
Study of a quantum nondemolition interferometer using ponderomotive squeezing

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1. Introduction

■ Sensitivity of GW detectors will be limited by quantum noise

» Shot noise

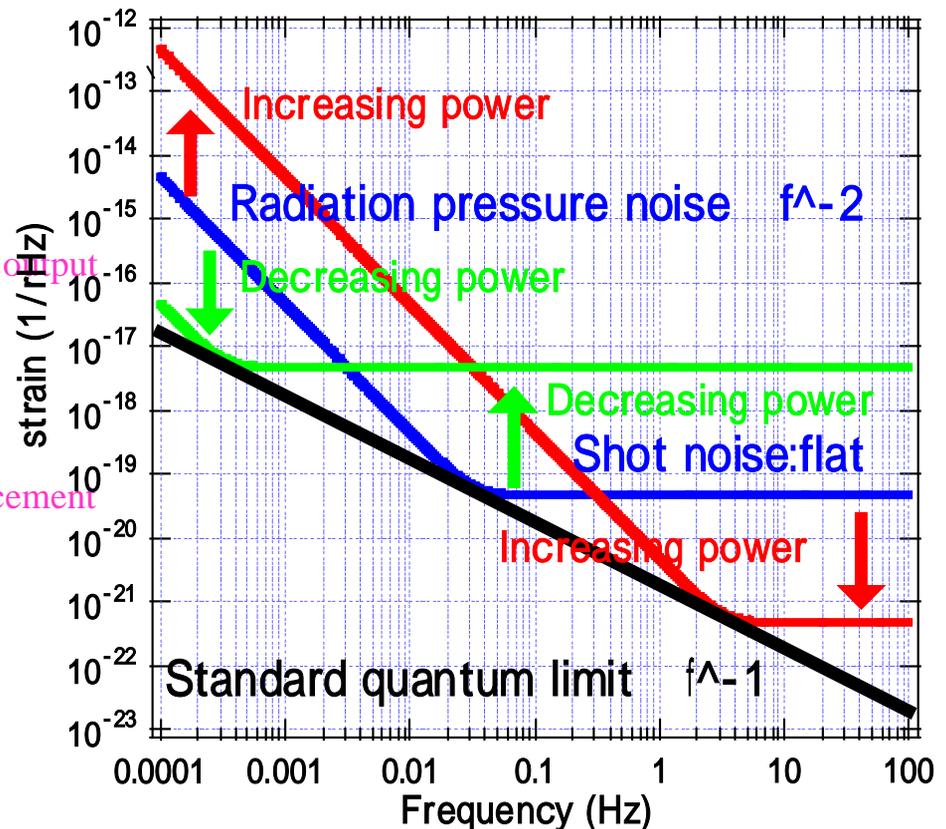
- Spectral density of the shot noise $1/P$,
- P : laser power
- The shot noise arises from uncertainty due to quantum mechanical fluctuations in the number of photons in the interferometer output

» Radiation pressure noise

- Spectral density of the radiation pressure noise P
- The radiation pressure arises from displacement induced by radiation-pressure fluctuations

