

Talk given at LIGO
Friday October 3, 1997

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DIFFRACTION EFFECTS AT SUSI

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LIGO - G970293-00-R

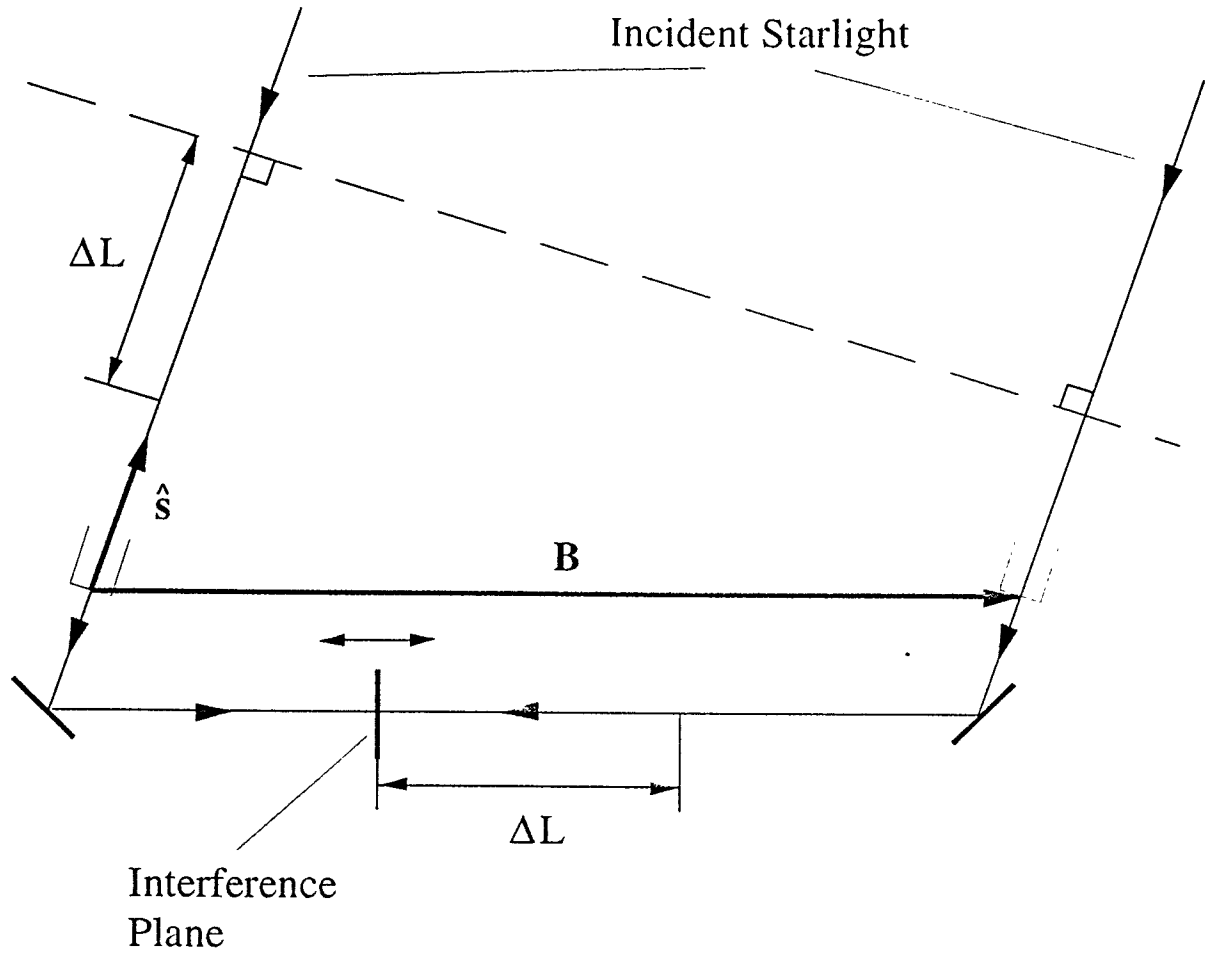


School of Physics
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Talk Outline

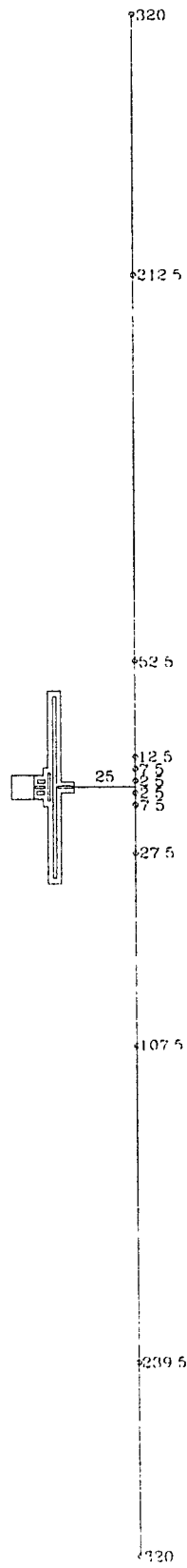
- Overview of SUSI
- Diffraction Theories
- The Fresnel Approximation
- Visibility Loss Calculation



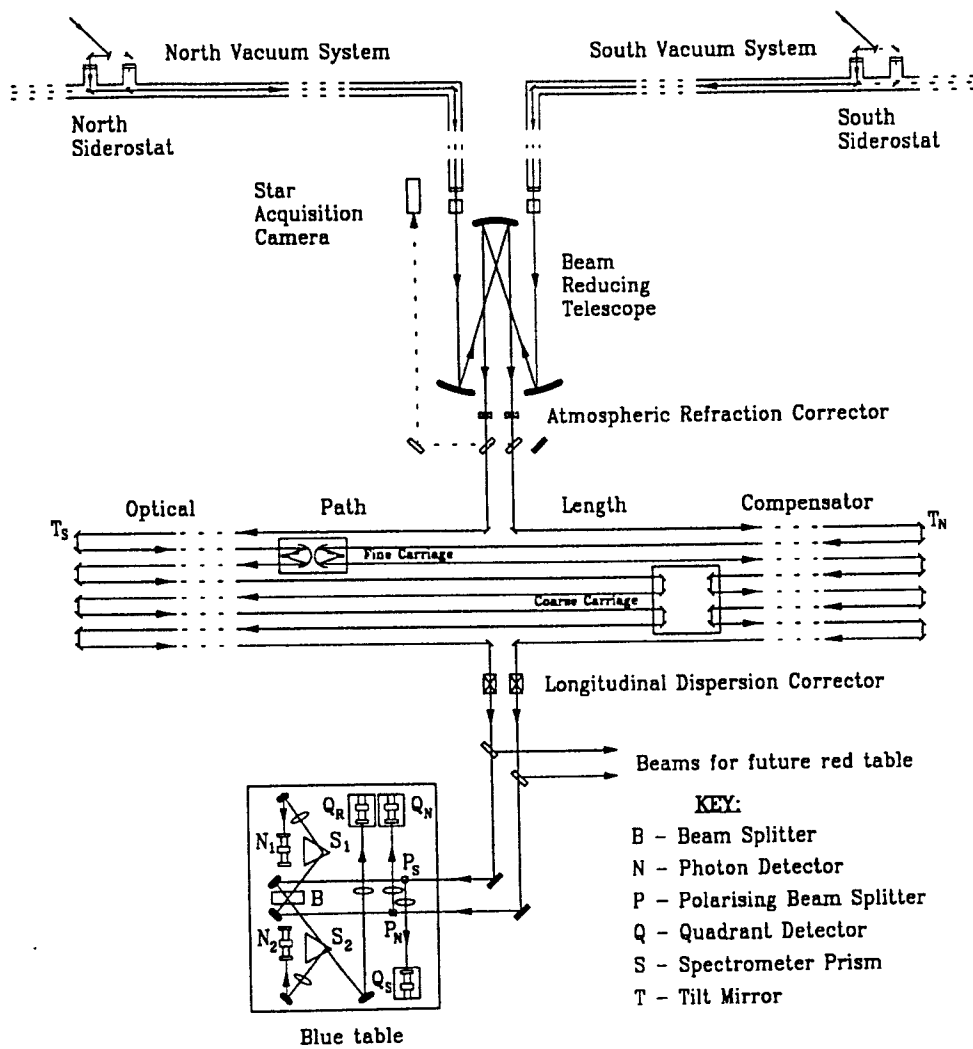
$$I(\mathbf{x}) = I_1(\mathbf{x}) + I_2(\mathbf{x}) + 2\sqrt{I_1(\mathbf{x})I_2(\mathbf{x})} \times |\gamma(\tau)| \cos((\phi_1(\mathbf{x}) - \phi_2(\mathbf{x})) + \delta)$$

$$\mathcal{V} = \frac{I_{max} - I_{min}}{I_{max} + I_{min}}$$

$$\mathcal{V} = \frac{2 \int_0^\infty \sqrt{I_1(\mathbf{x})I_2(\mathbf{x})} \cos(\phi_1(\mathbf{x}) - \phi_2(\mathbf{x})) d\mathbf{x}}{\int_0^\infty (I_1(\mathbf{x}) + I_2(\mathbf{x})) d\mathbf{x}} |\gamma(\tau)|$$



