}$$

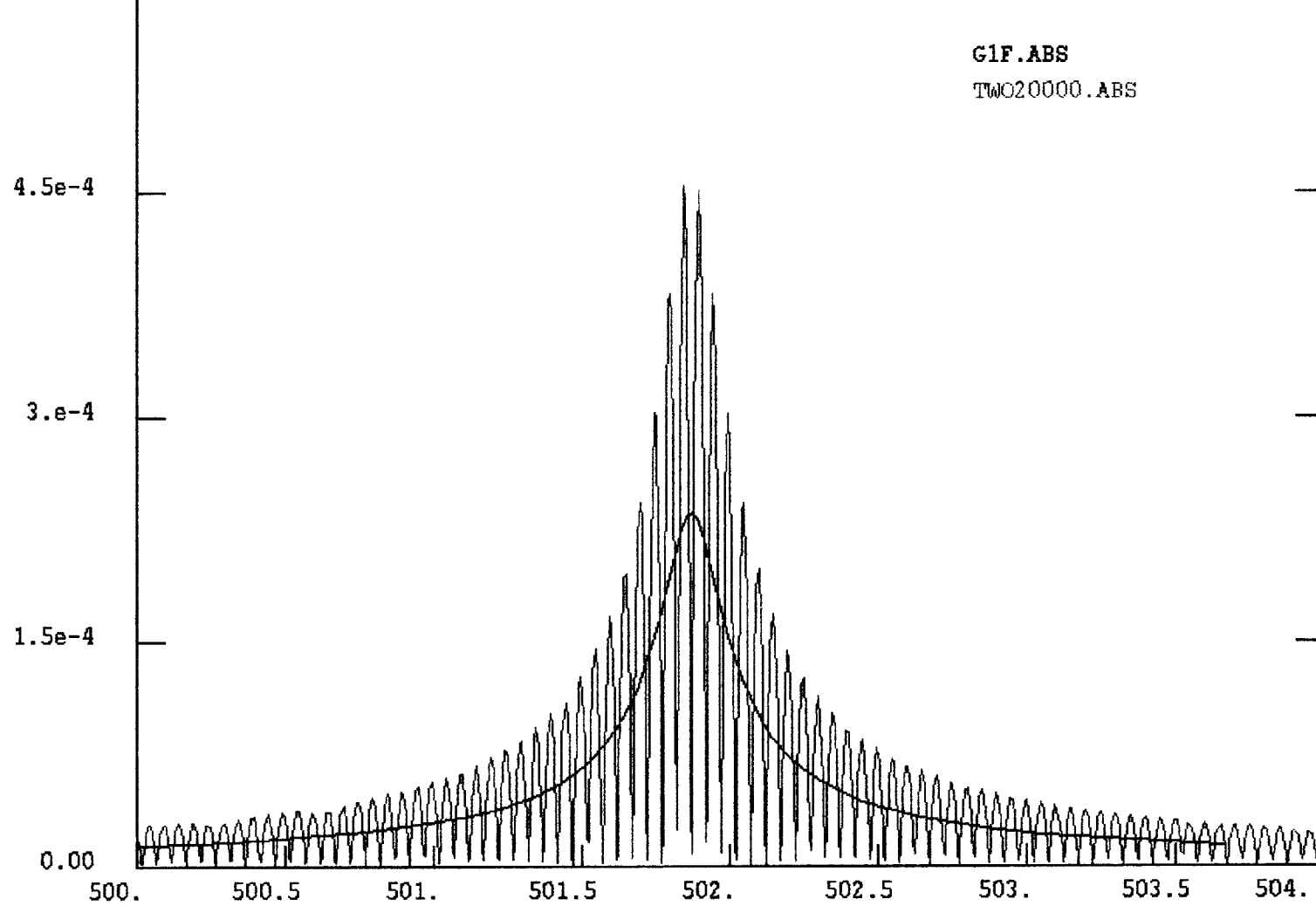
$$c_1 (-1)^m + 1 = 0; \quad c_1 = (-1)^{m+1} \quad . \quad \text{The value a short distance from } f_0 \text{ can be found from the expansion}$$

$$D_2(f) \cong D_2(f_0) + (f - f_0) \left. \frac{\partial D_2}{\partial f} \right|_{f=f_0} \cong j2\pi t_1 c_1 (f - f_0) (\cos(2\pi f_0 t_1)) G(j$$

