

Superconducting cavities for the detection of high frequency gw

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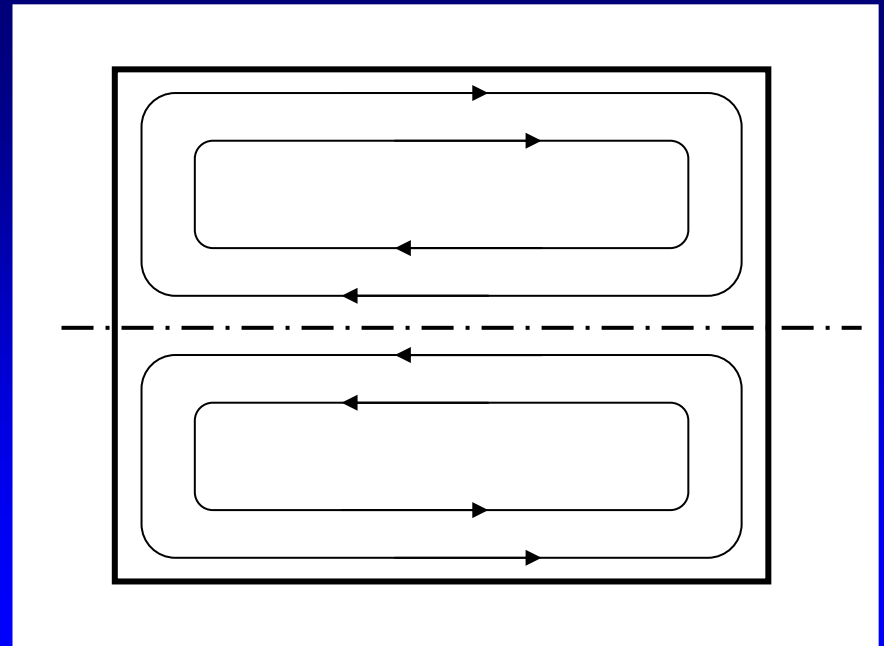
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E. Picasso