Dimensions shown in metres



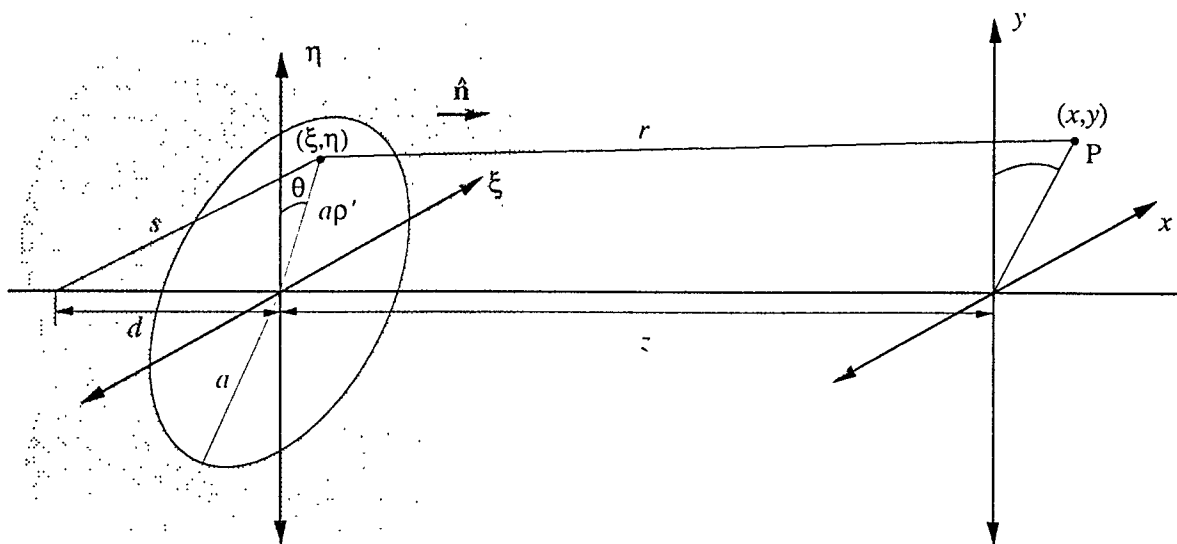
Diffraction Theories

- Electromagnetic (Vector)
- Scalar
 - Kirchhoff
 - Rayleigh-Sommerfeld
- Geometric Theory of Diffraction (GTD)

Practical Considerations

- Description of relevant features of wave propagation and diffraction effect.
- Relative ease of calculation.
- Possible need to incorporate vignetting from cascaded apertures.
- Would it be possible to incorporate an atmospherically distorted wavefront?

The Fresnel Approximation



$$r = \sqrt{z^2 + (x - a\rho \sin \theta)^2 + (y - a\rho \cos \theta)^2}$$

$$U(P) = \frac{-A}{2\lambda} \int \frac{e^{ikr}}{r} \left(1 + \frac{z}{r}\right) dS$$

Apply Fresnel Approximation

$$r \approx z \left(1 + \frac{(x - a\rho \sin \theta)^2 + (a\rho \cos \theta)^2}{2z^2} - \frac{((x - a\rho \sin \theta)^2 + (a\rho \cos \theta)^2)^2}{8z^4} + \dots \right)$$

Introduce the optical coordinates u, v

$$u = \frac{2\pi a^2}{\lambda z}$$

$$v = \frac{2\pi ax}{\lambda z}$$

$$\rho = \frac{v}{u}$$

$$U(P) = -iuAe^{i\left(kz + \frac{v^2}{2u}\right)} \int_0^1 e^{i\frac{u}{2}\rho'^2} J_0(v\rho') \rho' d\rho'$$

$$U(P) = iAe^{i\left(kz + \frac{u}{2} + \frac{v^2}{2u}\right)} [U_1(u, v) - iU_2(u, v)]$$

where the Lommel functions are defined as

$$U_n(u, v) = \sum_{p=0}^{\infty} (-1)^p \left(\frac{u}{v}\right)^{n+2p} J_{n+2p}(v)$$

