Laser Beam Pointing in LIGO II

and its consequences for the Input Optics

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LIGO-G010078-00-D

- 1. The Problem
- 2. Basic Tools of the Modal Model:
 - Modal Expansion of the light field
 - Tilted Mirrors
- 3. Assembling LIGO II with the Modal Model:
 - Tilted Mirrors in a cavity
 - Cavity Enhanced Michelson
 - LIGO-I
 - LIGO-II
- 4. Requirements for Pointing in LIGO II
 - Equivalent Shot Noise limit
 - Mode Cleaner specs



Distances between tilted mirrors depend on beam position:

 $\delta_1 - \delta_2 = \Delta \left(\Theta_1 - \Theta_2 \right)$

Beam pointing in a misaligned IFO generates noise.

Definitions:

Differential tilt: $\delta_1 = -\delta_2$

Common Tilt: $\delta_1 = \delta_2$

Involved Frequencies:

- (mis-) alignement of mirrors is static
- Pointing of laser happens at GW-frequencies



POINTING IN THE MODAL MODEL

TILTED MIRROR IN THE MODAL MODEL



It can be described with a 2x2 matrix acting on \hat{u}_0 and \hat{u}_1 :

 $E_r = \hat{M}_r * E_{in}$

$$\hat{M}_r = \begin{pmatrix} \sqrt{1 - 4\theta^2} & -2i\theta \\ -2i\theta & \sqrt{1 - 4\theta^2} \end{pmatrix}$$
 with $\theta = \frac{\pi w(z)}{\lambda} \Theta$

$$E_t = \hat{M}_t * E_{in}$$

$$\hat{M}_t = \begin{pmatrix} \sqrt{1-x^2} & x \\ -x & \sqrt{1-x^2} \end{pmatrix}$$
 with $x = \frac{d}{w} \frac{n-1}{n} \Theta$

In interferometric GW-detectors ($\Theta = O(10^{-8})$):

$$rac{\pi w(z)}{\lambda}pprox 10^5 \qquad \gg \qquad rac{d}{w}rac{n-1}{n}pprox 1$$



Roundtrip Propagator:

$$\hat{P}_{RT} = r_1 r_2 \hat{M}_1 \hat{L} \hat{M}_2 \hat{L}$$
 $\hat{L} = \left(egin{array}{c} e^{i\omega rac{L}{c}} & 0 \ 0 & e^{i\omega rac{L}{c} + i\phi_g} \end{array}
ight)$

matrices do NOT commute !

Be careful with the sequence.

$$E_c = it_1 \left(\hat{U} - \hat{P}_{RT} \right)^{-1} E_{in}$$

 $E_t = -t_1 t_2 \hat{L} \left(\hat{U} - \hat{P}_{RT} \right)^{-1} E_{in}$
 $E_r = \left(r_1 \hat{M}_1^{-1} - t_1^2 r_2 \hat{L} \hat{M}_2 \hat{L} \left(\hat{U} - \hat{P}_{RT} \right)^{-1} \right) E_{in}$

We need to know: How much of the 1-mode will show up as a 0-mode at the detector (dark port)?



- ETM tilt rolls off with cavity pol
- peak when 1-mode on resonance (Gouy phase)
- ITM tilt: $E = rE_r + tE_c \rightarrow$ Interference
 - overcoupled cavity: $f < Pol: tE_c > rE_r$
 - **f** > **Pol:** $rE_r > tE_c$
 - f \approx PoI: constructive or destructive Interference



Common Tilt:

Difference between RF-locking and DC-locking:

- 1. RF-locking
 - LO at dark port is RF-sideband
 - carrier *completely* dark
- 2. DC-locking
 - small differential offset in arm cavities ($\pm 10^{-11}m$)
 - carrier is its own LO
 - no RF-sideband necessary



RF-locking:

Cavity enhanced Michelson interferometer with power recycling.

LIGO-I Configuration **T**_{1->0} **Differential tilt 10⁻²** ETM (+f) ETM (-f) 10⁻³ ITM (+f) ITM (-f) 10⁻⁴ 10⁻⁵ 10² 10³ **10**⁴ **10**¹ 1**0**⁰ f [Hz]

Differential Tilt:

LIGO-I with LIGO-II parameter: ITM: T = 0.005PR: $T_{PR} = 0.08$

LIGO-I \approx Cavity enhanced MI x $2\sqrt{T_{PR}}$

PR-filters the pointing





- low frequencies f < 100Hz: SR-mirror filters noise
- above 100Hz: SR-cavity amplifies noise

In General:

Like simple Cavity in reflection filtered by PR-mirror times Signal Recycling Transferfunction



Common Tilt in DC-locking:

low frequencies: gain of PR-cavity boosts pointing signal

Summary:

- Disadvantages for DC-locking at low frequencies
- Still less problematic than differential tilt at all frequencies

Remark:

Additional asymmetries will boost Common Signals in RF-locking scheme.