SNS, Pisa

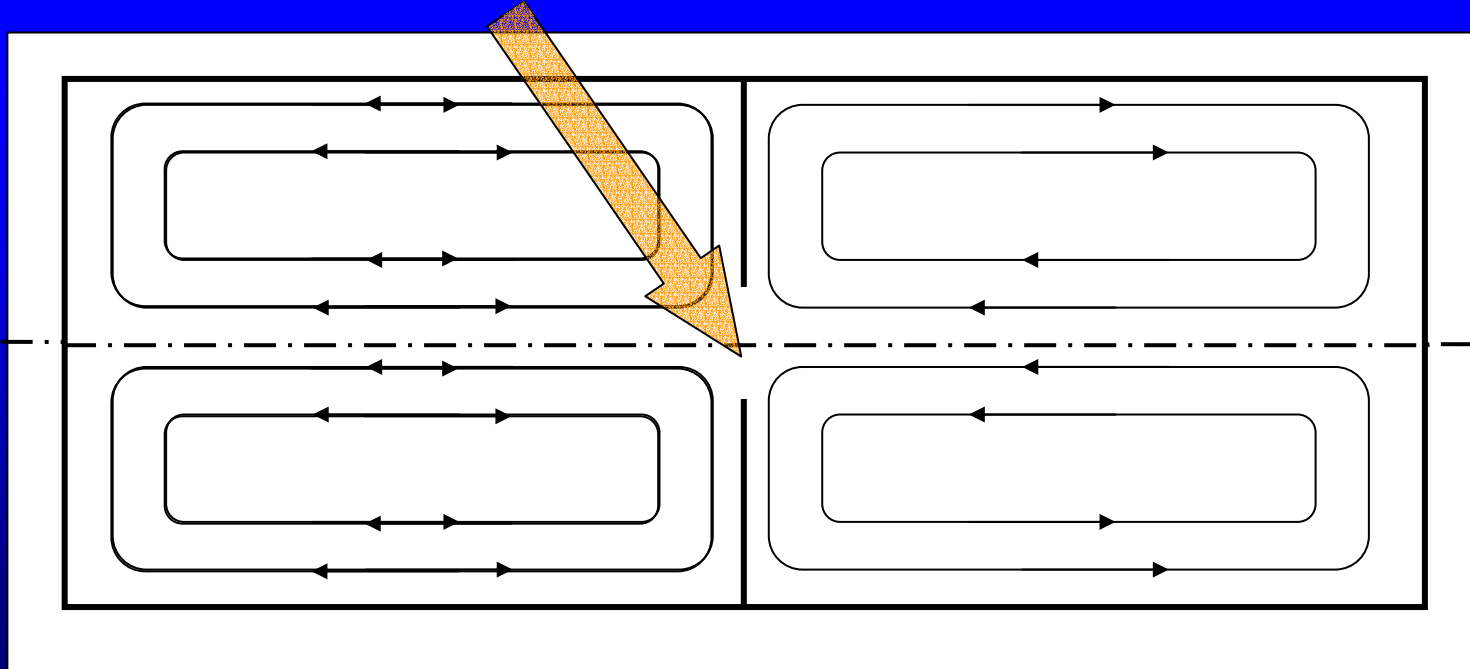
**Pill-box cavity
TE₀₁₁ mode**



Symmetric mode: ω_s

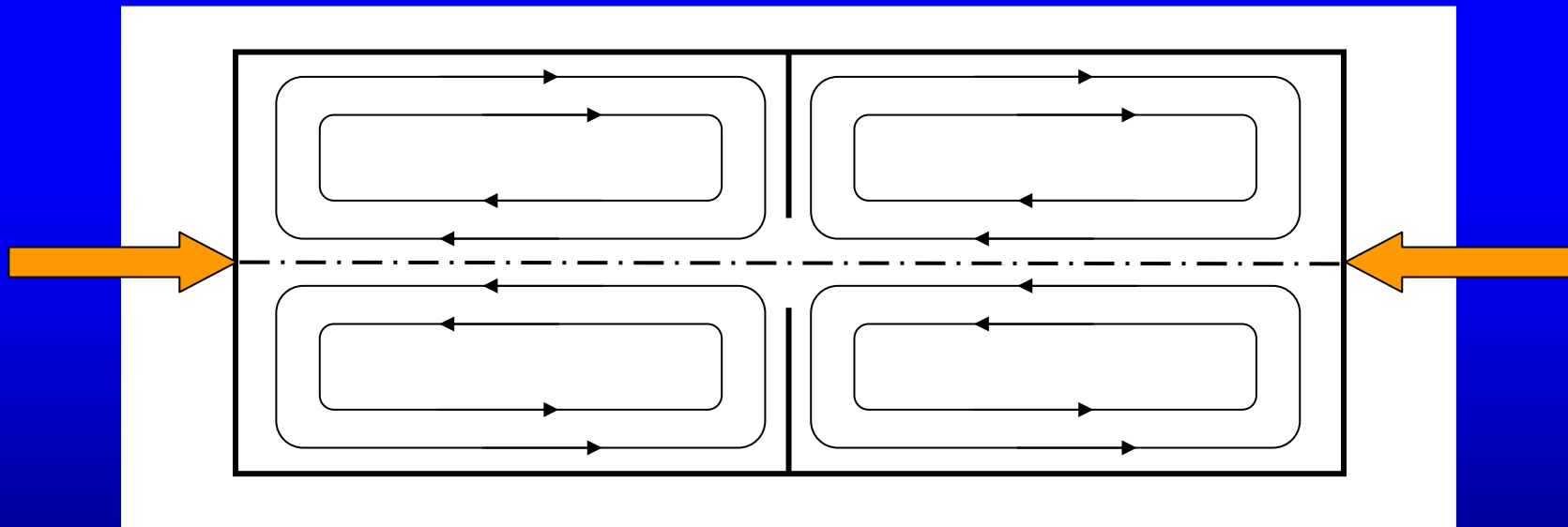
Antisymmetric mode: ω_a

**$\omega_a - \omega_s$ proportional to the
coupling strength (tunable)**



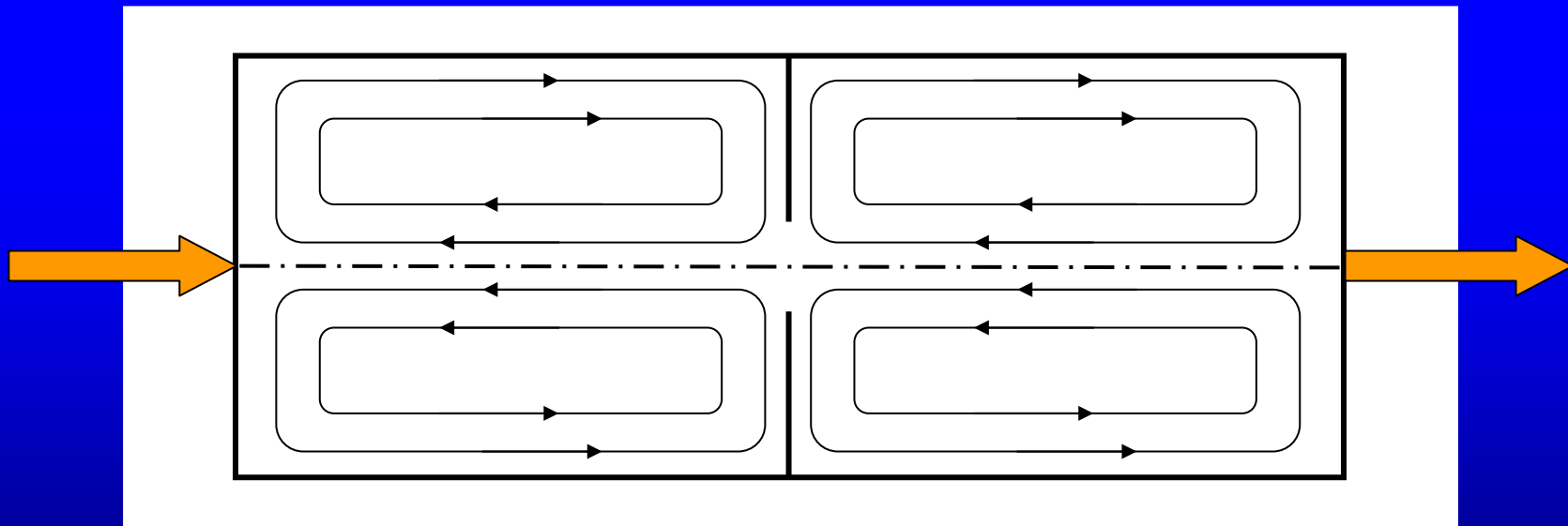
If the symmetric mode is initially excited and we perturb one system parameter (e.g. the length of the cavity) with a characteristic frequency much lower than the normal mode frequency ($\Omega \ll \omega_0$)...

...we can just have a modulation of the mode frequency, ***without any mode mixing...***



...or we can have a **coupling** between the two normal modes of the unperturbed system → there is transfer of energy from one mode to the other;

the energy transfer is maximum when the frequency of the external perturbation equals the normal modes frequency difference: $\Omega = \omega_a - \omega_s$



Let us note that this effect depends on **how we perturb** the system

Field equations in an e.m. resonator with perturbed boundaries

$$\vec{u}(\vec{r}, t) = \sum_{\alpha} q_{\alpha}(t) \vec{\xi}_{\alpha}(\vec{r})$$

$$\vec{E}(\vec{r}, t) = \sum_n \mathcal{E}_n(t) \vec{E}_n(\vec{r})$$

$$f_{\alpha}(t) = \int_{Vol} \vec{f}(\vec{r}, t) \cdot \vec{\xi}_{\alpha}(\vec{r}) dV$$

$$\vec{H}(\vec{r}, t) = \sum_n \mathcal{H}_n(t) \vec{H}_n(\vec{r})$$

$$H = \frac{1}{2} \left[\mathcal{E}_1^2 + \mathcal{H}_1^2 + \mathcal{E}_2^2 + \mathcal{H}_2^2 + \frac{p_m^2}{M} + M\omega_m^2 q_m^2 \right] + \text{Free fields}$$

$$- \frac{q_m}{2} \left(C_{11}^m \mathcal{H}_1^2 + C_{22}^m \mathcal{H}_2^2 + 2C_{12}^m \mathcal{H}_1 \mathcal{H}_2 \right) + \text{Interaction}$$

$$- q_m f_m \text{ External force}$$

Coupling coefficient

$$C_{nm}^{\alpha} = \int_S \left(\vec{H}_n \cdot \vec{H}_m - \vec{E}_n \cdot \vec{E}_m \right) \vec{\xi}_{\alpha} \cdot d\vec{S}$$

PARAmetric COnverter (1998-2000)



Two **pill-box** niobium cavities mounted end-to-end and coupled through a small aperture on the axis

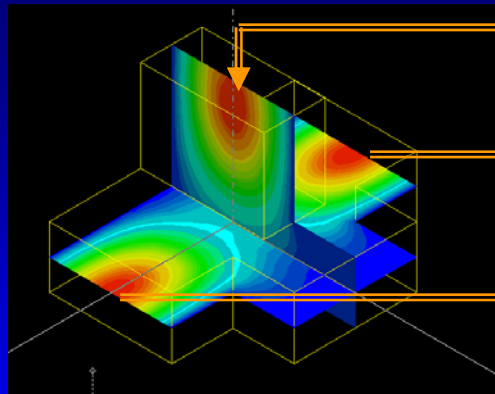
Wall movement induced by a piezoelectric crystal

Working frequency ≈ 3 GHz

Mode splitting ≈ 500 kHz

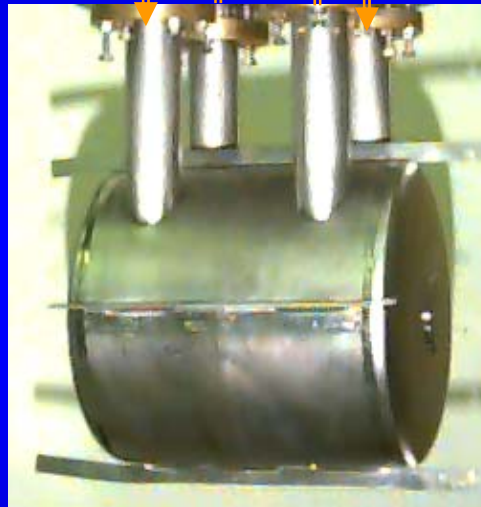
Quality factor (e.m.) 2×10^9 @ 1.8 K

Stored energy **1.8 J**



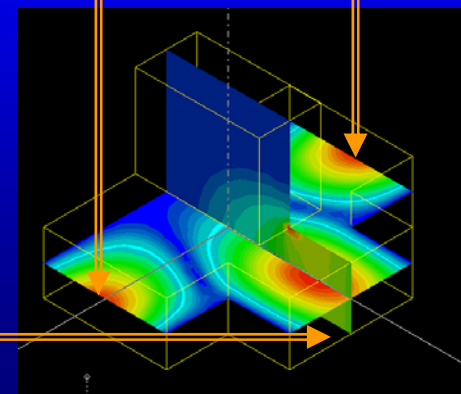
Input signal (ω_s)