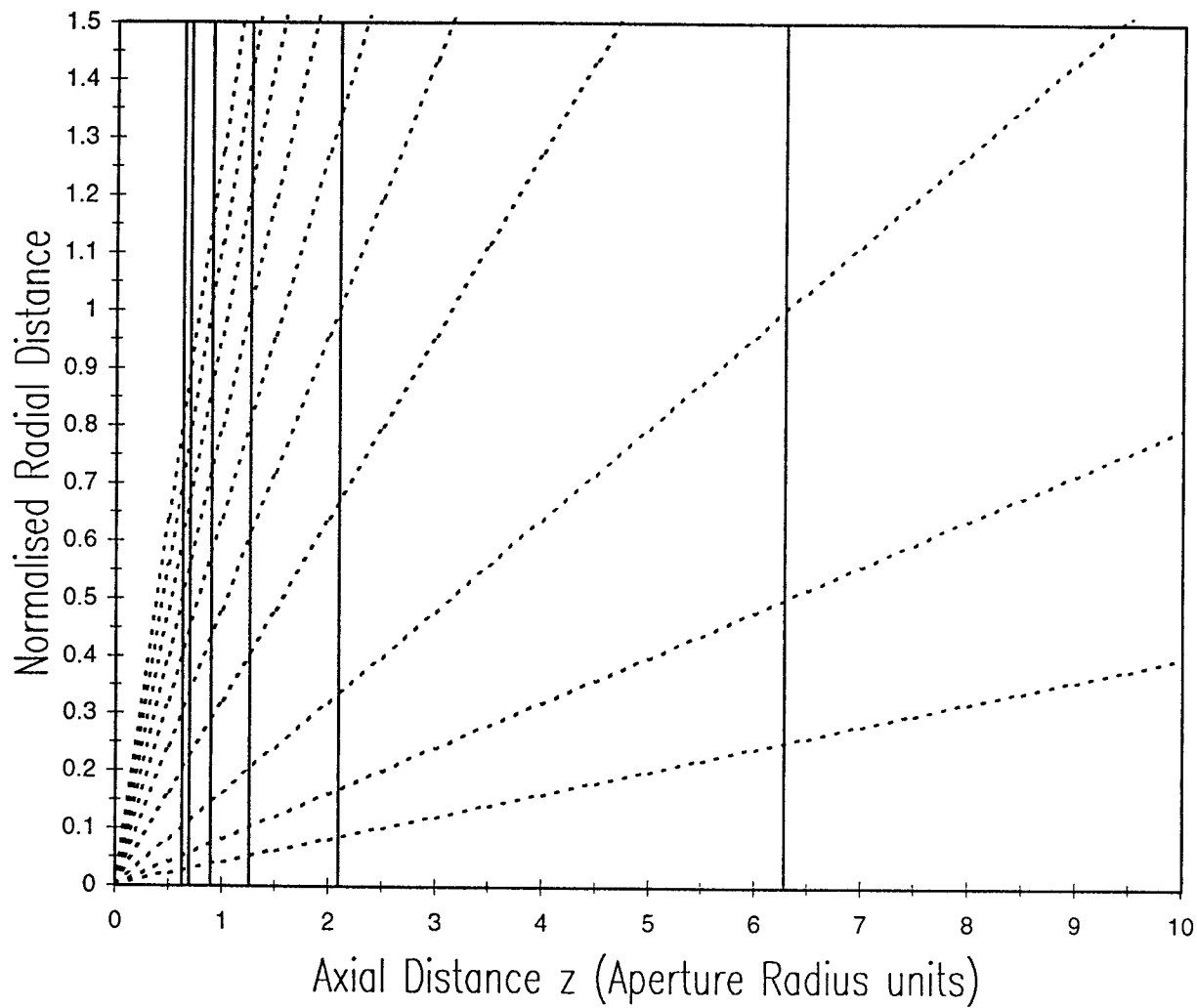
Intensity is simply expressed as

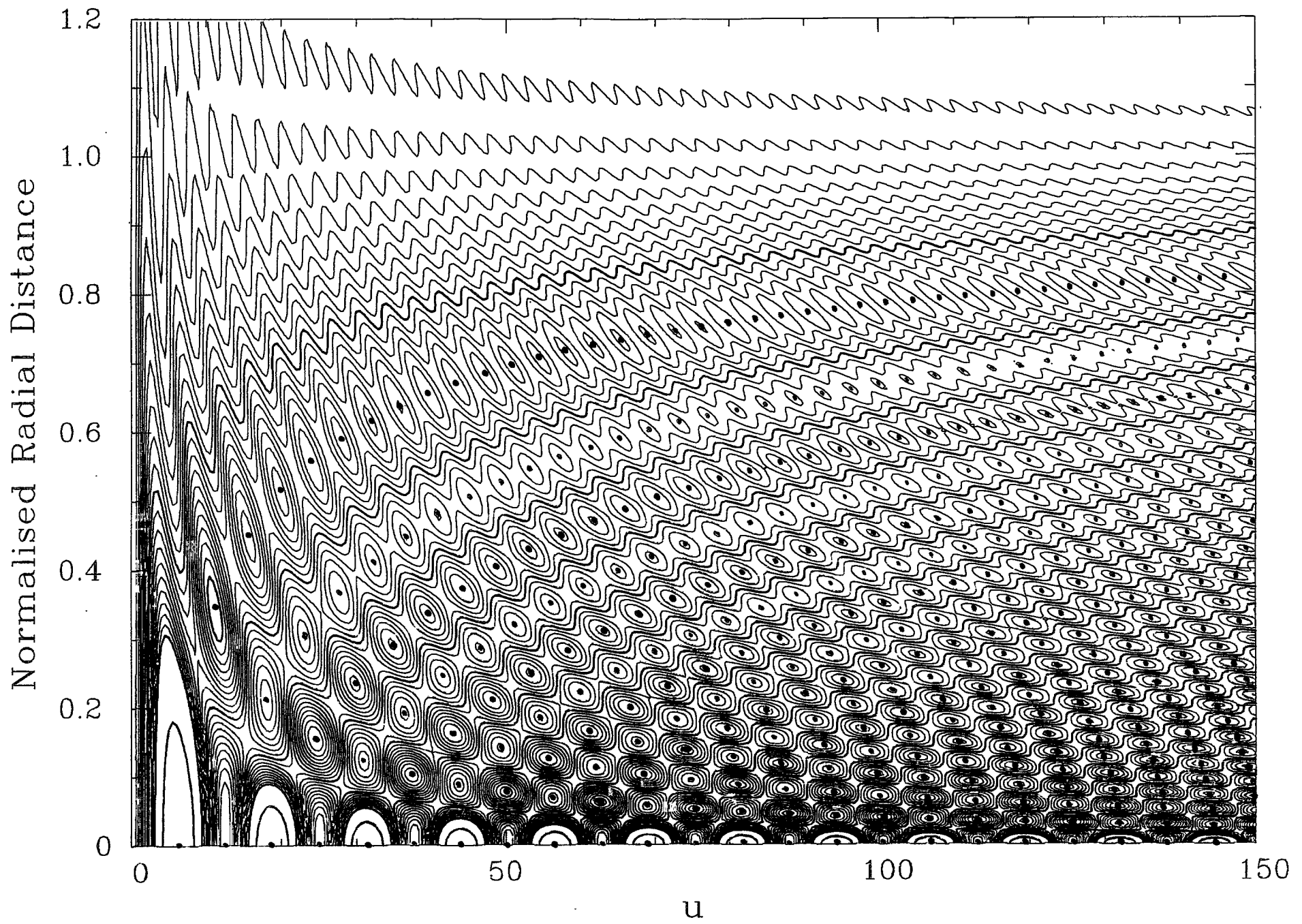
$$I(u, v) = \left(U_1^2(u, v) + U_2^2(u, v)\right) I_0$$

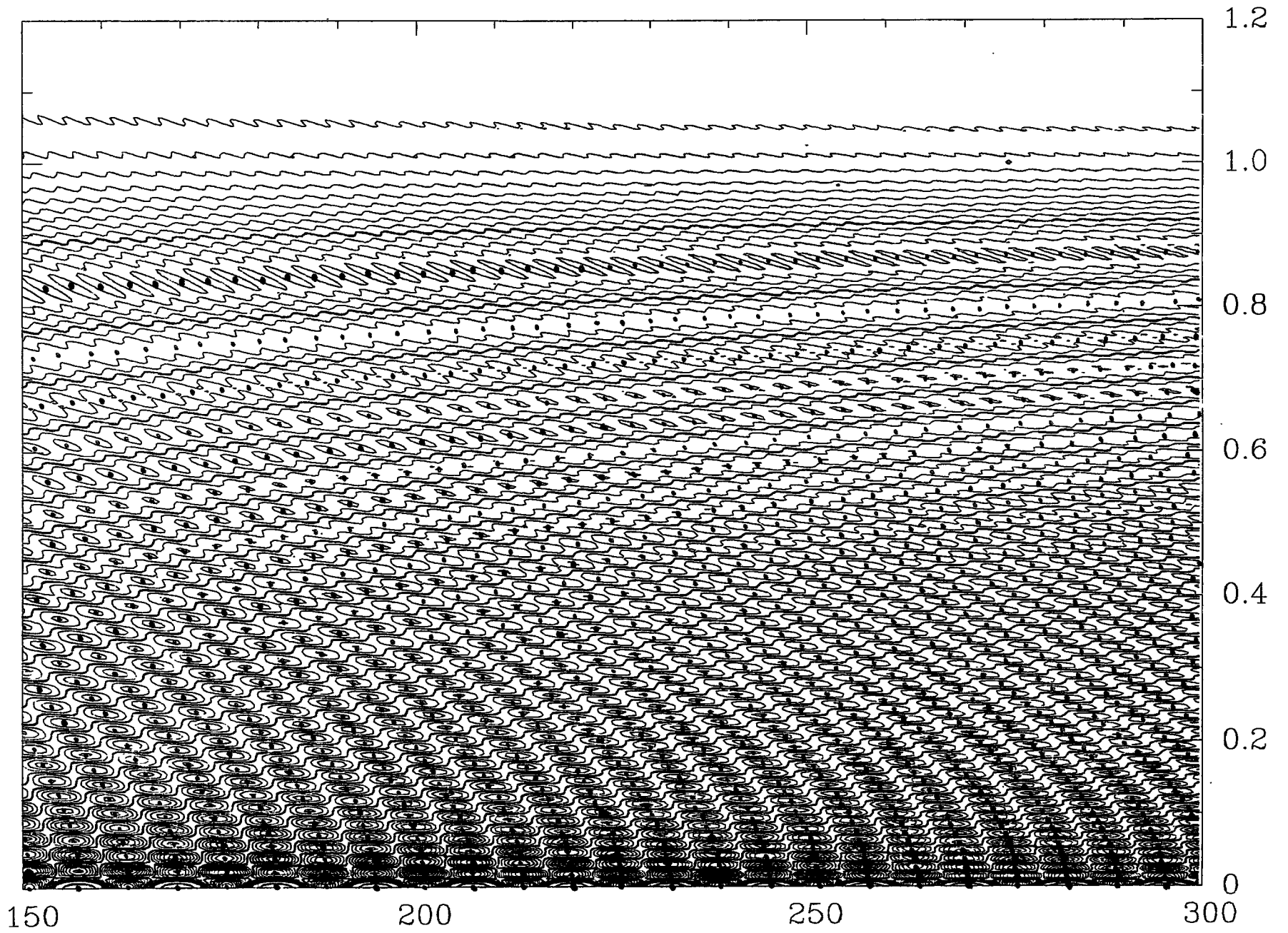
and Phase by

$$\phi(u, v) = kz + \frac{u}{2} + \frac{v^2}{2u} - \tan^{-1} \left(\frac{U_2(u, v)}{U_1(u, v)} \right) + \frac{\pi}{2}$$