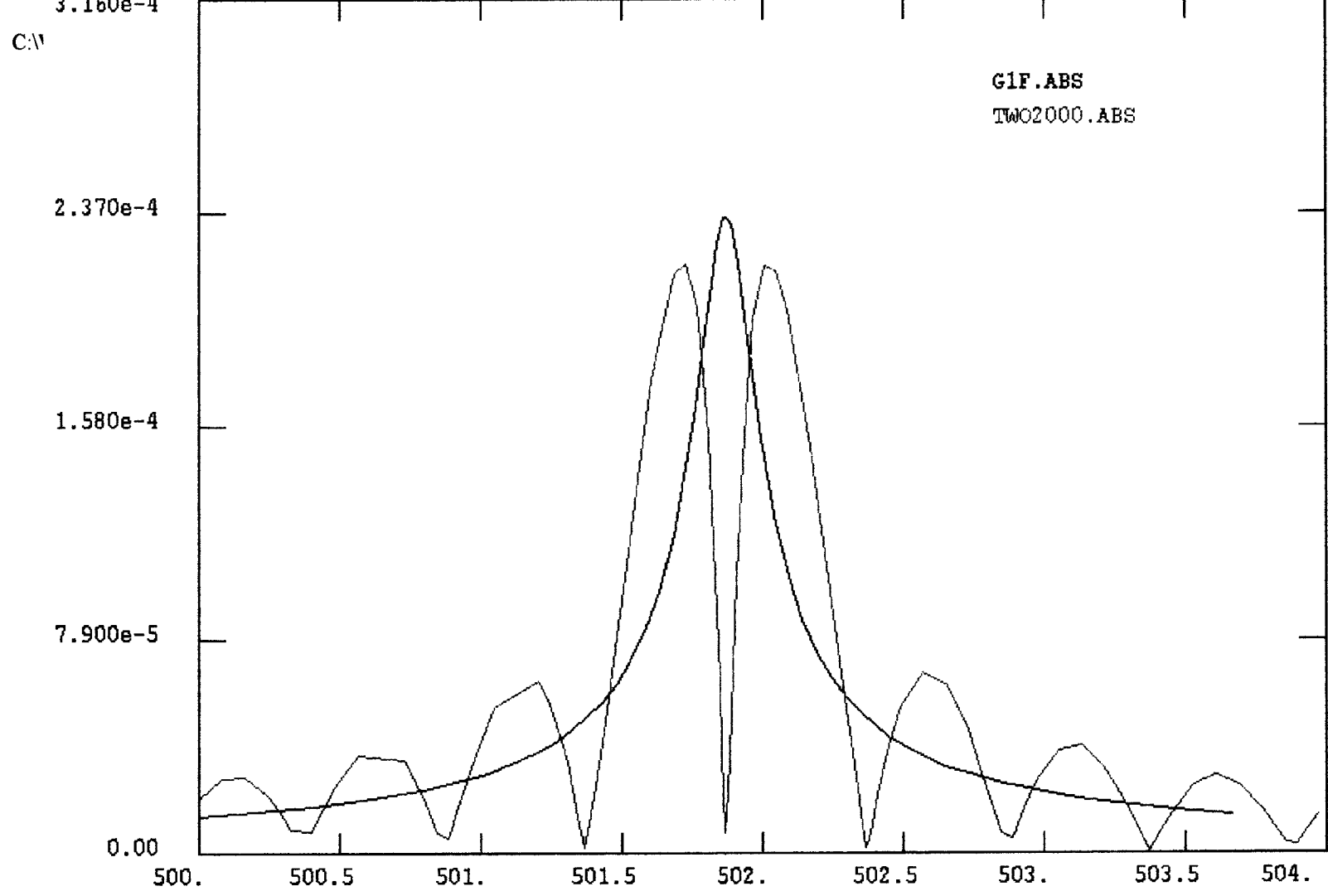
$$j2\pi t_1 c_1 (f - f_0) (-1)^m G(f_0) e^{-(f_0 w_t)^2} = -j2\pi m \frac{(f - f_0)}{2f_0} G(f_0)$$

The last term is the peak from a delta function response. The derivative of this with respect to  $f$  at  $f=f_0$  is zero owing to  $f_0$  being the value at the peak. Note that a two-term attempt to eliminate the peak will have the most success for small  $m$ . In fact  $m=0$  is a positive and negative peak at exactly the same time, which totally eliminates the peak for all values of  $f$ .