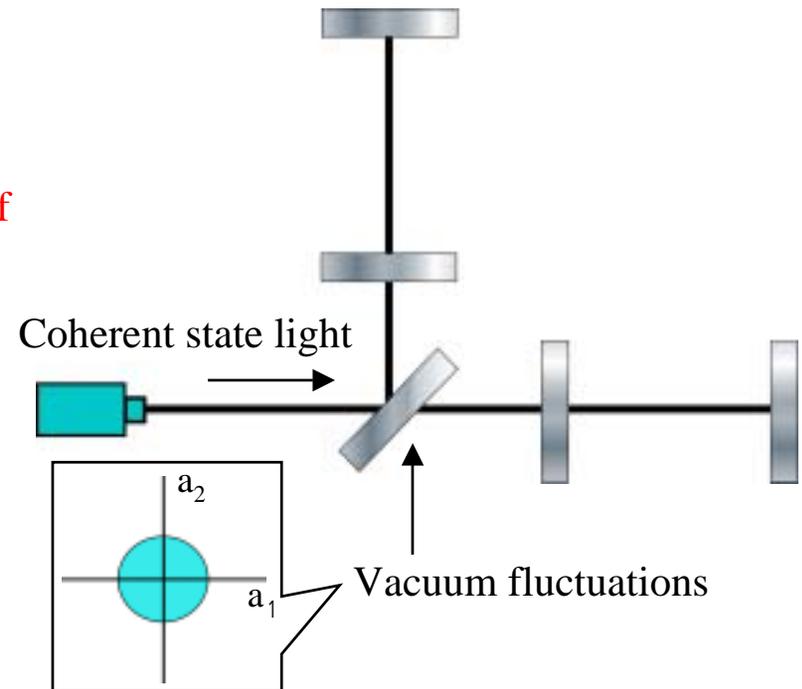
» Standard quantum limit

- A point at which the shot noise equals radiation pressure noise



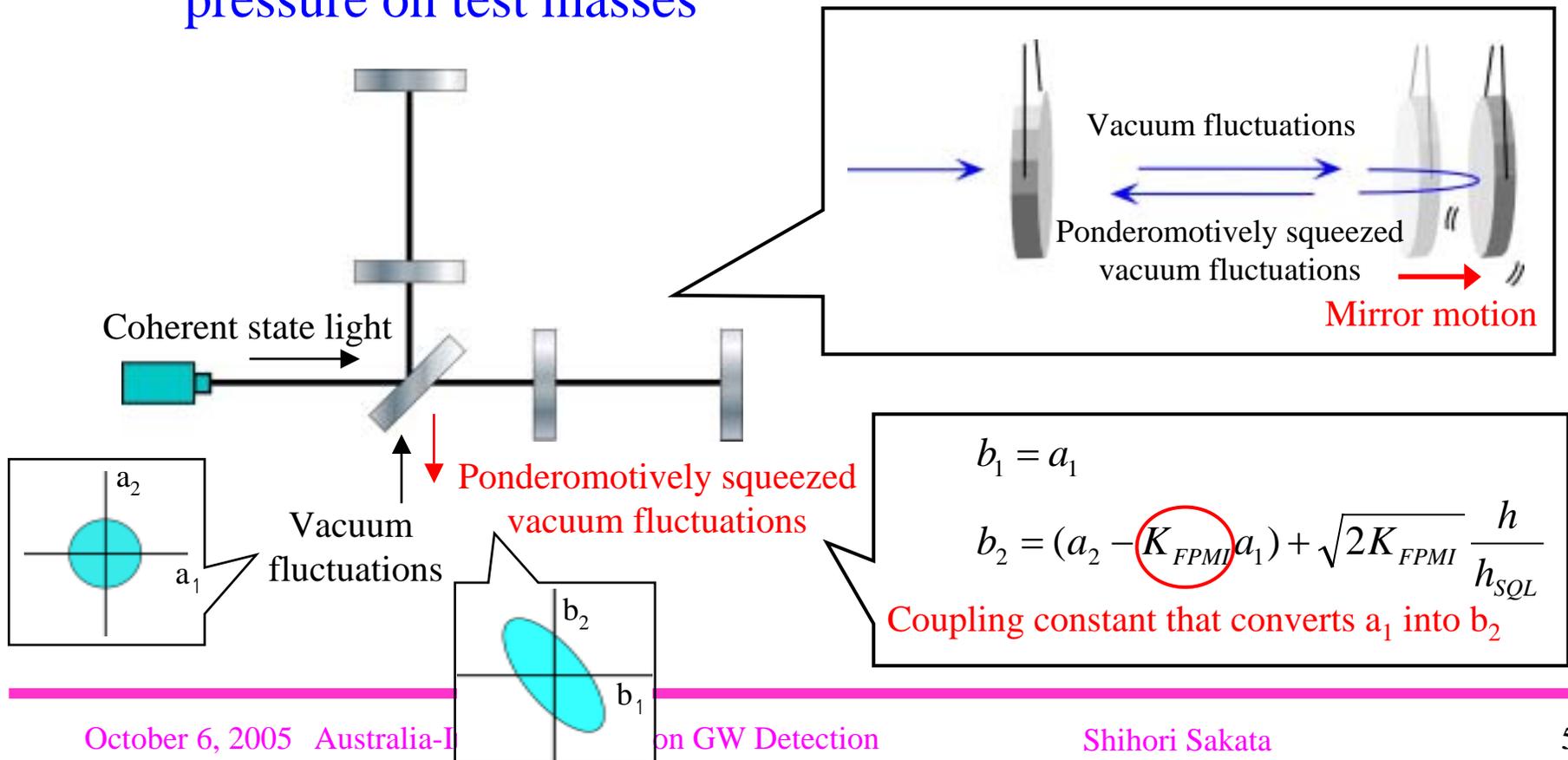
Vacuum fluctuations

- In the next generation, the sensitivity of the interferometer will be limited by the quantum noise at most of frequencies
 - » Vacuum fluctuations entering an anti-symmetric port of conventional interferometer