Contours of Constant u and v







150

200

250

300

Lommel Functions

$$U_n(u, v) = \sum_{p=0}^{\infty} (-1)^p \left(\frac{u}{v}\right)^{n+2p} J_{n+2p}(v)$$

$$V_n(u, v) = \sum_{p=0}^{\infty} (-1)^p \left(\frac{v}{u}\right)^{n+2p} J_{n+2p}(v)$$

$$U_1(u, v) = \sin \left(\frac{u}{2} + \frac{v^2}{2u} \right) - V_1(u, v)$$

$$U_2(u, v) = -\cos \left(\frac{u}{2} + \frac{v^2}{2u} \right) + V_0(u, v)$$

At the shadow boundary (where $u = v$)

$$U_1(u, u) = \frac{1}{2} \sin u$$

$$U_2(u, u) = \frac{1}{2} (J_0(u) - \cos u)$$

Along the shadow boundary

$$I(u, u) = \frac{I_0}{4} \left(1 - 2J_0(u) \cos u + J_0^2(u) \right)$$

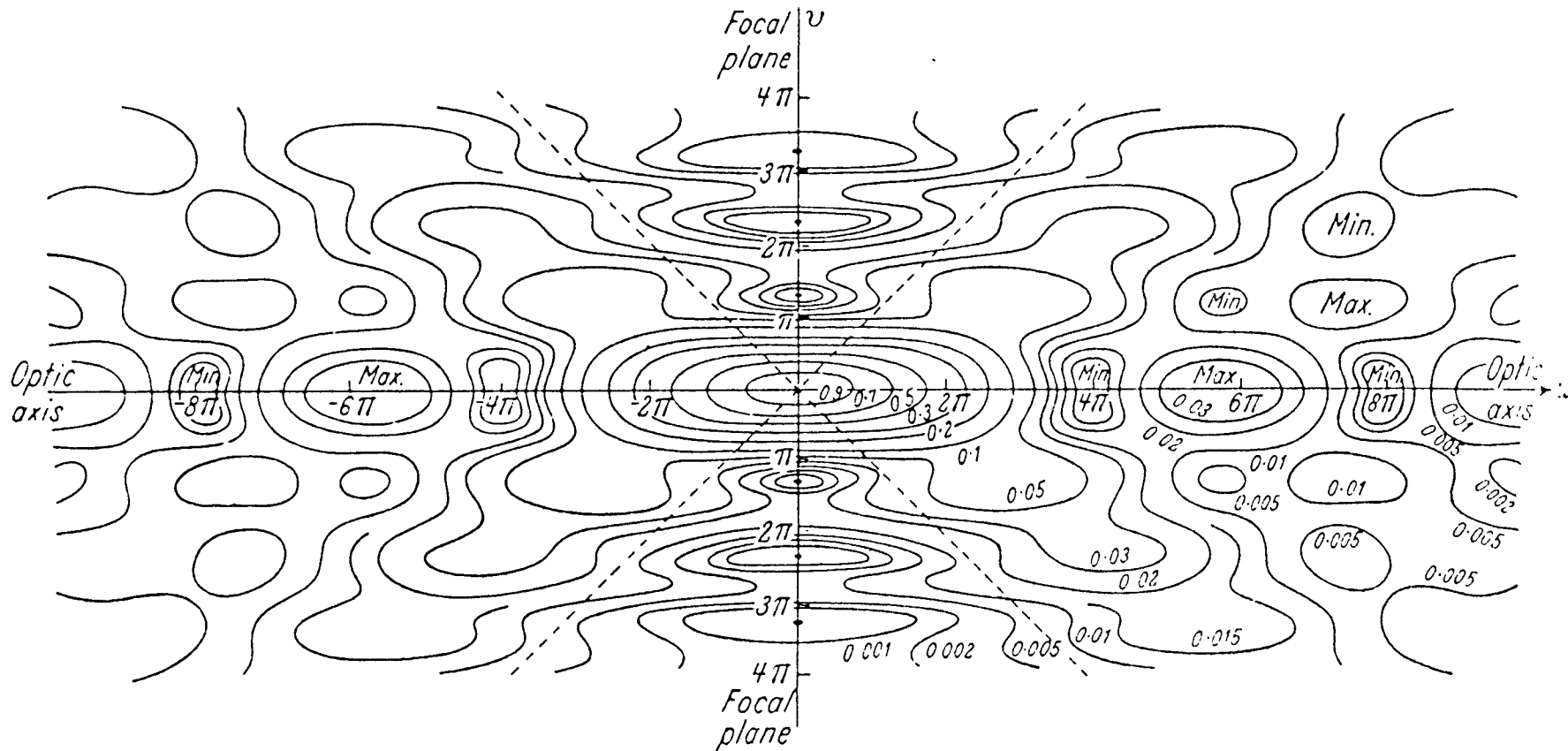
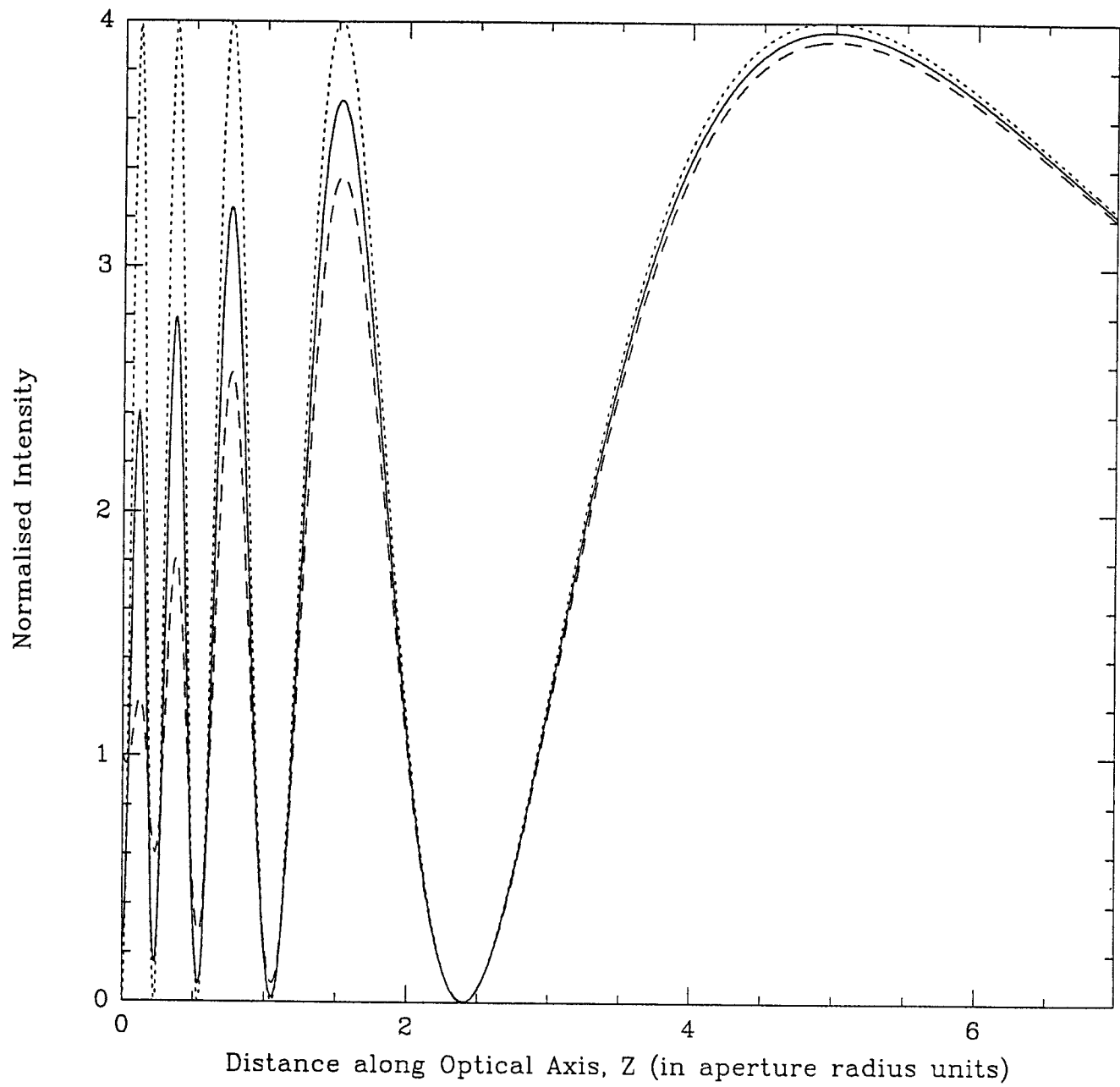


Fig. 8.41. Isophotes [contour lines of the intensity $I(u, v)$] in a meridional plane near focus of a converging spherical wave diffracted at a circular aperture. The intensity is normalized to unity at focus. The dotted lines represent the boundary of the geometrical shadow. When the figure is rotated about the u -axis, the minima on the v -axis generate the AIRY dark rings.