PR-tilt:

- at low frequencies: factor 1-2 smaller than Common ITM tilt
- strong roll off to higher frequencies

SR-tilt:

- at low frequencies: far below Common ITM
- reaches Common ITM transferfunction at peak frequency

Comments on this (not backed up by any analysis):

- PR-mirror not a probem
- SR-mirror could be, will be hard to detect SR-tilt

Detected Pointing Signal:

$$S_p(f) = E_{LO}^* \left(E_p(+f) e^{i2\pi ft} + E_p(-f) e^{-i2\pi ft} \right) + c.c.$$

= $\Re \left[E_{LO}^* \left(T_{1->0}(+f) e^{i2\pi ft} + T_{1->0}(-f) e^{-i2\pi ft} \right) E_1(f) \right]$

 $E_1(f) = a_1(f)\sqrt{n_{in}}, \qquad a_1(f) = x(f) + i\alpha(f)$

Don't know the phase (?). Assume worst case:

 $S_p(f) \le |E_{LO}| \left(|T_{1->0}(+f)| + |T_{1->0}(-f)| \right) |a_1(f)| |E_{in}|$

equivalent Shot Noise Limit (P=125W, $n_{in} = 6.3 \cdot 10^{20}/s$)

$$|a_1(f)| \le \frac{1}{(|T_{1->0}(f)| + |T_{1->0}(-f)|)\sqrt{n_{in}}}$$



The limit for differential ITM (the most crucial) tilt:

$$\begin{aligned} |a_1(30Hz)| &< \frac{1.0 \cdot 10^{-7}}{\sqrt{Hz}} \frac{[10^{-8}]}{\Theta_D^{ITM}} \sqrt{\frac{[125W]}{P_{in}}} \\ |a_1(100Hz)| &< \frac{3.4 \cdot 10^{-8}}{\sqrt{Hz}} \frac{[10^{-8}]}{\Theta_D^{ITM}} \sqrt{\frac{[125W]}{P_{in}}} \\ |a_1(300Hz)| &< \frac{5.7 \cdot 10^{-9}}{\sqrt{Hz}} \frac{[10^{-8}]}{\Theta_D^{ITM}} \sqrt{\frac{[125W]}{P_{in}}} \end{aligned}$$

Both dimensions and PSL specs:

f	$a_{10/01}^{max}(f)[Hz^{-1/2}]$	$a_{10/01}^{PSL}(f)[Hz^{-1/2}]$	a_1^{max}/a_1^{PSL}
30Hz	$5.0 \cdot 10^{-8}$	$1.3 \cdot 10^{-5}$	$3.8 \cdot 10^{-3}$
100Hz	$1.7 \cdot 10^{-8}$	$4 \cdot 10^{-6}$	$4.3 \cdot 10^{-3}$
300Hz	$2.8 \cdot 10^{-9}$	$1.3 \cdot 10^{-6}$	$2.2 \cdot 10^{-3}$

Mode cleaner as a passive filter: Finesse: 2026, g-factor 0.407 Amplitude Transmission $TEM_{10} < 10^{-3}$

- Leaves a safety factor of 2 (compared to 10 required for LIGO I).
- Requires that pointing in the GW-band is dominated by PSL.

Safety factor of 2 is not acceptable. Three solutions:

- active steering system
- lower offsets in mirror tilts
- additional (low finesse) mode cleaner

Parameter:

Arm Cavities: 4 kmDe $T_{ITM} = 0.005$ T_E Rayleigh Range: 2000 mGinline short MI: 3.21 morPR-arm: 5.34 mSI $T_{PR} = 0.08$ T_S

DC-detuning: $\pm 10^{-11}m$ *T_{ETM}*=100ppm **Gouy phase:** 0.5 rad **outline short MI:** 2.79 m **SR-arm:** 6.134 m *T_{SR}* = 0.06

SR-detuning: 87.4deg from SR

beamsize: 4 cm only scales the mixing angles of the tilted mirrors.

Literature:

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