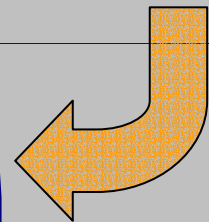
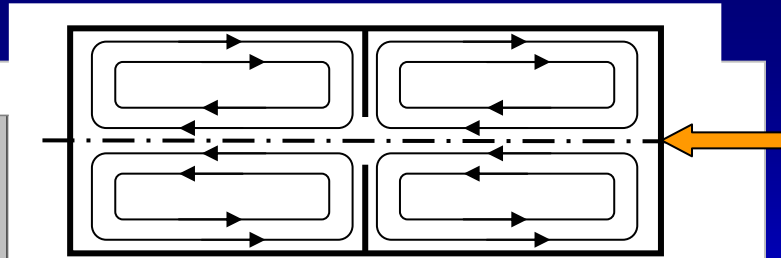
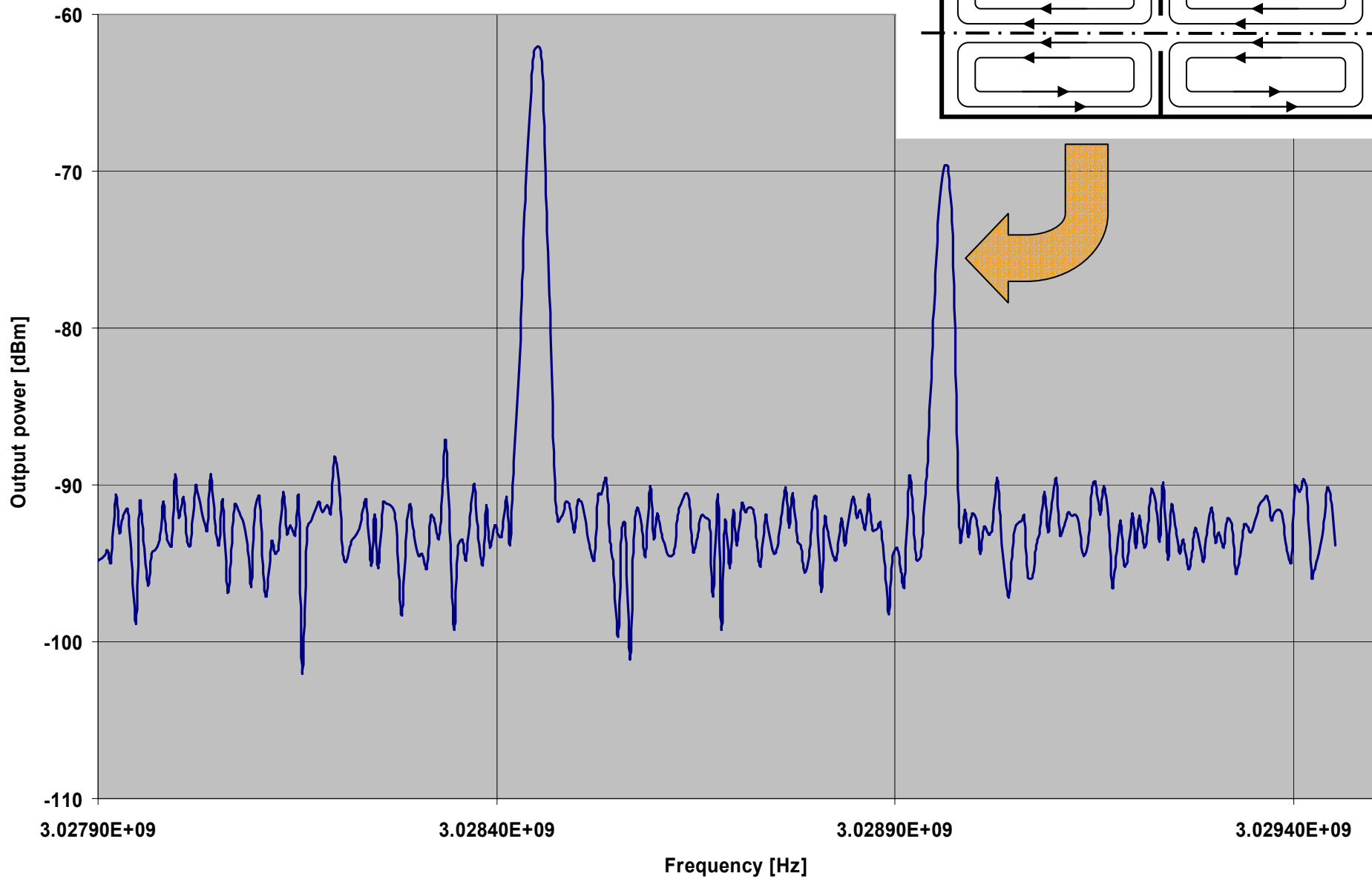
Symmetric
(common) mode