C

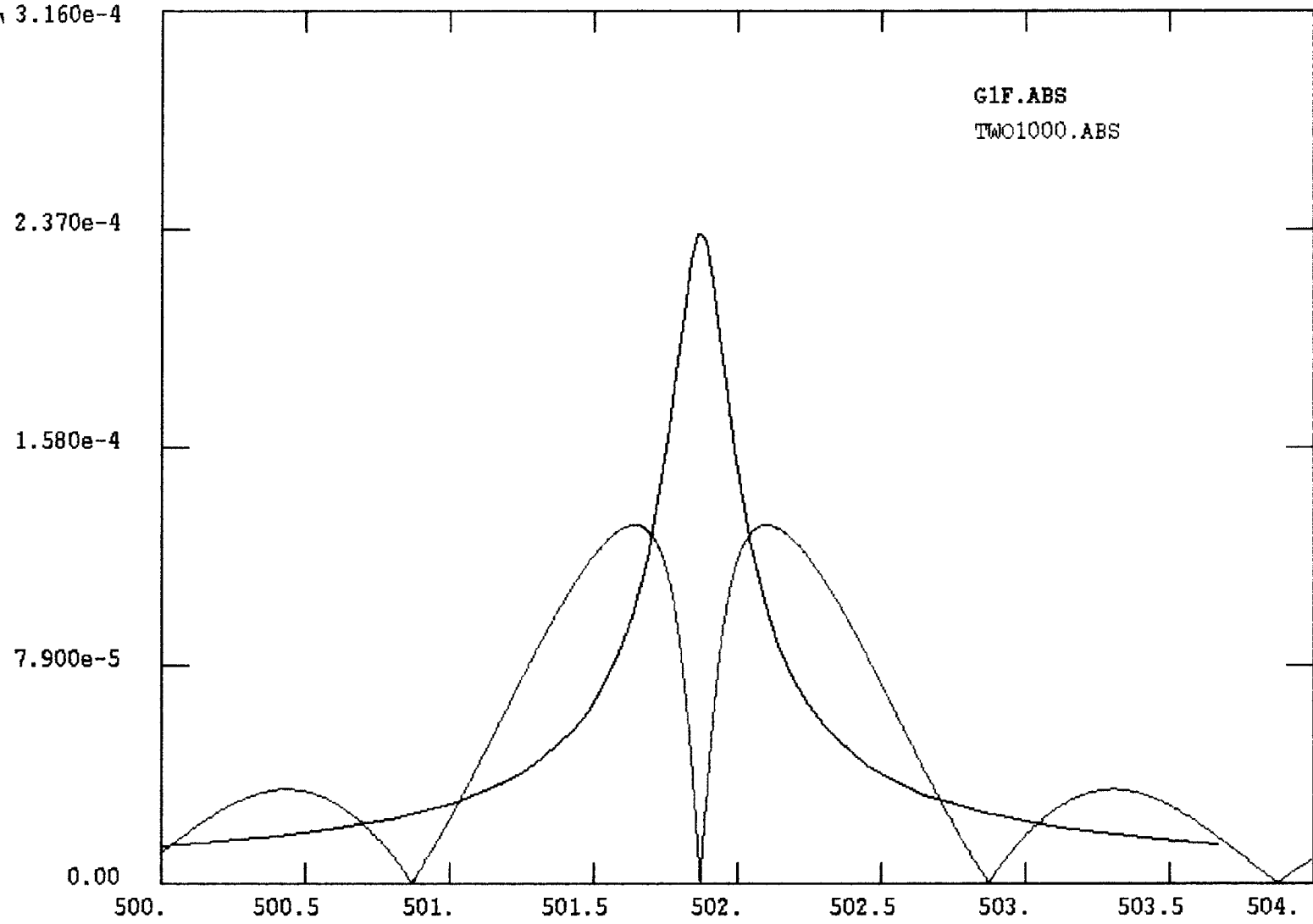


**Figure 1 Two pulses ~40 seconds apart. The peaks are narrow relative to the size of the natural line.**



**Figure 2 Two pulses ~4 seconds apart. The peaks are about the same width as those in the natural line.**

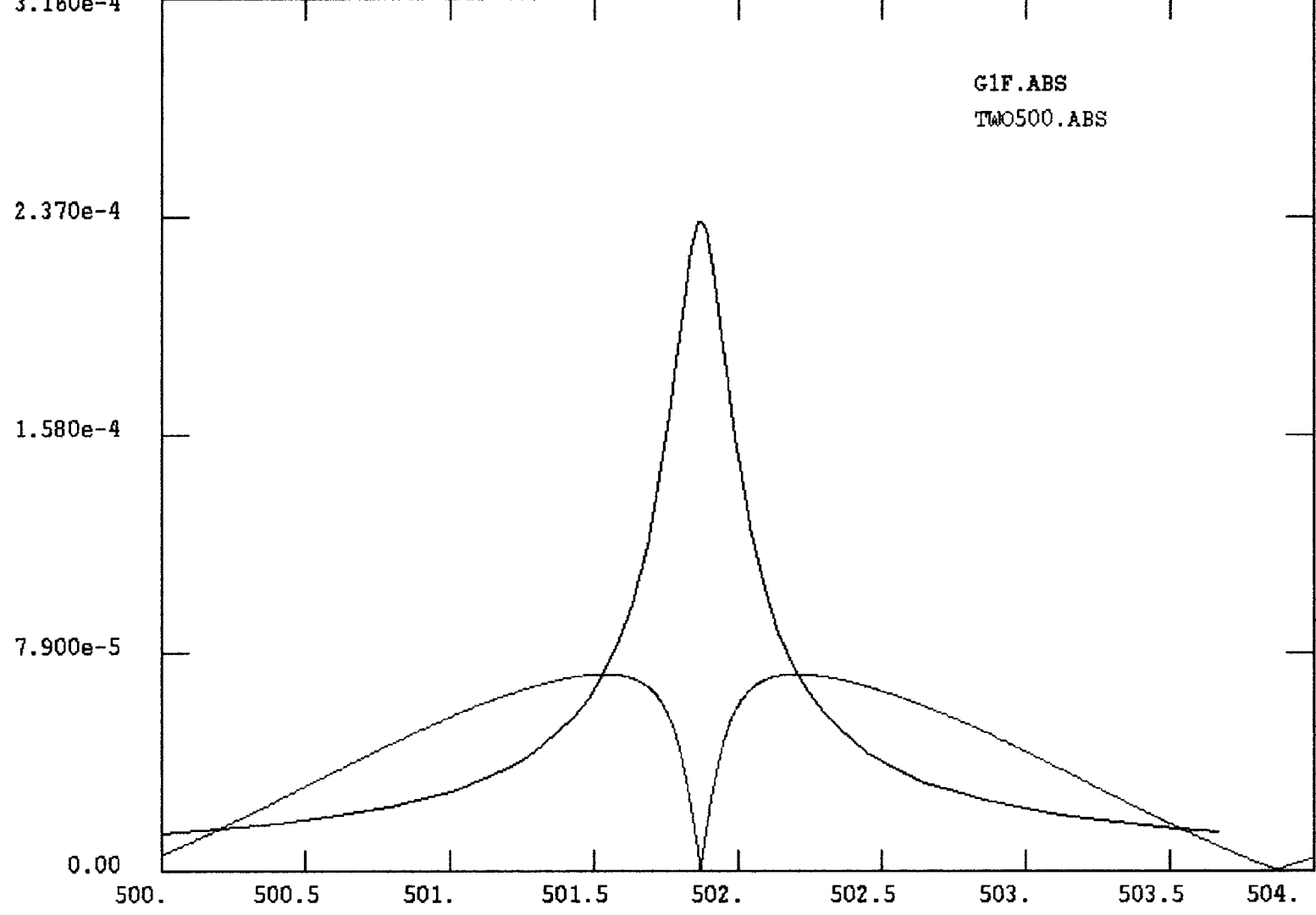
C:\ 3.160e-4



**Figure 3 Two pulses on the order of 2 seconds apart are shown. The peaks are much wider than the natural shape.**

The width of a pulse is the change in frequency required to go through  $\pi$  in the imaginary exponent. Thus  $\Delta f t_0 = 1$  or  $\Delta f = 1/t_0$ , with a natural width of 0.1 cycles per second, the time required to get to this width is  $t_0 = 10 \text{ sec}$ . Thus a good approximation to the peak may be a series of  $t_0$ 's between 1 and 10 seconds. If the times are allowed to adjust so that  $t_i = t_{0,i} + \delta_i$ , the adjustments will change the value of  $e^{j2\pi f(t_{i,0} + \delta_i)}$ . These changes approximately repeat for  $f\delta_i = 1$  meaning that the changes will be found to vary the value of  $\delta_i$  only by  $1/f$ .