(Adapted from E. H. LINFOOT and E. WOLF, *Proc. Phys. Soc., B*, **69** (1956), 823.)



Axial Comparisons

Exact Axial Solutions

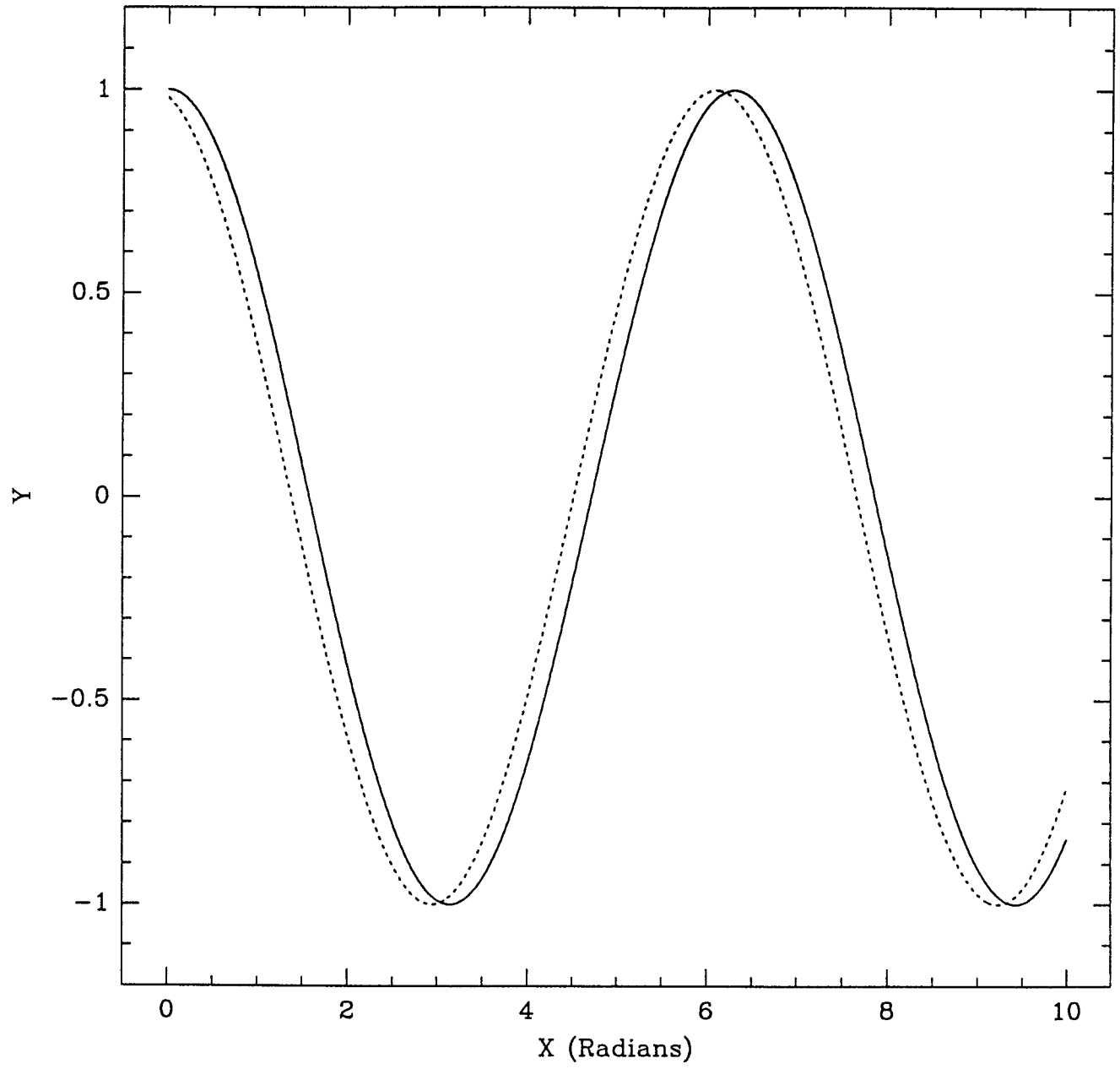
$$I_K(z) = 1 - \left(\frac{z}{\sqrt{z^2 + a^2}} + 1 \right) \cos \left(k(\sqrt{z^2 + a^2} - z) \right) + \frac{1}{4} \left(\frac{z}{\sqrt{z^2 + a^2}} + 1 \right)^2$$