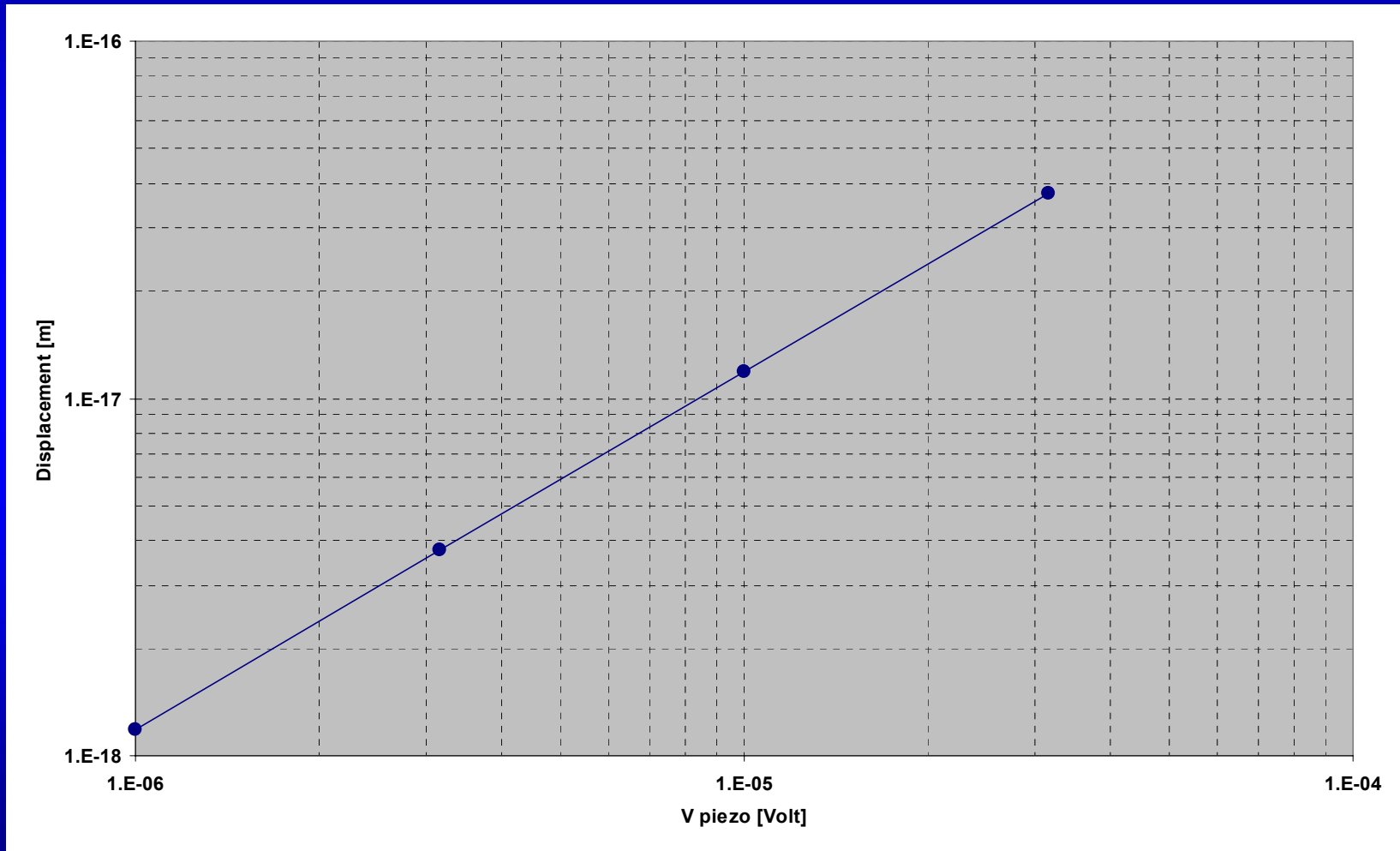
Antisymmetric (differential) mode

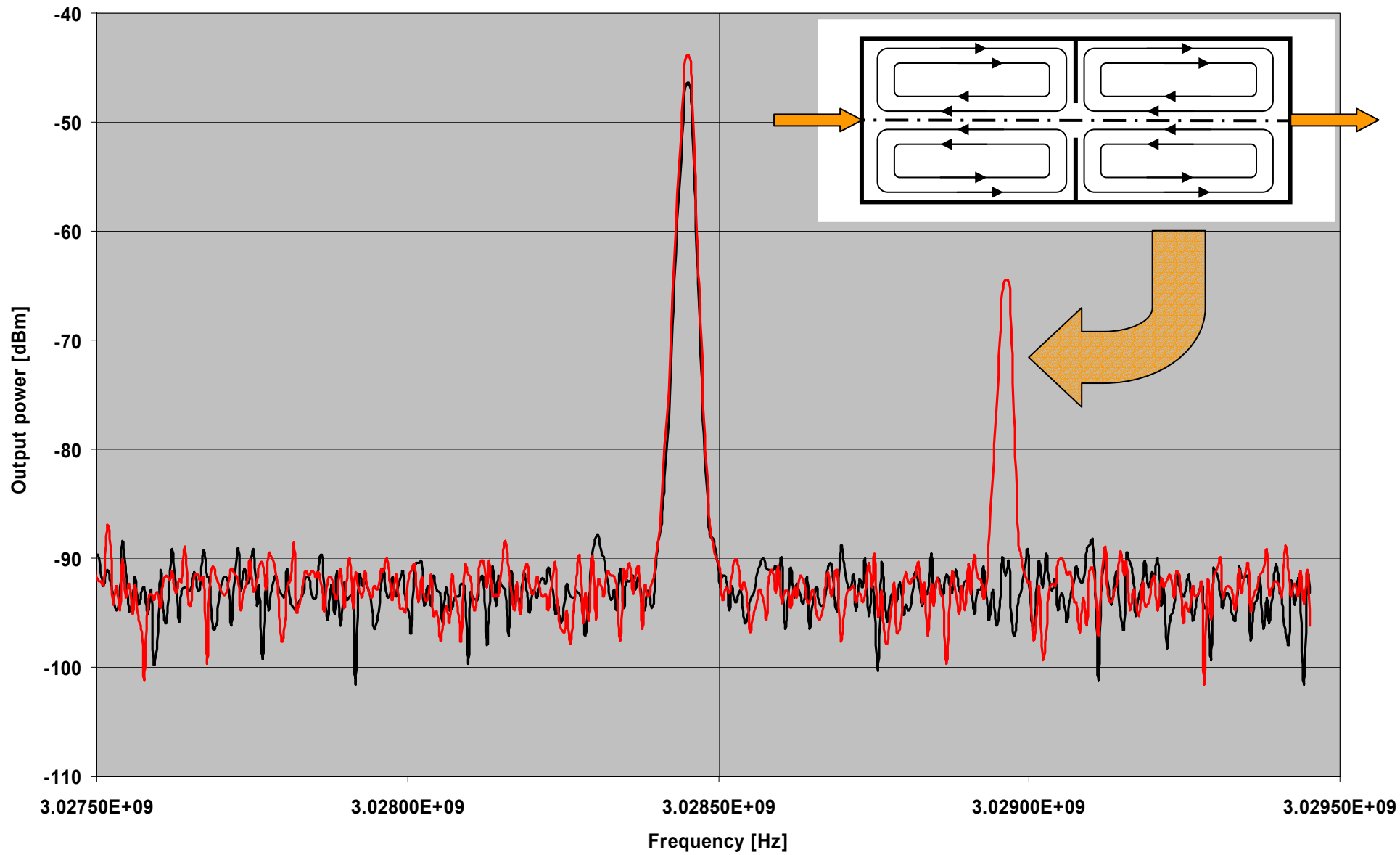
Output signal (ω_a)






Response to piezo excitation



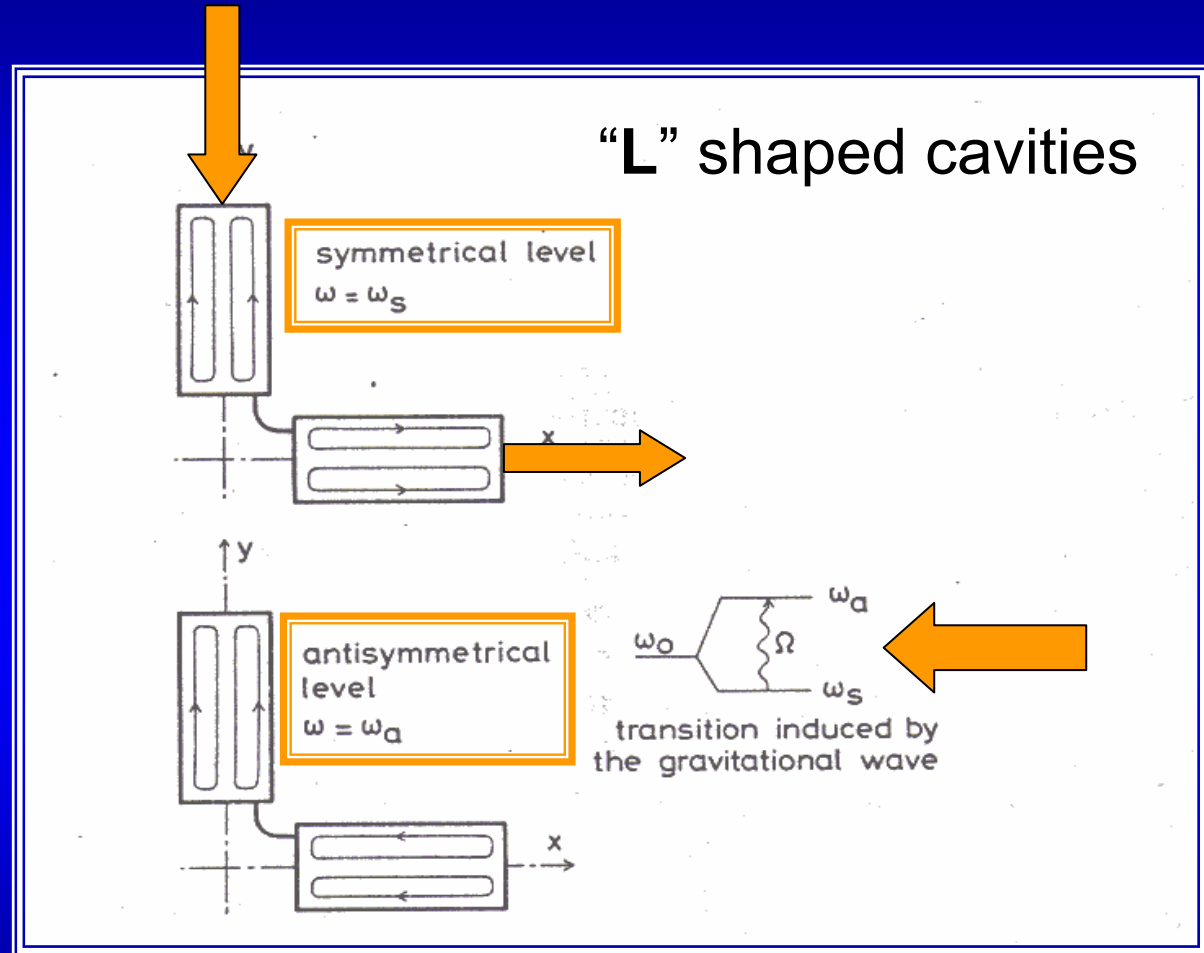


PACO - 2 (2001-2003)

- Lower detection frequency (10 kHz)
- Variable coupling  tuning system
- Spherical cavities development (in collaboration with CERN)

When we take into account the quadrupolar character of the gw...