**Figure 4 Two peaks ~1 second apart.**

## Finding $\rho(t)$

For any observed data  $D(f)$ , there exists a  $P(f)$  given by

$$P(f) = -(2\pi)^2 [f^2 + 2jwf - f_0^2] D(f)$$

### **Example 1 A displaced violin mode**

In particular suppose that  $D(f)$  is a violin mode at a frequency other than the original

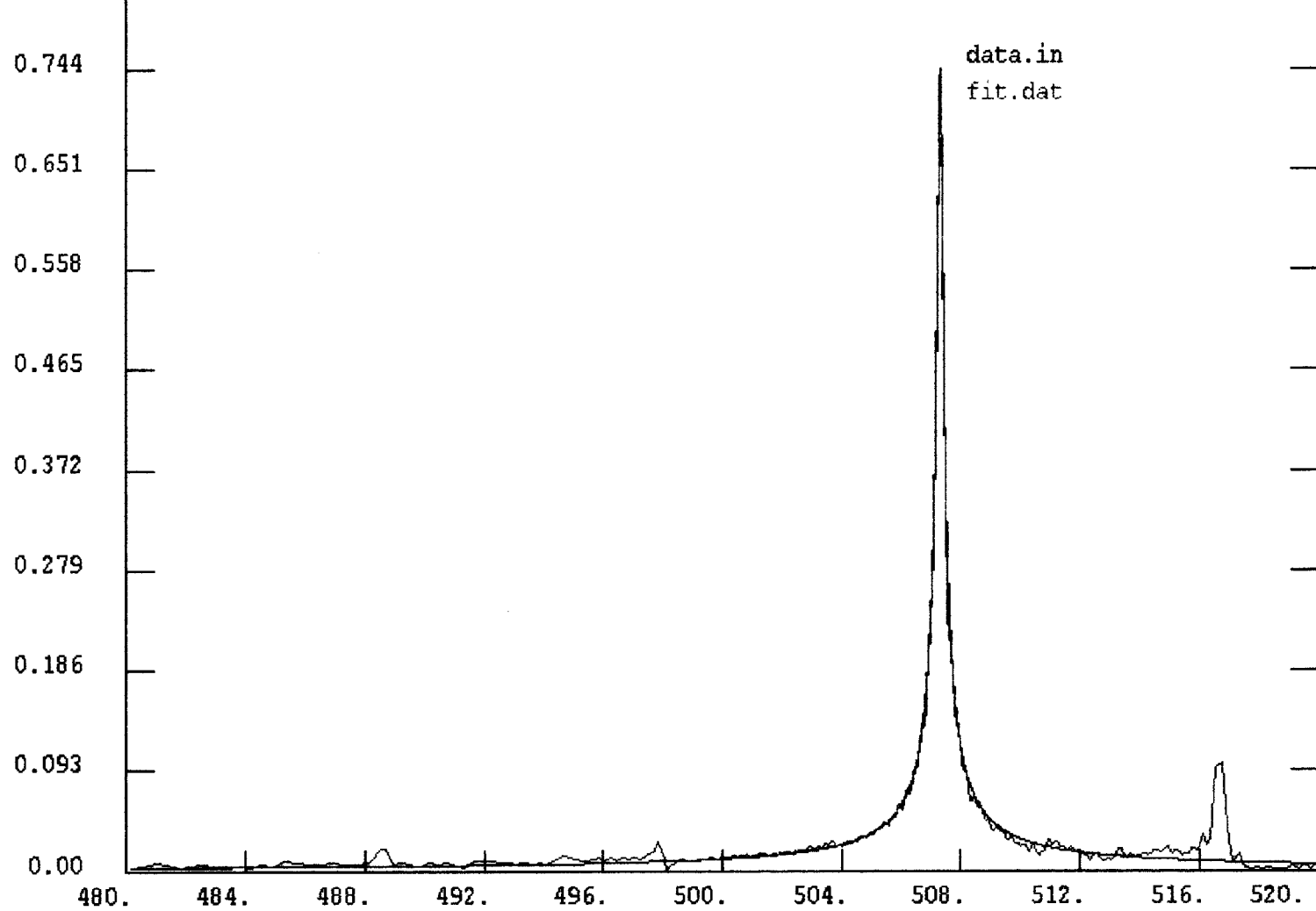
$$D(f) = e^{j2\pi t_0} V(f; f_0', w')$$

$$\Rightarrow P(f; f_0, w) = e^{j2\pi f t_0} \frac{f_0' [(w - jf)^2 + f_0'^2]}{f_0 [(w' - jf)^2 + f_0'^2]}$$