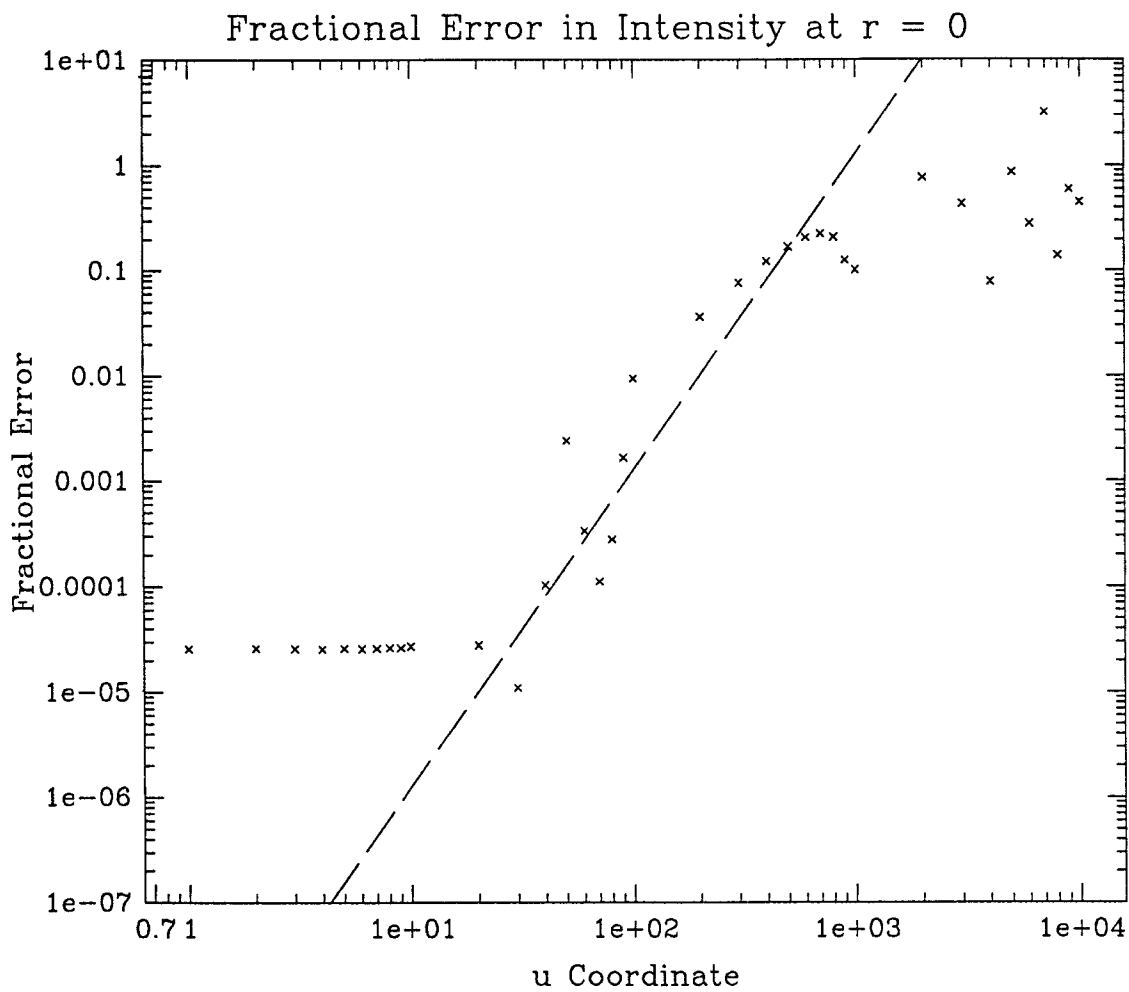
Number of axial maxima is given by a/λ

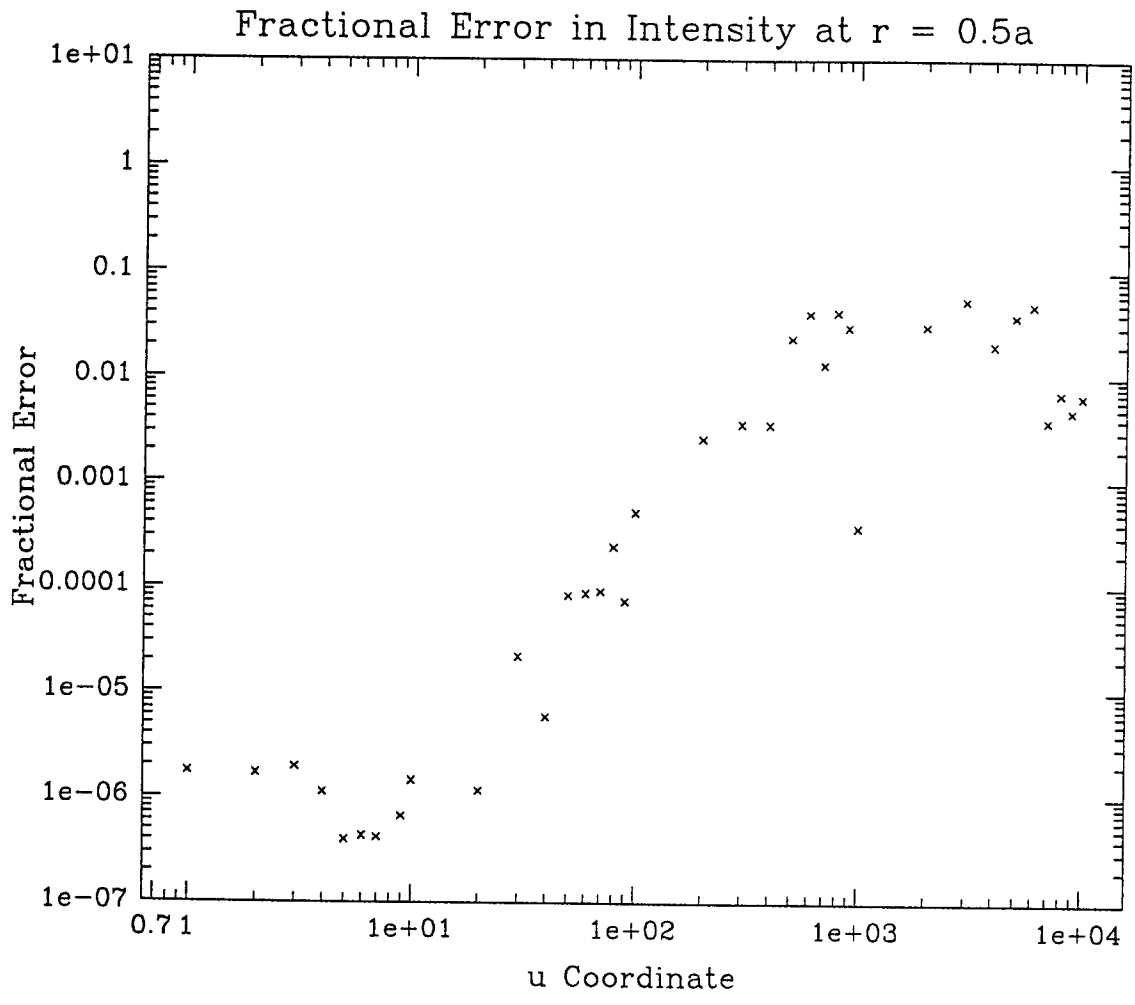
Fresnel Axial Solution

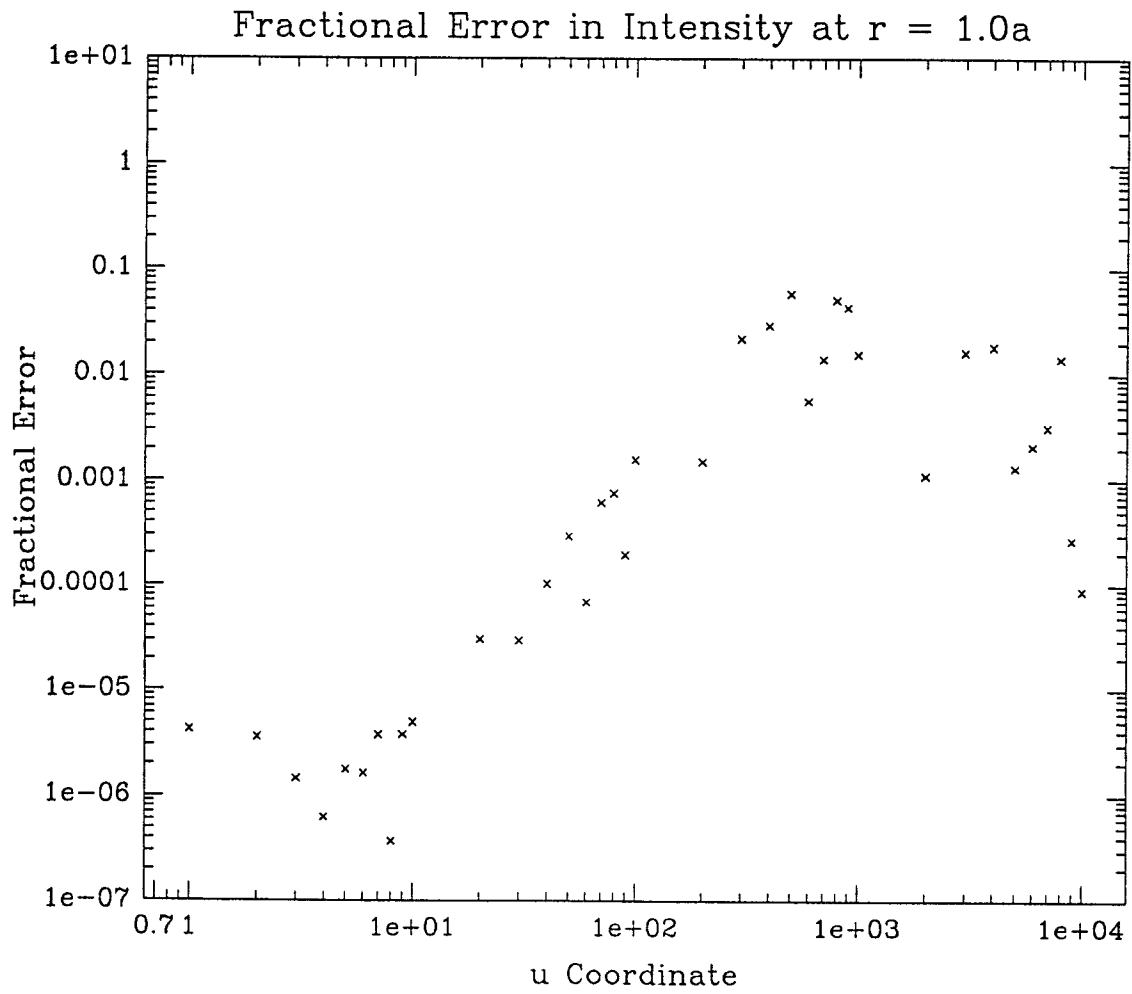
$$\begin{aligned} I(u, 0) &= 4 \sin^2 \left(\frac{u}{4} \right) \\ &= 4 \sin^2 \left(\frac{\pi a^2}{2\lambda z} \right) \end{aligned}$$

$$\epsilon \approx \frac{ka^4}{z^3}$$









Improving the Approximation

$$\begin{aligned} kr &= k\sqrt{z^2 + a^2\rho^2 + a^2\rho'^2 - 2a^2\rho\rho' \sin \theta} \\ &\approx \frac{1}{2}U\rho'^2 - V\rho' \sin \theta + W \end{aligned}$$

Exterior stationary points at $\rho' = 1, \theta = 0$.