...we realize that the cavity shape has to be chosen in order to maximize the energy transfer between the two resonant modes



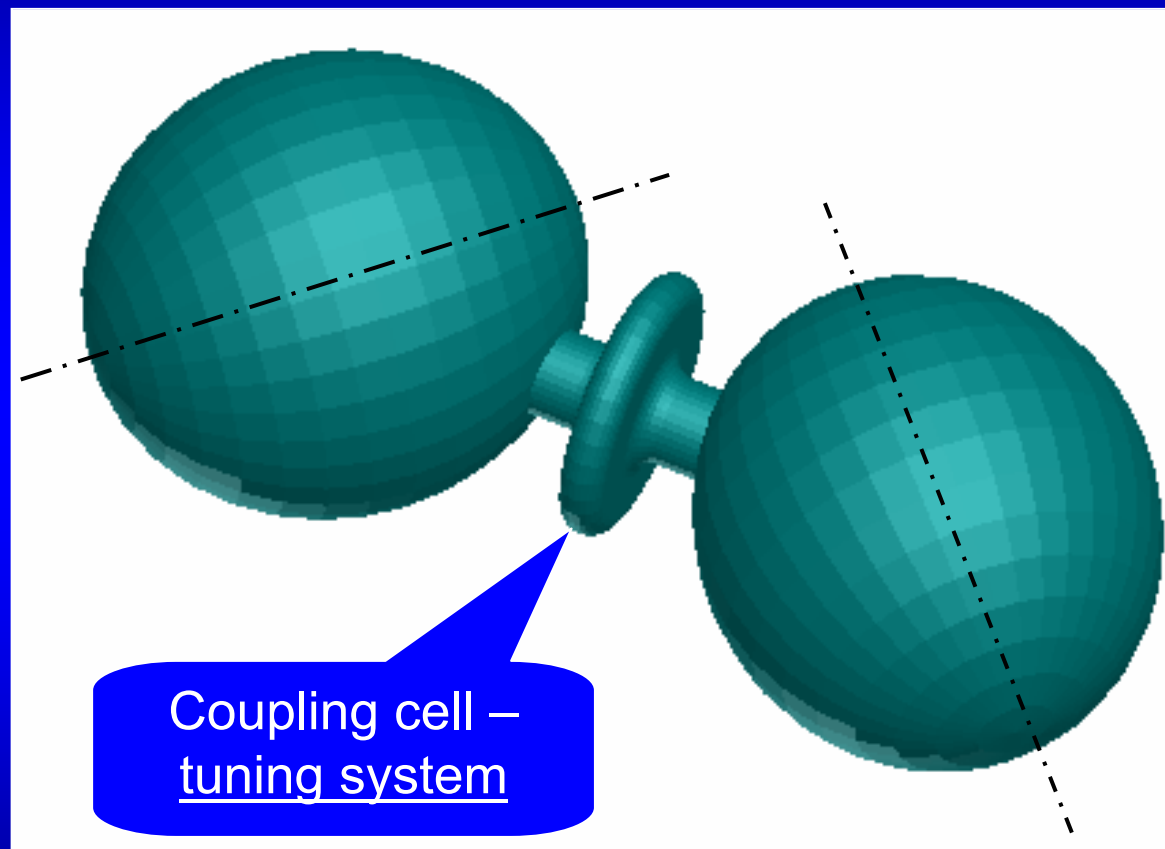
PACO-2 conceptual layout

Cavity internal radius: 100 mm

Operating rf frequency
(TE₀₁₁ mode) \approx 2 GHz

Mode splitting \approx 10 kHz

Stored energy \approx 10 J



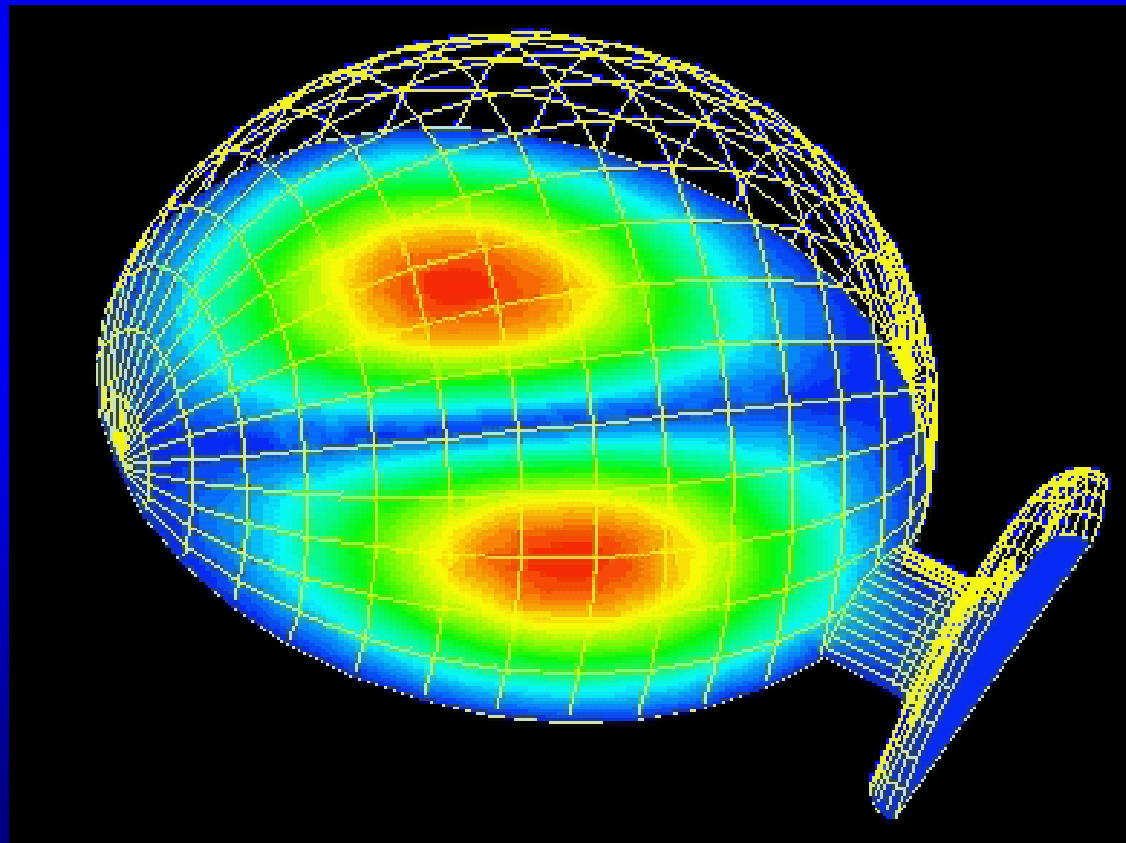
Why spherical cavities?

- Highest e.m. geometrical factor → **highest e.m. quality factor** for a given surface resistance ($Q = G/R_s$)
 - For the TE_{011} mode of a **sphere** $G \sim 850 \Omega$,
 - For the TM_{010} mode of a standard **elliptical** accelerating cavity, $G \sim 250 \Omega$
- Typical values of quality factor of accelerating cavities (TM modes) are in the range **$10^{10} - 10^{11}$**
- The quality factor of the TE_{011} mode of a spherical cavity may well exceed 10^{11}