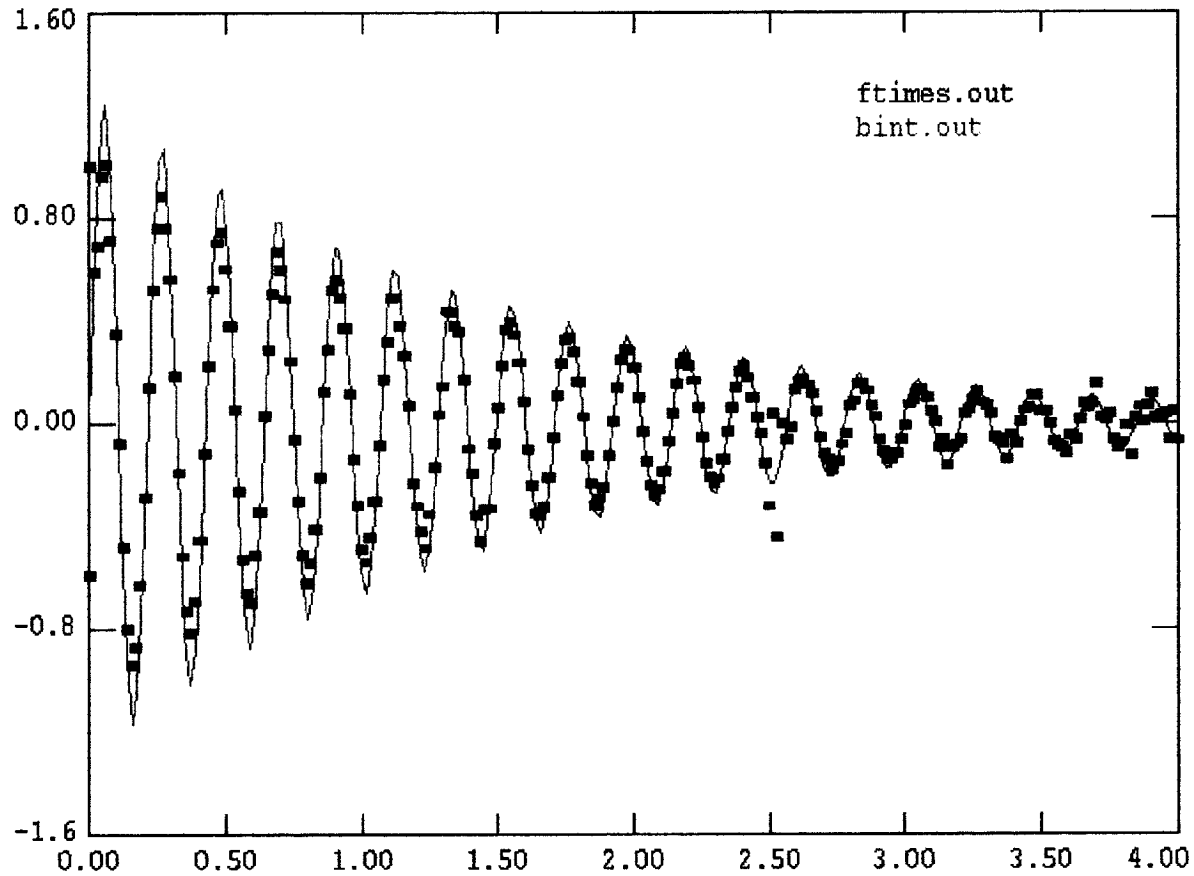
This can be solved analytically to yield

$$\rho(t; t_0, f_0, f'_0, w, w') = \frac{f'_0}{f_0} \delta(t_0 - t) - 2\pi\theta(t - t_0) \frac{e^{-2\pi w(t-t_0)}}{f_0} \left\{ \begin{array}{l} \left[ (w - w')^2 + f_0^2 - f_0'^2 \right] \sin(2\pi f'_0(t - t_0)) \\ - \left[ 2j f'_0 (w - w') \right] \cos(2\pi f'_0(t - t_0)) \end{array} \right\}$$

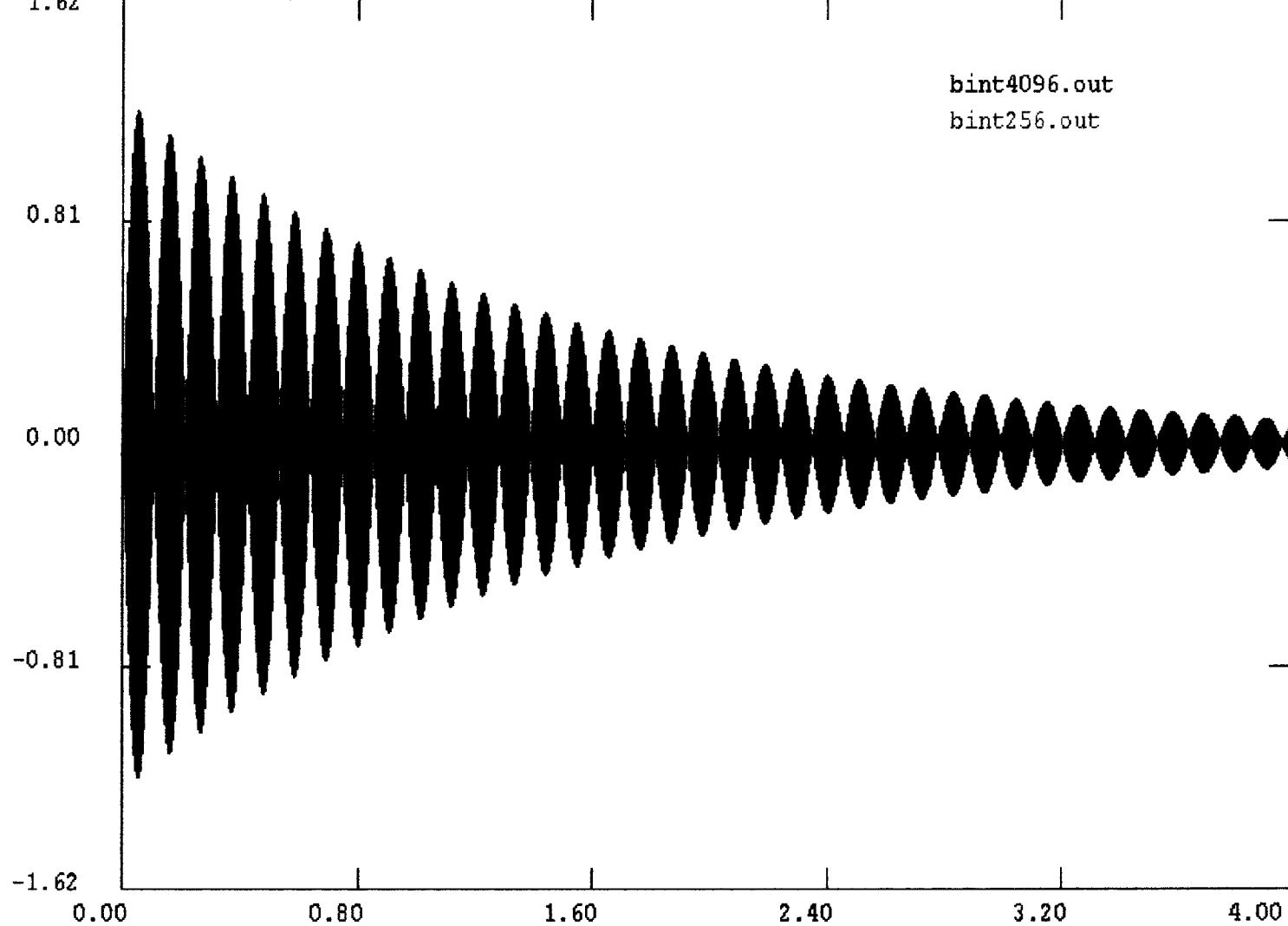
The data can also be fitted as a series of delta functions. The fit below is to a series of 257 delta functions in which the heights and times are both varied.