Interior stationary point at $d(kr) = 0$

$$\begin{aligned} d(kr) &= \frac{\partial(kr)}{\partial\rho'}d\rho' + \frac{\partial(kr)}{\partial\theta}d\theta \\ &= (U\rho' - V \sin \theta)d\rho' - (V \cos \theta)d\theta \\ \rho' &= \rho = \frac{V}{U}, \theta = 0 \end{aligned}$$

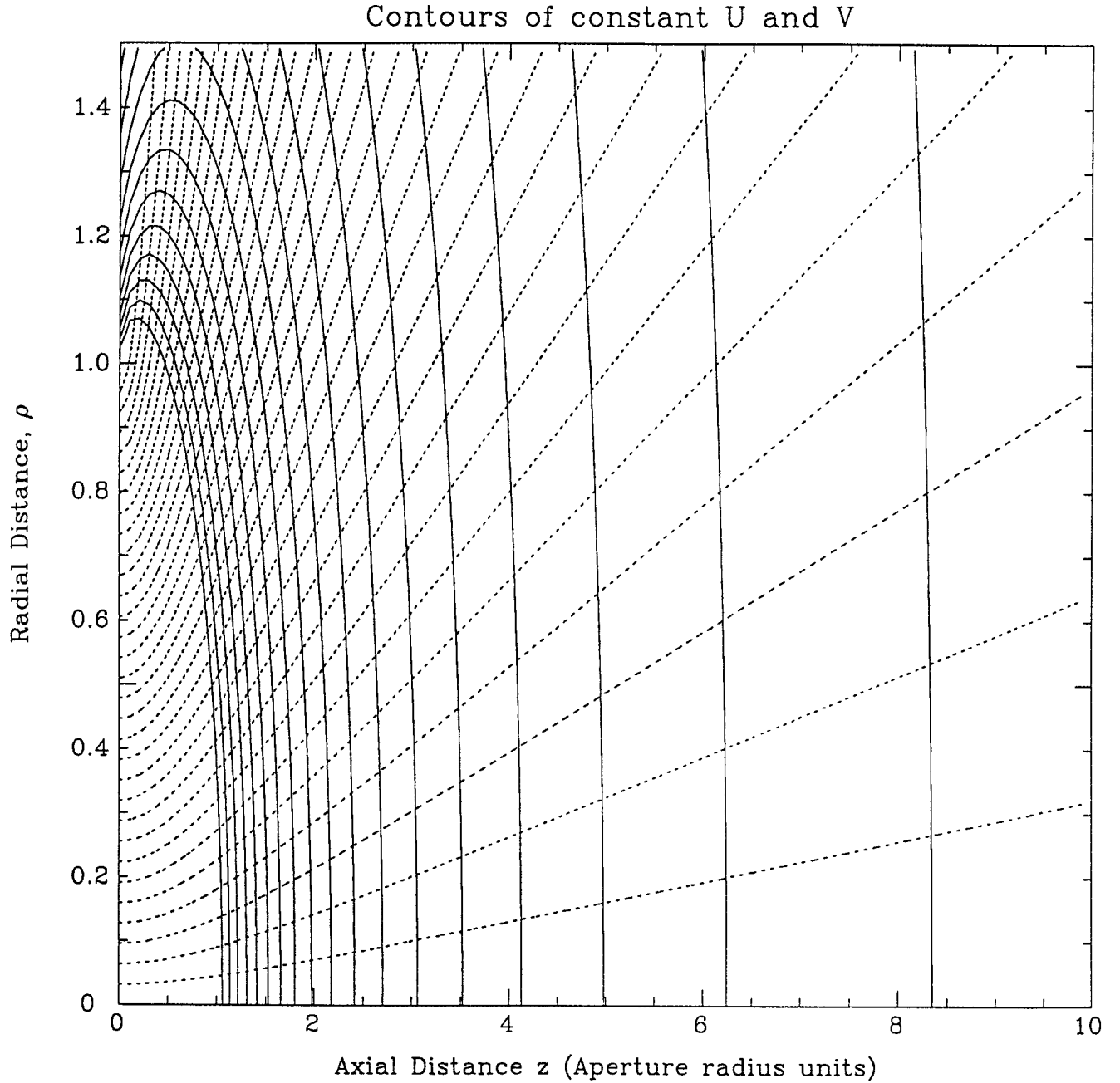
$$U = \frac{k}{2}(p + q)^2,$$

$$V = \frac{k}{2}(p^2 - q^2),$$

$$W = \frac{k}{4}(p^2 - q^2) + kz$$

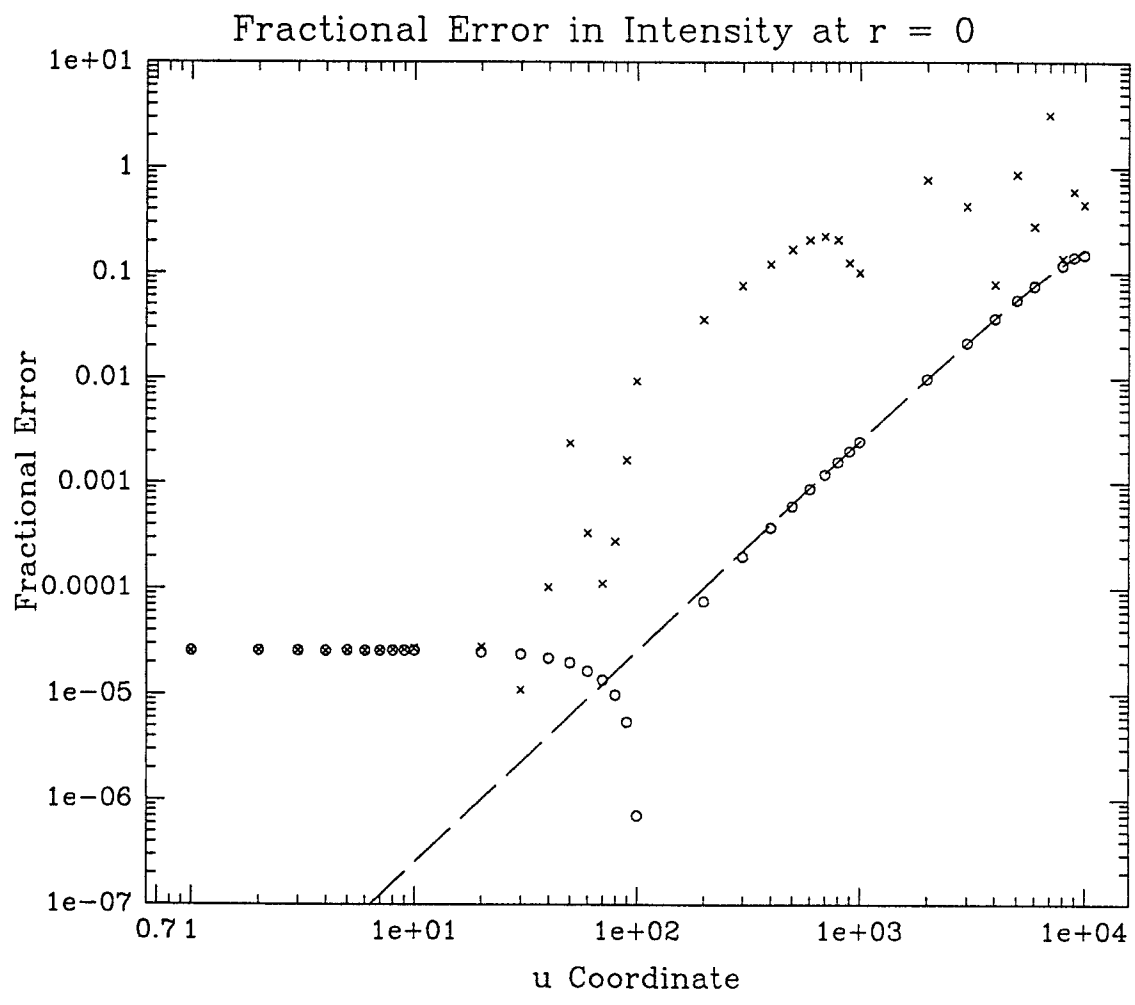
$$p = \sqrt{\sqrt{z^2 + a^2(1 + \rho)^2} - z}$$