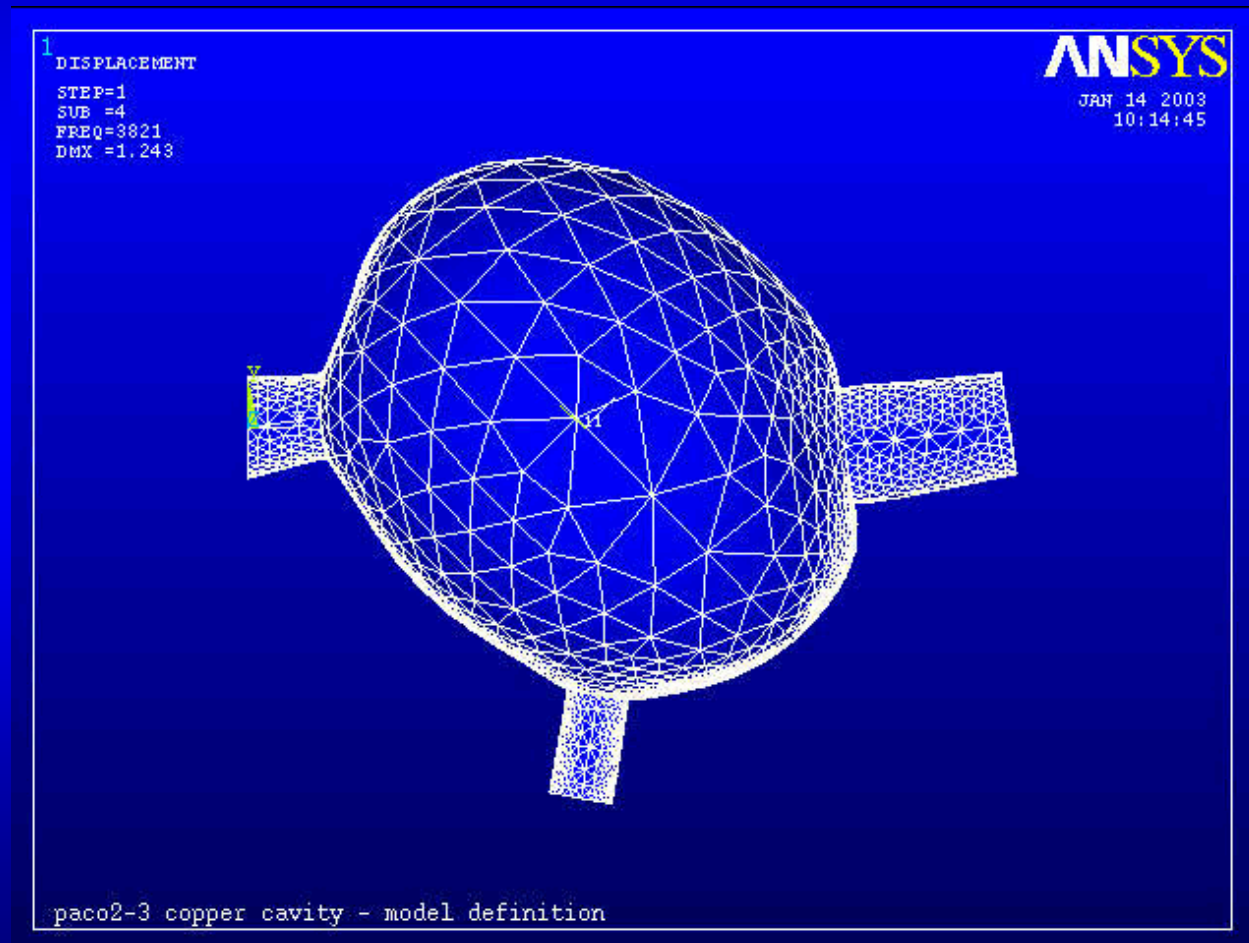
- The spherical cell can be easily deformed in order to remove the e.m. modes degeneracy and to induce the field polarization suitable for g.w. detection
- The interaction between the stored e.m. field and the time-varying boundary conditions depends both on **how the boundary is deformed** and on **the spatial distribution of the fields** inside the resonator
- The optimal field spatial distribution is with the field axis in the two cavities **orthogonal** to each other

