### Some Ideas on Coatingless all-reflective ITF

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#### Why Total Reflecting Mirrors?

## Experimental data from Yamamoto and Numata show a very high coating loss angle even at low temperature.



The Virgo design Thermal noise curve is evaluated with a Q~10<sup>6</sup>, and the maximum expected Q, due to coating losses, is Q~2.5 10<sup>7</sup>. Standard Quantum Limit is about a factor 100 below Virgo TN i.e. equivalent to a Q=10<sup>10</sup> ; consequently, for reaching SQL sensitivity Q should improve by  $10^{10}/ 2.5 \ 10^7 \sim 400$ - For this reason it is interesting to explore Coating-less Mirrors



#### **Some History**

Toraldo di Francia in 1965 proposes flat Roof Prism for creating a stable Radio Frequency cavity. Stability was succesfully experimentally tested.



In 1970 Robert Forward (Hughes Lab.) build the first Interferometer for GW detection. It was equipped with Roof Prisms Mirrors. Arm length ~2m

Braginsky et al. Recently Published a paper on the use of Roof Prism and Corner Cube mirrors in Fabry Perot cavities using Antireflective Coatings.

# Rotation Parabolas as exact reflectors for closed geometrical optical trajectories

![](_page_4_Figure_1.jpeg)

![](_page_4_Figure_2.jpeg)

## Rotation Parabolas as reflectors for closed geometrical optics trajectories

![](_page_5_Figure_1.jpeg)

### **Reflective cavity with Asymmetric Arms**

![](_page_6_Figure_1.jpeg)

**Cavity Beam Injection One of the main problems we faced was to find the way to inject a beam in the All Reflective Cavity.** 

![](_page_7_Figure_1.jpeg)

#### **Example of Power Injection in a Parabolic Coatingless all-Reflective Cavity**

![](_page_8_Figure_1.jpeg)

#### **A REALISTIC CAVITY**

![](_page_9_Figure_1.jpeg)

**Some Peculiar Properties-1** If channel L3-L4 is antiresonating 1) Then  $\frac{\partial U}{\partial L_3} = \frac{\partial U}{\partial L_4} = 0$  i.e. Thermal noise of mirror M<sub>3</sub> does not affect U. Hence for stopping input of vacuum we can use a normally coated mirror **2)**Consequently the amplitude C=0. i.e. there is no power on  $M_3$ 3)Since  $U = e^{i\varphi}$  the Finesse F. defined as  $F = \frac{\partial \varphi}{\partial L_1} \frac{\lambda}{2\pi} \Rightarrow -\frac{i}{U} \frac{\partial U}{\partial L_1} \frac{\lambda}{2\pi}$ ,  $L_1$ M<sub>3</sub>  $-R_3, T_3 R_3, T_3$ becomes F d  $L_3$  $L_4$ b g  $R_4, T_4$  $U = e^{i\varphi}$  $-R_1, T_1, R_1, T_1$ 

![](_page_11_Figure_0.jpeg)

#### Silica Reflectivity

![](_page_12_Figure_1.jpeg)

#### DISCUSSION

By considering that  $R_1^2 \ll 1$ , by stopping expansion to  $\chi^2$  terms and by putting in evidence  $F = \frac{2}{R^2}$ , we obtain:

$$\delta\varphi = \frac{2}{R_1^2} \left[ \left( 1 + \frac{1 - R_1^2}{1 + R_1^2} \chi^2 \right) \frac{\omega_0}{c} \delta L_1 - R_1^2 \frac{\left( -1 + R_1^2 \right)}{\left( 1 + R_1^2 \right)^2} \chi^2 \frac{\omega_0}{c} \delta L_4 \right] \right]$$

This equation shows that we may change the finesse by the factor  $\Delta F=1+\chi^2$  and still the phase shift due to thermal noises of M<sub>3</sub> in L<sub>4</sub> is depressed by the factor  $\frac{R_1^2\chi^2}{1+\chi^2} \ll 1$  because R<sub>1</sub><sup>2</sup> $\ll 1$ .

Then it seems that in this cavity system we may change finesse, and consequently storage time, without affecting :

1)Thermal noise from channel  $L_3$ - $L_4$ , which for  $\chi^2=0$  is zero.

2)Cavity geometry.

The independence of geometry from finesse seems to be a very important feathure of this system.

![](_page_14_Figure_0.jpeg)

#### **Thermal Noise**

With this optical configuration it is evident that every photon traverses a large amount of material; here the problem is transferred from normal mirror reflective coating bad loss angle to thermorefractive noise (TRN) of coatingless cavity. TRN spectral density is:

$$S_{\phi}(\omega) = \frac{16\pi L_{bulk} k_B T^2 \kappa}{\lambda^2 \rho^2 C^2 w_0^4 \omega} \beta^2$$

Where:  $L_{bulk}$  is the thickness of traversed material,  $w_0$  the beam waist,  $\rho$  the bulk density,  $\kappa$  the bulk thermal conductivity coefficient,  $\lambda$  the wavelength, C the bulk thermal capacity and  $\beta$  is:

$$\beta = \frac{\partial n}{\partial T} = \frac{\left(n^2 - 1\right)\left(n^2 + 1\right)}{6n} \left[\frac{1}{\alpha_p}\frac{\partial \alpha_p}{\partial T} - 3\alpha \left(1 + \frac{\rho}{\alpha_p}\frac{\partial \alpha_p}{\partial \rho}\right)\right]$$

Where  $\alpha_p$  is the bulk polarizability, **n** the bulk refractive index and  $\alpha$  is the bulk linear thermal expansion coefficient.

# We tried to find data on $\beta$ at different temperature and we only found 300K data for Silica, which has large C and small $\kappa$ .

![](_page_16_Figure_1.jpeg)

Silica 300K

Data for Sapphire give bad results since C is small and κ is very large compared with Silica

Cristaline Quartz could be a good candidate for going to low temperature but we could not find any data.

### Conclusions

- This coatingless optical configuration has some interesting feathures:
- 1) Since all the surfaces are optically matched, antireflective coatings are not needed.
- 2) Power can be injected in the cavity in a straightforward way; use of tunnel effect injection or transversal cavity are unecessary.
- **3)** The optical scheme can be made with variable finesse. This property can also applied also to "normal" mirror configuration.
  - It is evident that the draw back is the Thermo Refractive noise; an investigation for selecting cristaline bulks with low thermal conductivity and large thermal capacity is instrumental.