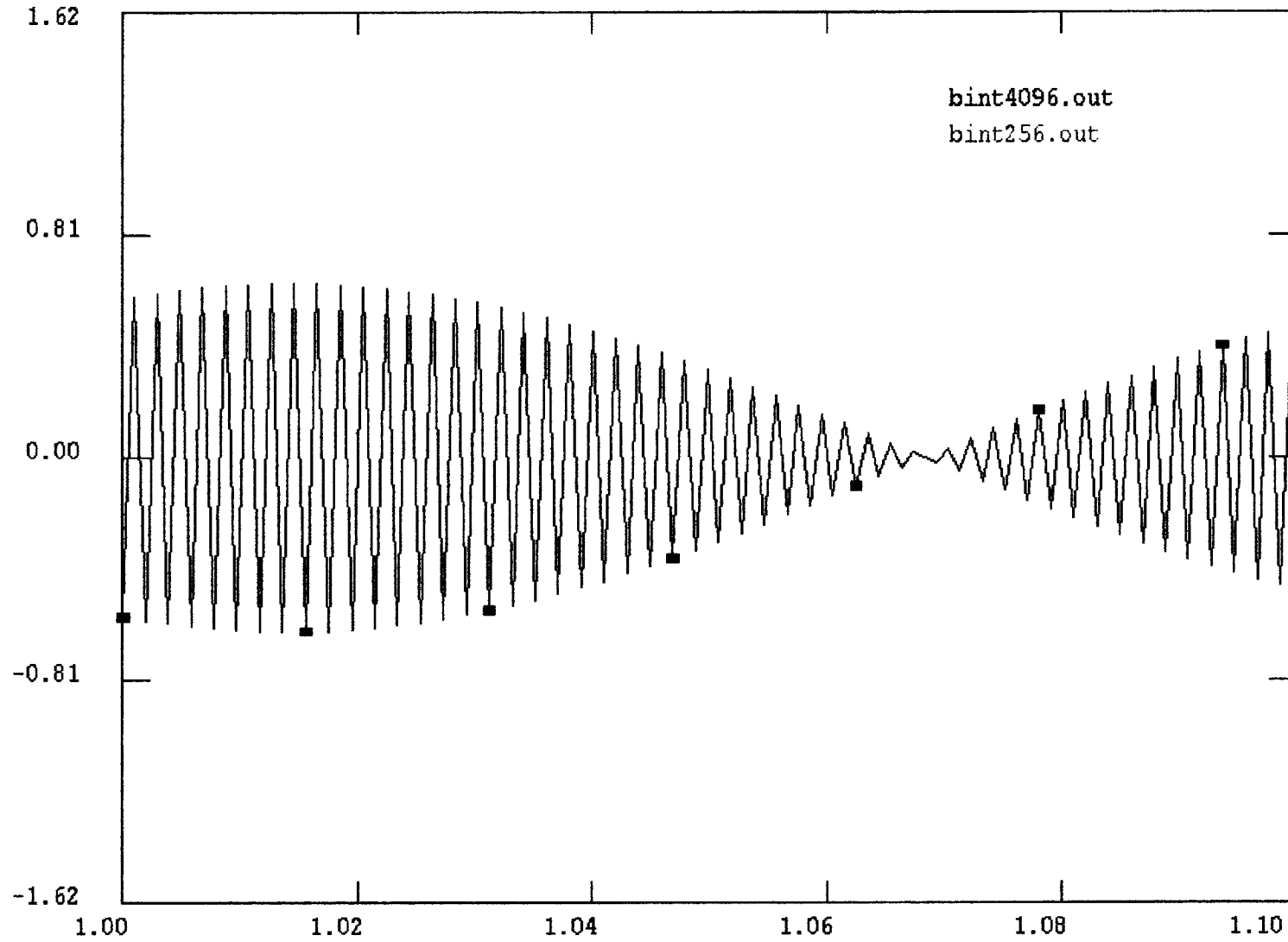
**Figure 5 This is the fit produced by 257 starting times. The violin mode is at 500.87 cycles/second. The data is assumed to consist of a single violin mode at 507 cycles per second.**