TE₀₁₁ mode @ 2 GHz

Electric field magnitude



Quadrupolar mode @ 4 kHz



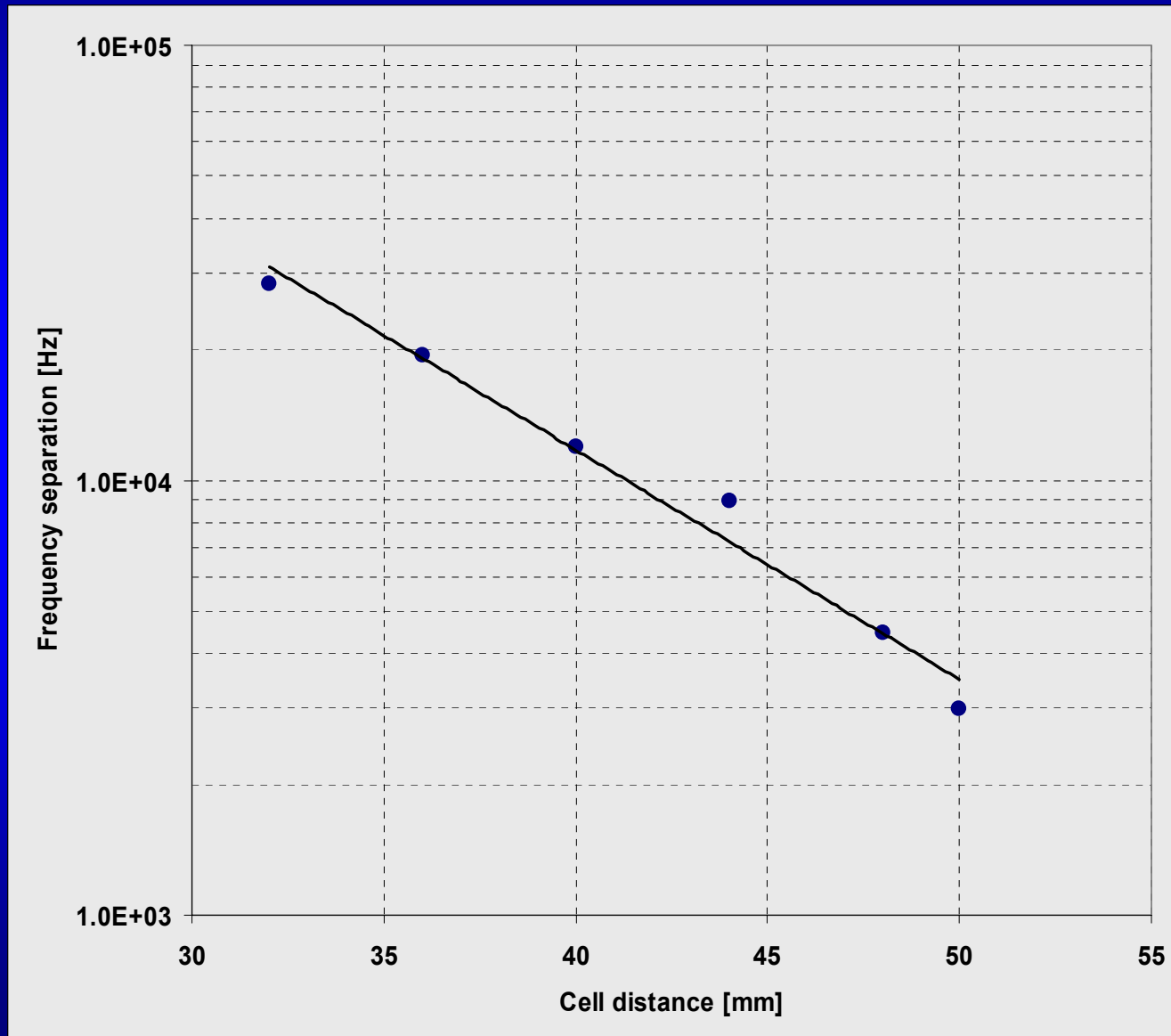


Tuning system

Tuning cell



Mode splitting vs. coupling cell length

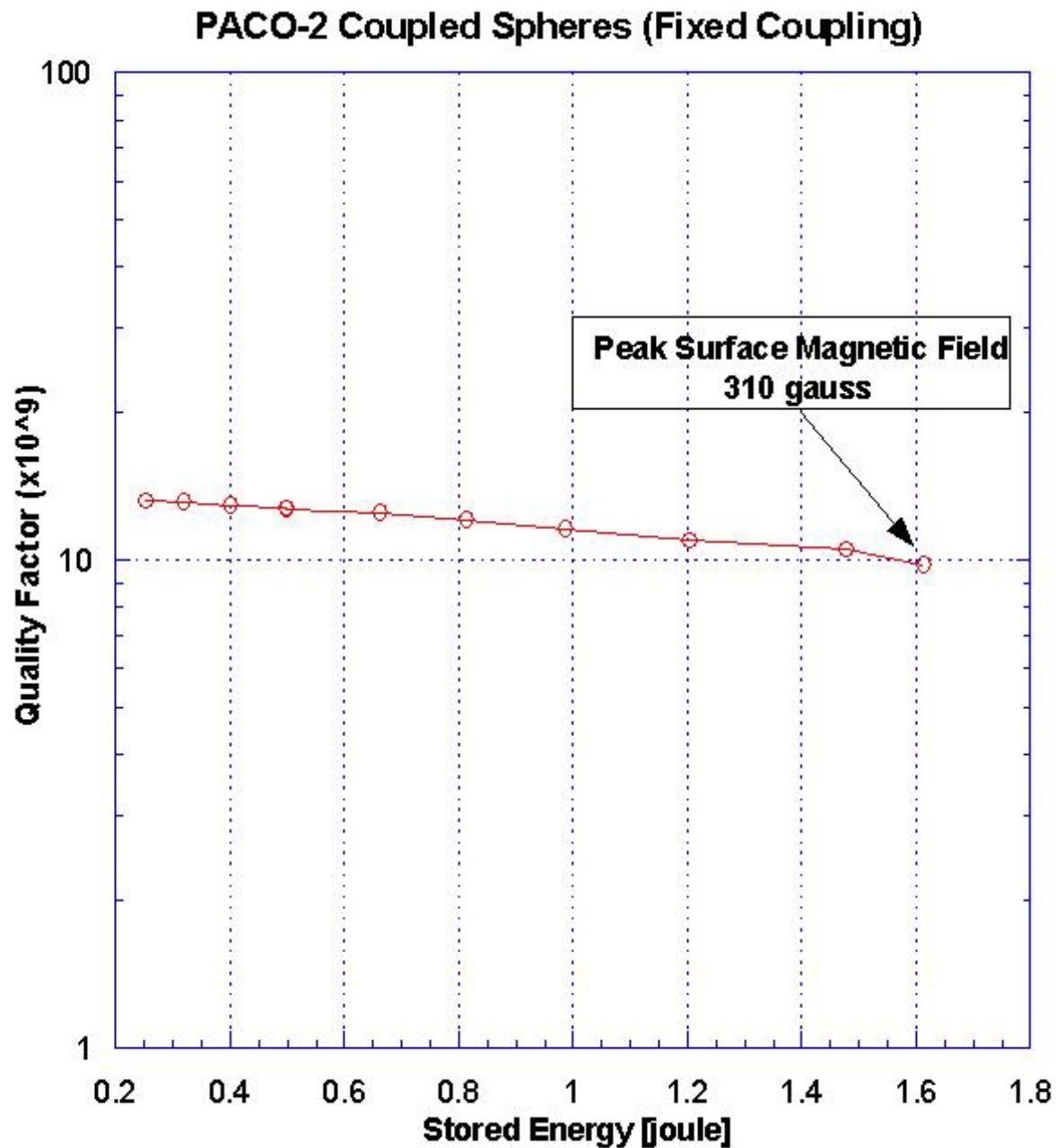




Fixed coupling

Niobium cavity built and tested at CERN

Electromagnetic
test of the niobium
cavity



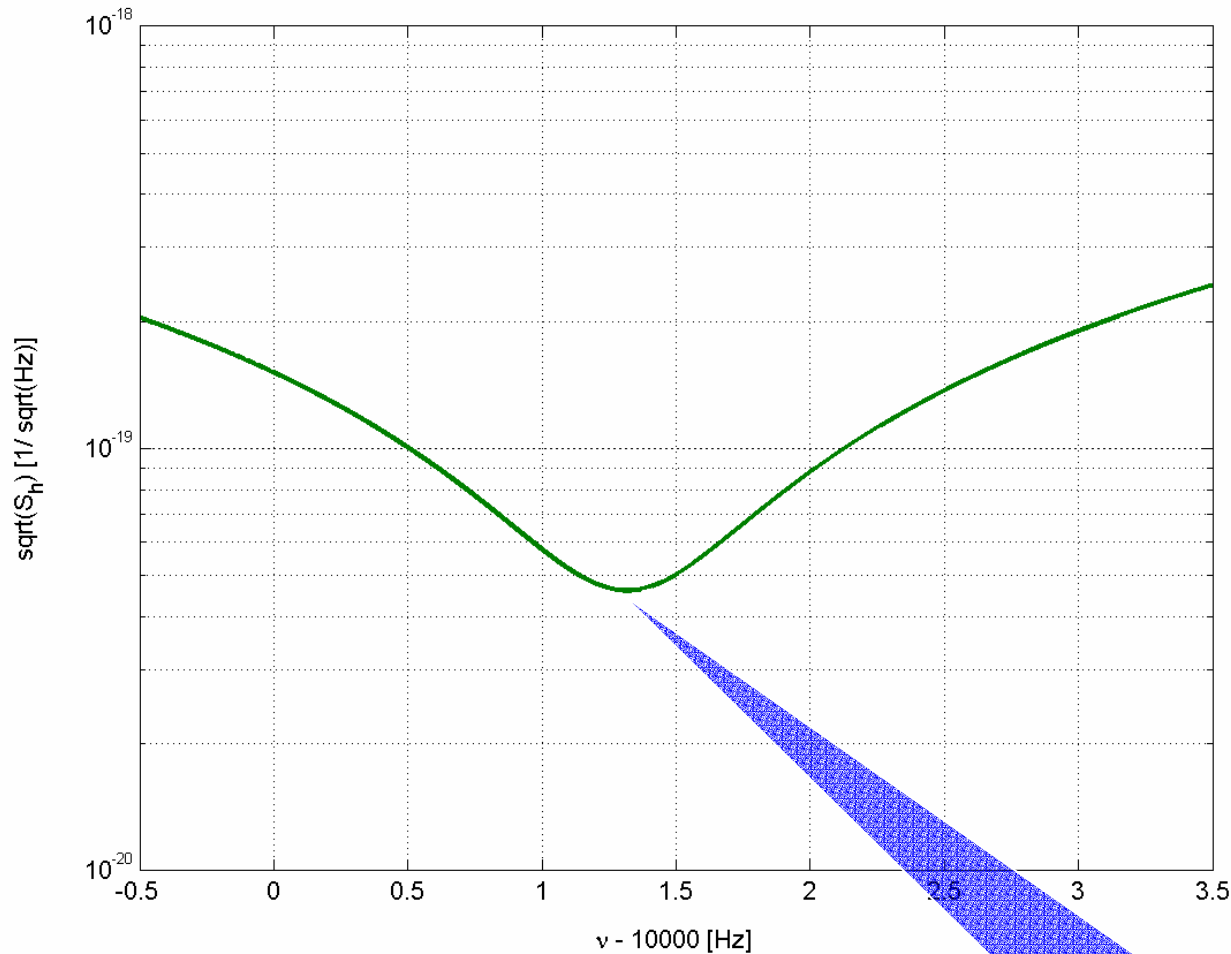
Tunable cavity at CERN

(Jan, 31th 2003)