 - Expression of the vacuum fluctuations using two quadrature amplitudes $a_1(\omega)$ and $a_2(\omega)$ which are made by the combination of annihilation operators $a_{\pm}(\omega) = \frac{1}{\sqrt{2}}(a_1 \pm ia_2)$



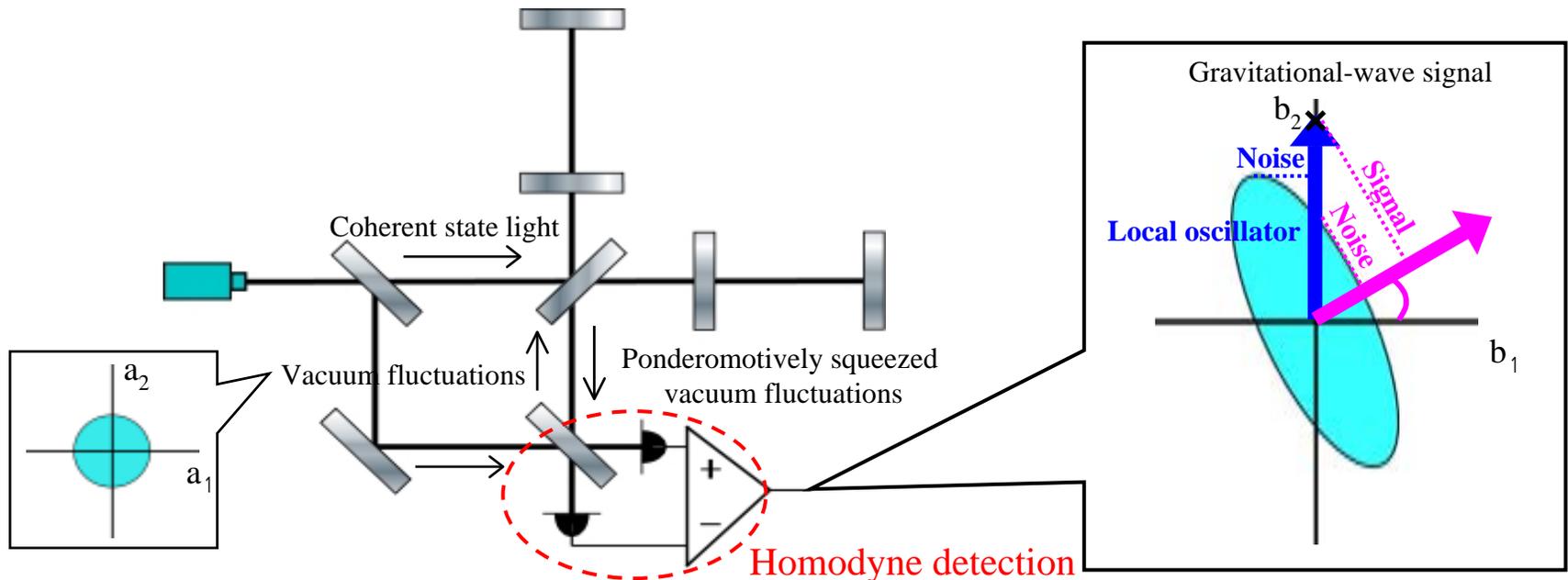
Ponderomotive squeezing

- The vacuum fluctuations is ponderomotively squeezed by back action of mirror motion due to fluctuating radiation pressure on test masses