**Figure 6 This is a plot  $\rho(t)$  as found by the fit - Black dots and the analytical "result" line connecting points.**

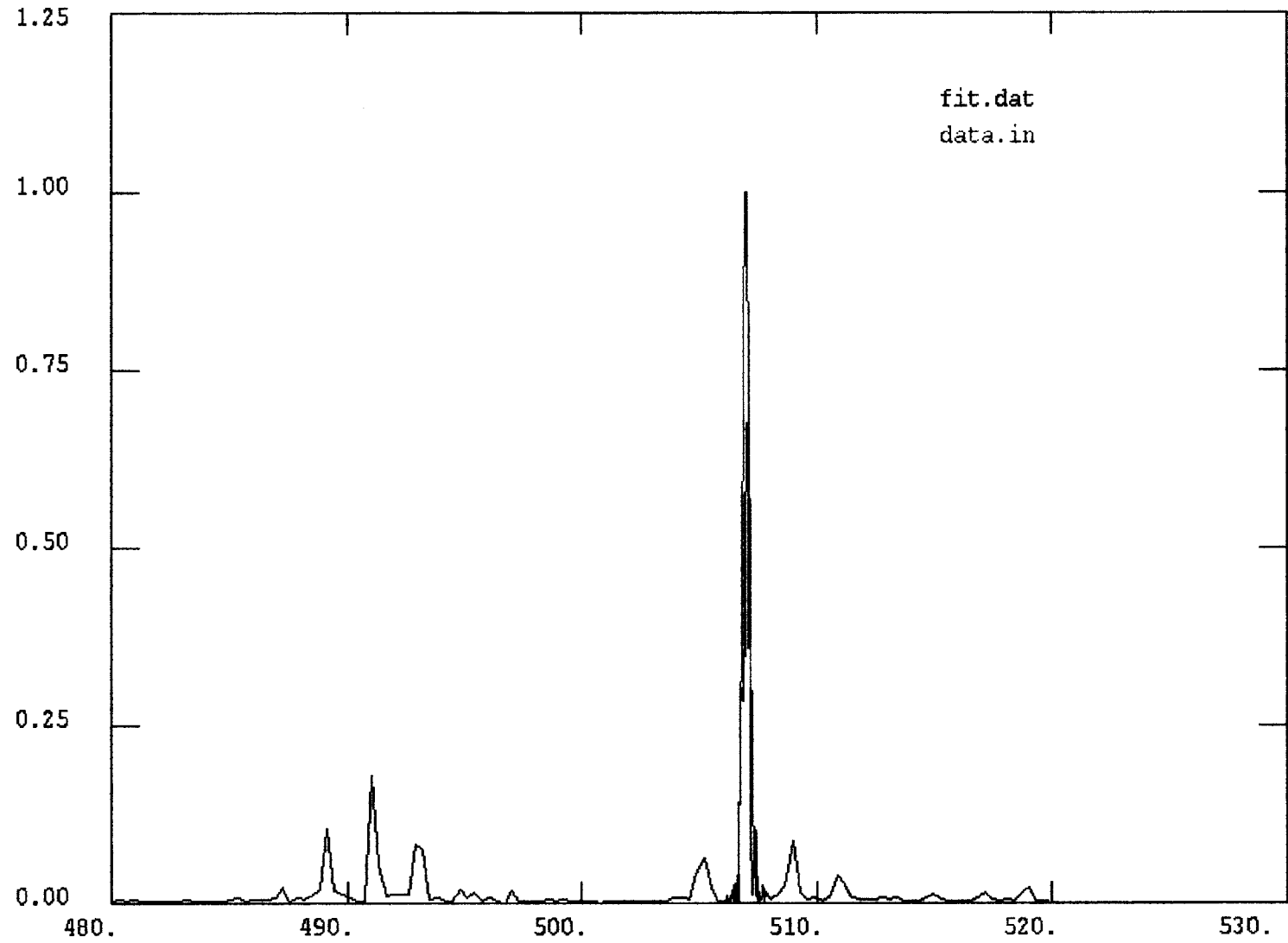


**Figure 7** Now all of the points in the analytical result are shown, rather than just the uniformly spaced 257 points between 0 and 4.