Tuning cell



Expected sensitivity



Detection frequency
(mode splitting) = **10 kHz**

Mechanical resonant
frequency = 4 kHz

$$U_1 = 10 \text{ J}$$

$$Q = 10^{10}$$

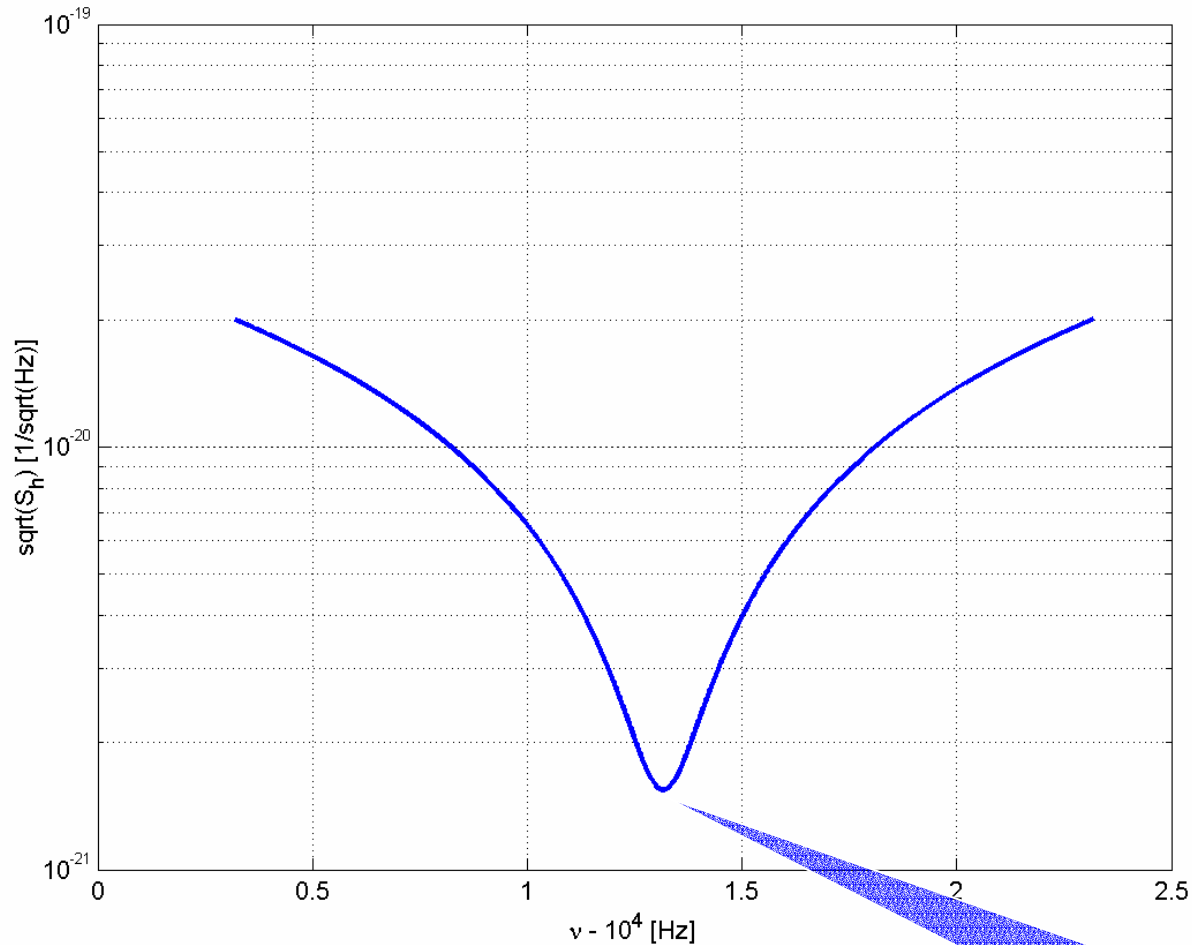
$$Q_m = 10^3$$

$$T = 1.8 \text{ K}$$

$$T_n = 30 \text{ K}$$

$$(\mathbf{S_h})^{1/2} \approx 4 \times 10^{-20} \text{ (Hz)}^{-1/2}$$

Expected sensitivity



Detection frequency
(mode splitting) = **10 kHz**

Mechanical resonant
frequency = 4 kHz

$$U_1 = 10 \text{ J}$$

$$Q = 10^{10}$$

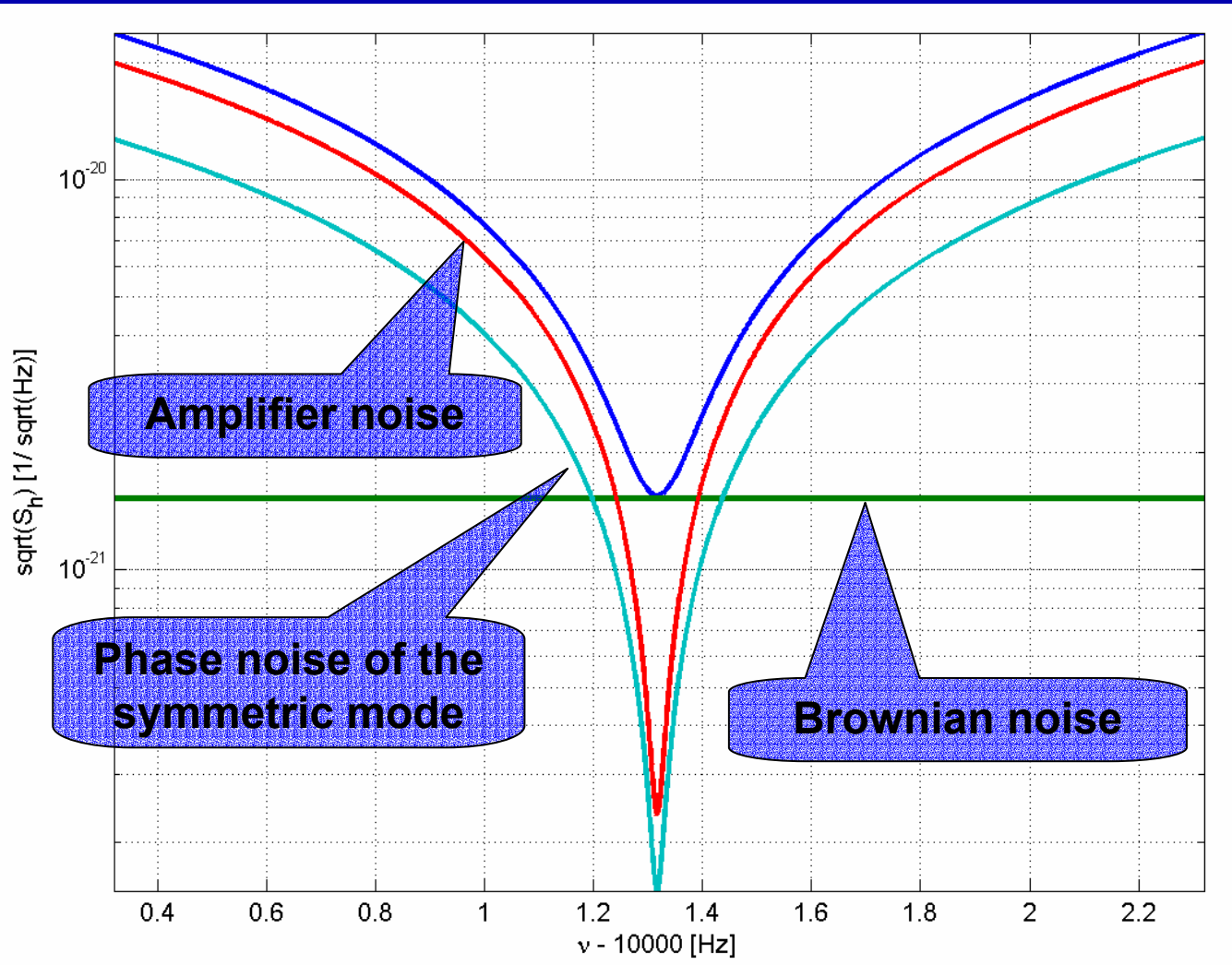
$$Q_m = 10^6$$

$$T = 1.8 \text{ K}$$

$$T_n = 1 \text{ K}$$

$$(\mathbf{S_h})^{1/2} \approx \mathbf{2 \times 10^{-21} \text{ (Hz)^{-1/2}}$$

Expected sensitivity



Detection frequency
(mode splitting) = **10 kHz**

Mechanical resonant
frequency = 4 kHz

$$U_1 = 10 \text{ J}$$

$$Q = 10^{10}$$

$$Q_m = 10^6$$

$$T = 1.8 \text{ K}$$

$$T_n = 1 \text{ K}$$

Conclusions

- Design and realization of an experiment based on the two existing cavities:
 - $\omega \approx 2 \text{ GHz}$
 - **detection frequency $\approx 10 \text{ kHz}$ (tunable between 7 - 20 kHz)**
 - $(S_h)^{1/2} \approx 10^{-21} - 10^{-20}$
- Design of the cryogenic system;
- Design of the suspension system;
- Low noise electronics;
- Data analysis
- Timescale: four years (2004-2007)