Homodyne detection

- Using conventional readout scheme, in case of detecting the ponderomotively vacuum fluctuations, the sensitivity is limited by the SQL
- However, by using homodyne detection that is one of the quantum nondemolition devices, the sensitivity will be able to beat the SQL
 - » Optimization of the homodyne phase make it the best signal-to-noise ratio



2. Interferometer to generate ponderomotively squeezed vacuum fluctuations

- We calculate generation of the ponderomotive squeezing using Michelson interferometer and Fabry-Perot Michelson interferometer
 - » In the calculation conditions is that:
 - Parameters
 - Laser power $I_0=1\text{W}$
 - : Phase modulation frequency due to the mirror motion
 - =1kHz
 - - : cavity pole
 - Arm length=10cm
 - The sensitivity assumes to be limited by only the quantum noise
 - Loss is zero
 - End mirror reflectivity: 1
 - The sensitivity achieves the SQL, then beat the SQL by extracting the ponderomotive squeezing at the SQL
 - The homodyne phase is independent of frequency
 - » Reference: H.J.Kimble et al. Phys.Rev.D **65**, 022002 (2002)

Generation of ponderomotive squeezing using Michelson interferometer

- In case of measuring the output quadrature field by means of the conventional readout scheme,
 - » The noise spectral density: $S_h = \frac{h_{\text{SQL}}^2}{2} \left(\frac{1}{\mathcal{K}_{\text{MI}}} + \mathcal{K}_{\text{MI}} \right)$
 - » where $\mathcal{K}_{\text{MI}} = \frac{4I_0\omega_0}{m_r c^2 \Omega^2}$, $m_r = \frac{Mm}{M+m}$
 - The sensitivity achieves the SQL, where $\mathcal{K}_{\text{MI}} = 1$
 - m_r : reduced mass, M_{BS} : BS mass, ω_0 : laser angular frequency
- how much mass of the end mirror is needed to achieve the SQL?
 - » The end mirror mass: 1 ng
 - BS mass: 500 g
 - » The reduced mass: $m_r = \frac{4I_0\omega_0}{c^2\Omega^2}$
 - Experimentally generation of the ponderomotive squeezing is very difficult
 - » $m_r \propto I_0$

Generation of ponderomotive squeezing using Fabry-Perot Michelson interferometer

- In case of measuring the output quadrature field by means of the conventional readout scheme,

» The noise spectral density: $S_h = \frac{h_{\text{SQL}}^2}{2} \left(\frac{1}{\mathcal{K}_{\text{FPMI}}} + \mathcal{K}_{\text{FPMI}} \right)$

» where $\mathcal{K}_{\text{FPMI}} = \frac{4I_0\omega_0}{m_r r_F} \left(\frac{T_F \mathcal{F}}{\pi c \Omega (1 - r_F)} \right)^2$, $m_r = \frac{Mm}{M+m}$

– $m_E \propto \mathcal{F}^2$

- The sensitivity achieves the SQL, where $K_{\text{MI}} = 1$
- m_r : reduced mass, M_F : front mirror mass
- r_F : amplitude reflectivity of front mirror, T_E : intensity transitivity of end mirror

– Making finesse \mathcal{F} ten times is the same effect as making end mass m_E 100 times

- $\mathcal{F}: 10000$ $m_E: 500\text{g}$
- $\mathcal{F}: 1000$ $m_E: 5\text{g}$
- $\mathcal{F}: 100$ $m_E: 50\text{mg}$