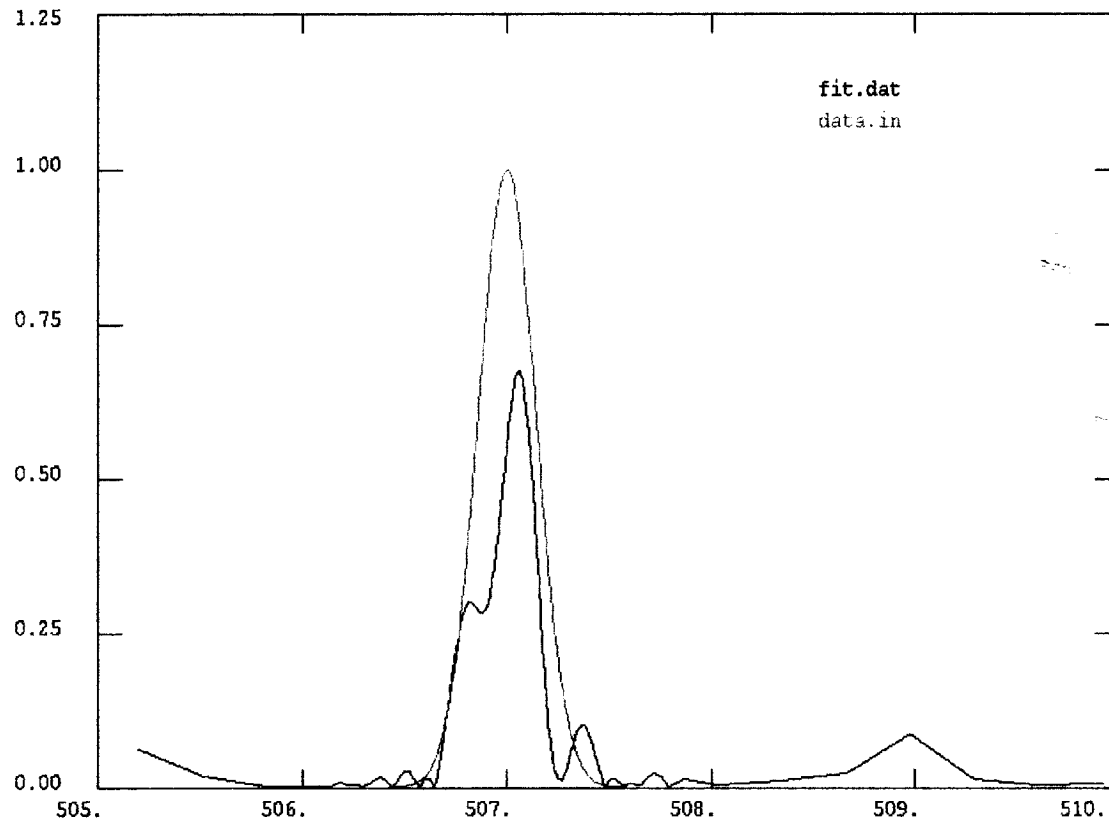


**Figure 8 Greater detail of the fit results versus the analytical solution.**

### Example 2 A Gaussian data peak

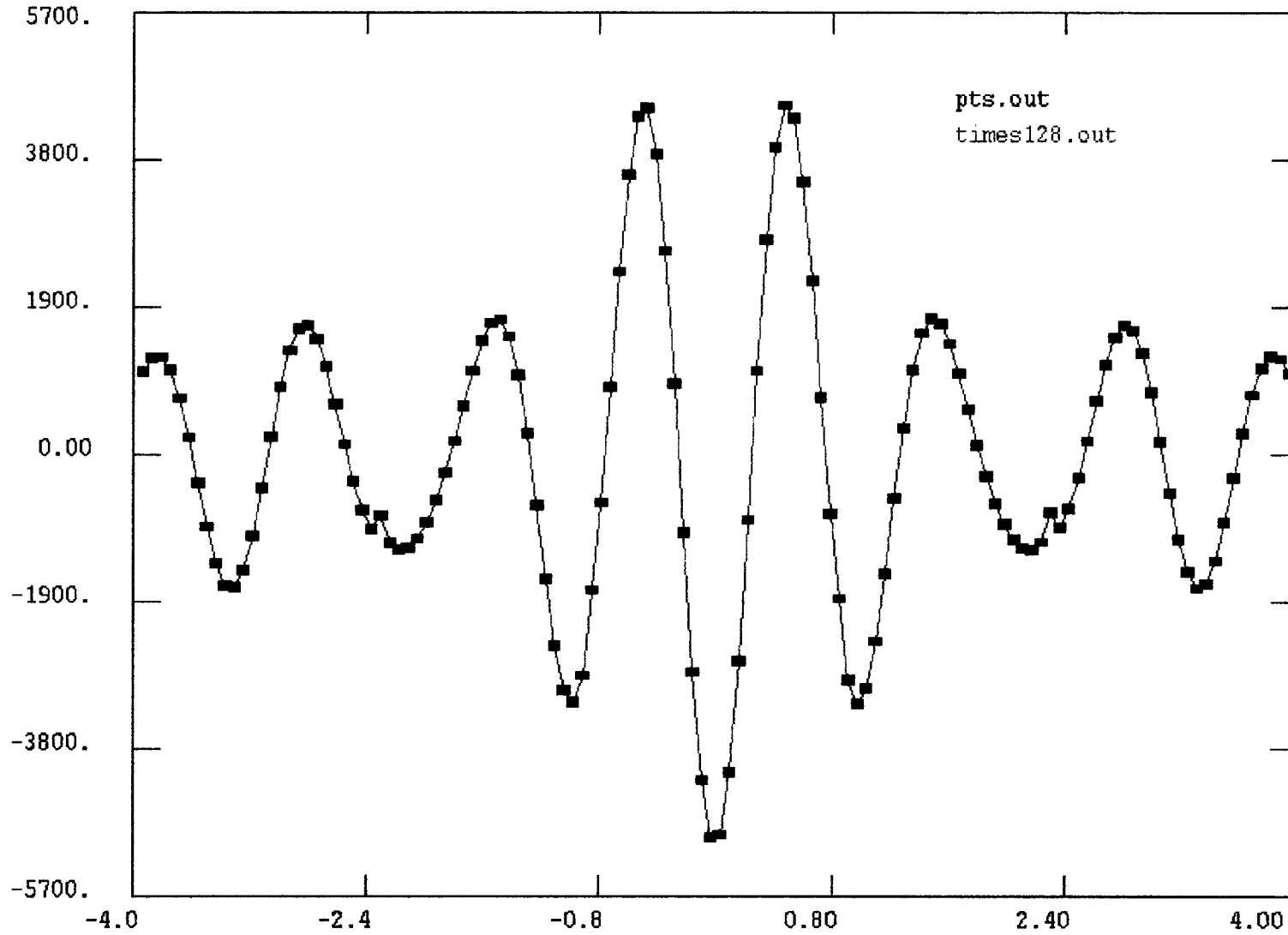






In this case the data is a Gaussian peak centered at 507 cycles/sec near a violin mode at 500.87 cycles/second.

**Figure 9 Expansion of the peak region.**



**Figure 10** Distribution of times leading to above fit.

*Note 1, Linda Turner, 05/09/00 10:21:59 AM*  
LIGO-G000106-00-D