Experimental parameter

■ Discussion of experimental parameter

» The sensitivity achieves the SQL at $\omega_{\text{SQL}} = 1 \text{ kHz}$

» The noise spectral density:
$$S_h = \frac{h_{\text{SQL}}^2}{2} \left(\frac{1}{\mathcal{K}_{\text{FPMI}}} + \mathcal{K}_{\text{FPMI}} \right)$$

– Finesses \mathcal{F} 1/10 times and mass m 1/100 times does not change the frequency

• Radiation pressure noise $1/m$, \mathcal{F}

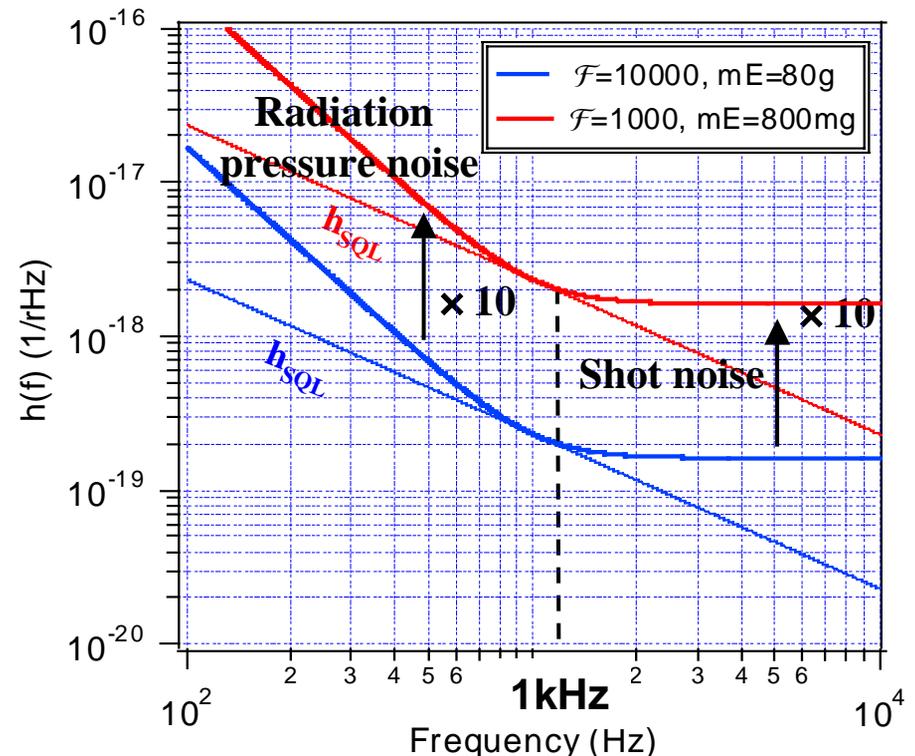
• Shot noise $1/\mathcal{F}$

– By making the sensitivity worse, achievement of the SQL is more possible

■ Noises other than the quantum noise

» Suspension thermal noise becomes worse

– Roughly $1/(m)^{1/2}$



Beat the SQL

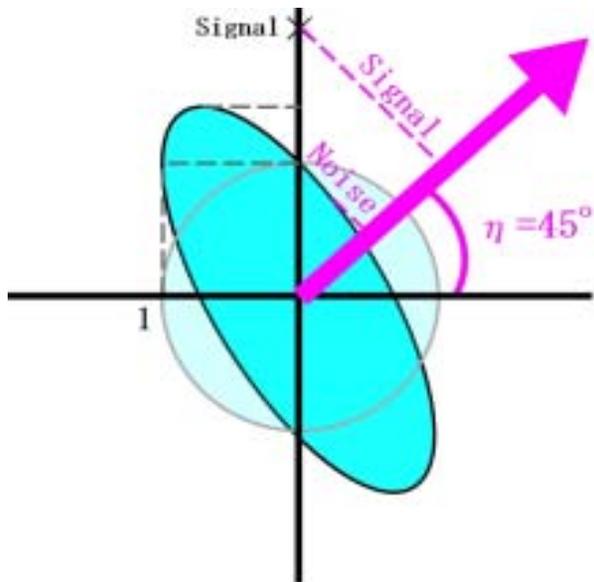
- In case of measuring the output quadrature field by homodyne detection,

» The sensitivity is obtained by $S_h = \frac{h_{\text{SQL}}^2}{2\mathcal{K}_{\text{FPMI}}} \left(1 + (\cot \eta - \mathcal{K}_{\text{FPMI}})^2\right)$

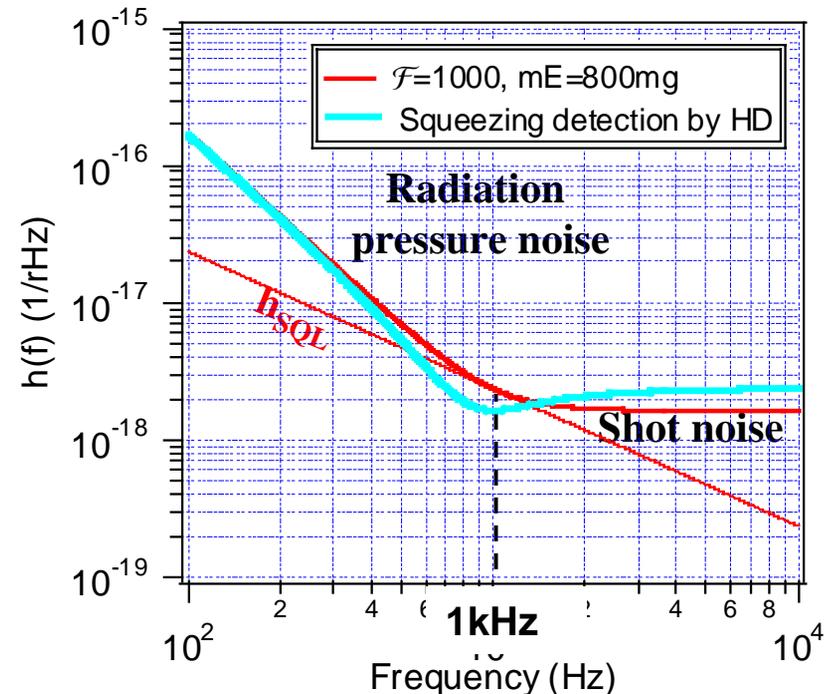
» At the SQL, the homodyne phase is optimized by

– Ponderomotive squeezing of 4dB @

• Squeezing factor: $\text{arcsinh}(\mathcal{K}_{\text{FPMI}}/2)$



$S_{\text{QL}} = 1\text{kHz}$



3. Summary

- By extraction of the ponderomotive squeezing at the SQL, the sensitivity will beat the SQL
 - » Discussion of interferometer to generate the ponderomotive squeezing
 - Michelson interferometer and Fabry-Perot Michelson interferometer
 - Loss is zero
 - The sensitivity assumes to be limited by only the quantum noise
 - Finesses \mathcal{F} 10 times is equivalent to the mass m 100 times
 - » Experimental parameter
 - The sensitivity achieves the SQL at $\omega_{\text{SQL}}=1\text{kHz}$
 - Finesses \mathcal{F} 1/10 times and mass m 1/100 times does not change the frequency
 - At the SQL, $K_{\text{FPMI}} = 1$
 - By making the sensitivity worse, achievement of the SQL is more possible
 - » At the SQL, the homodyne phase is optimized by
 - Ponderomotive squeezing of 4 dB @ $\omega_{\text{SQL}} = 1\text{kHz}$
 - Homodyne phase $\phi = 45^\circ$
 - $m_E = 800\text{mg}$, $\mathcal{F} = 1000$

4. Future plan

- Design of end mirrors
 - » Design of suspension
 - » Thermal noise
- Photo detection
 - » Homodyne detection or DC readout