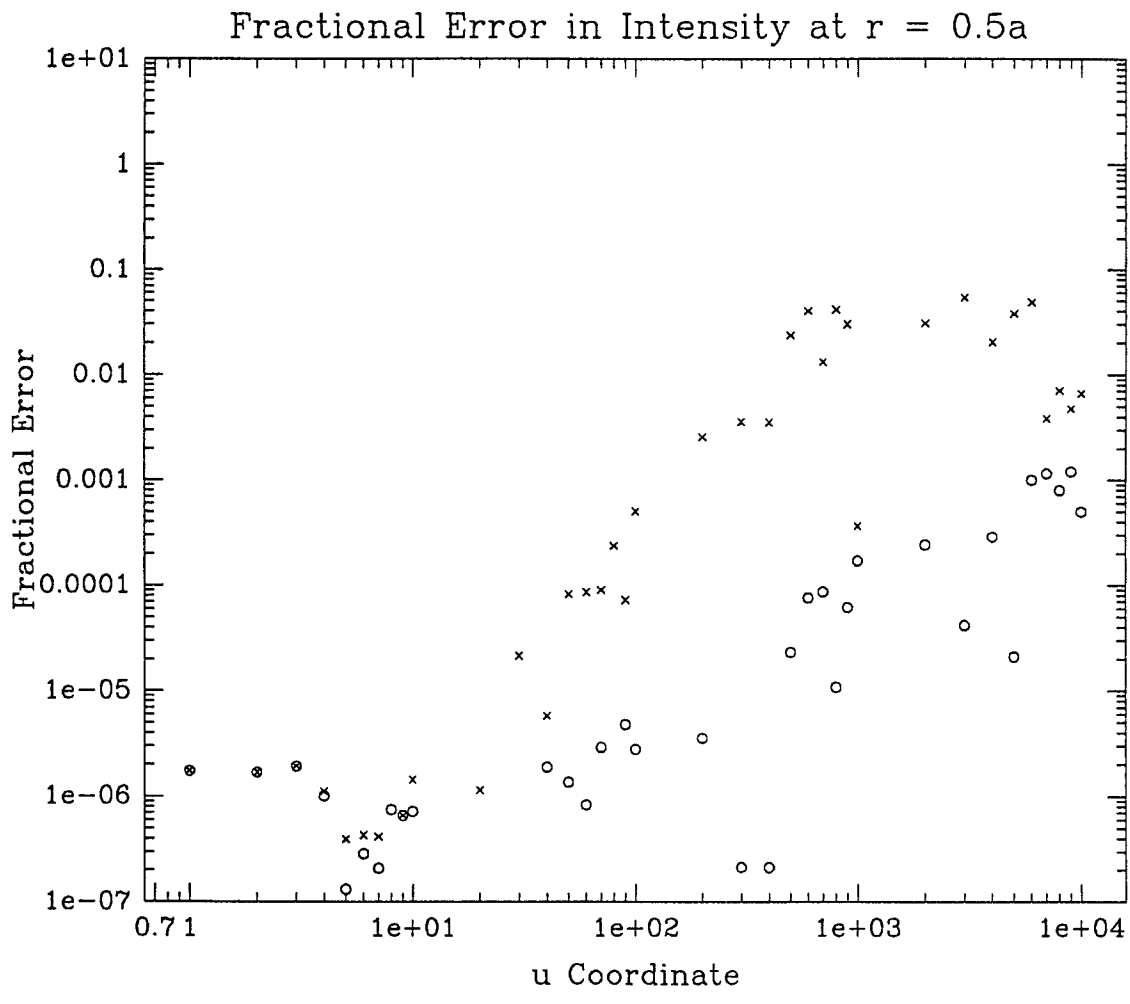
$$q = \sqrt{\sqrt{z^2 + a^2(1 - \rho)^2} - z}$$

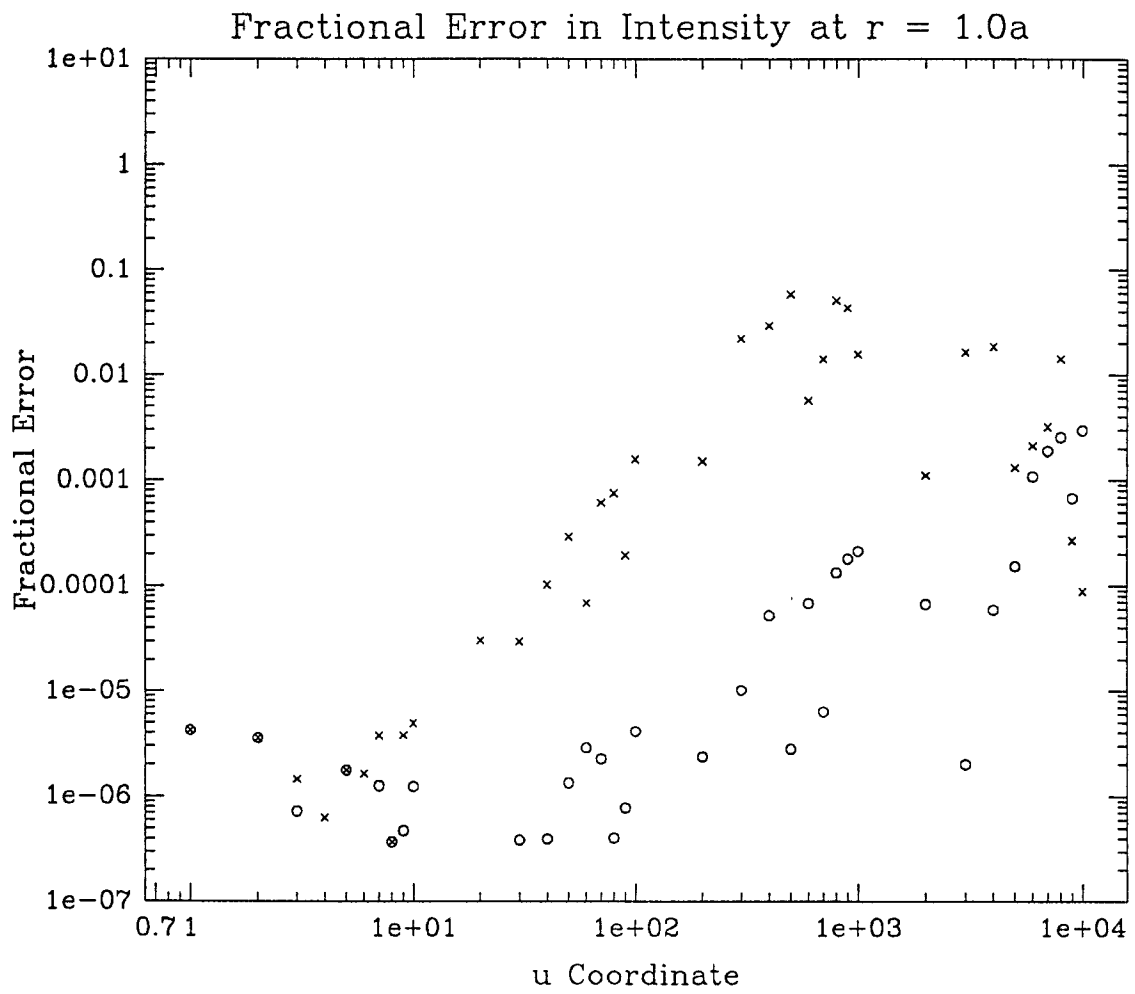


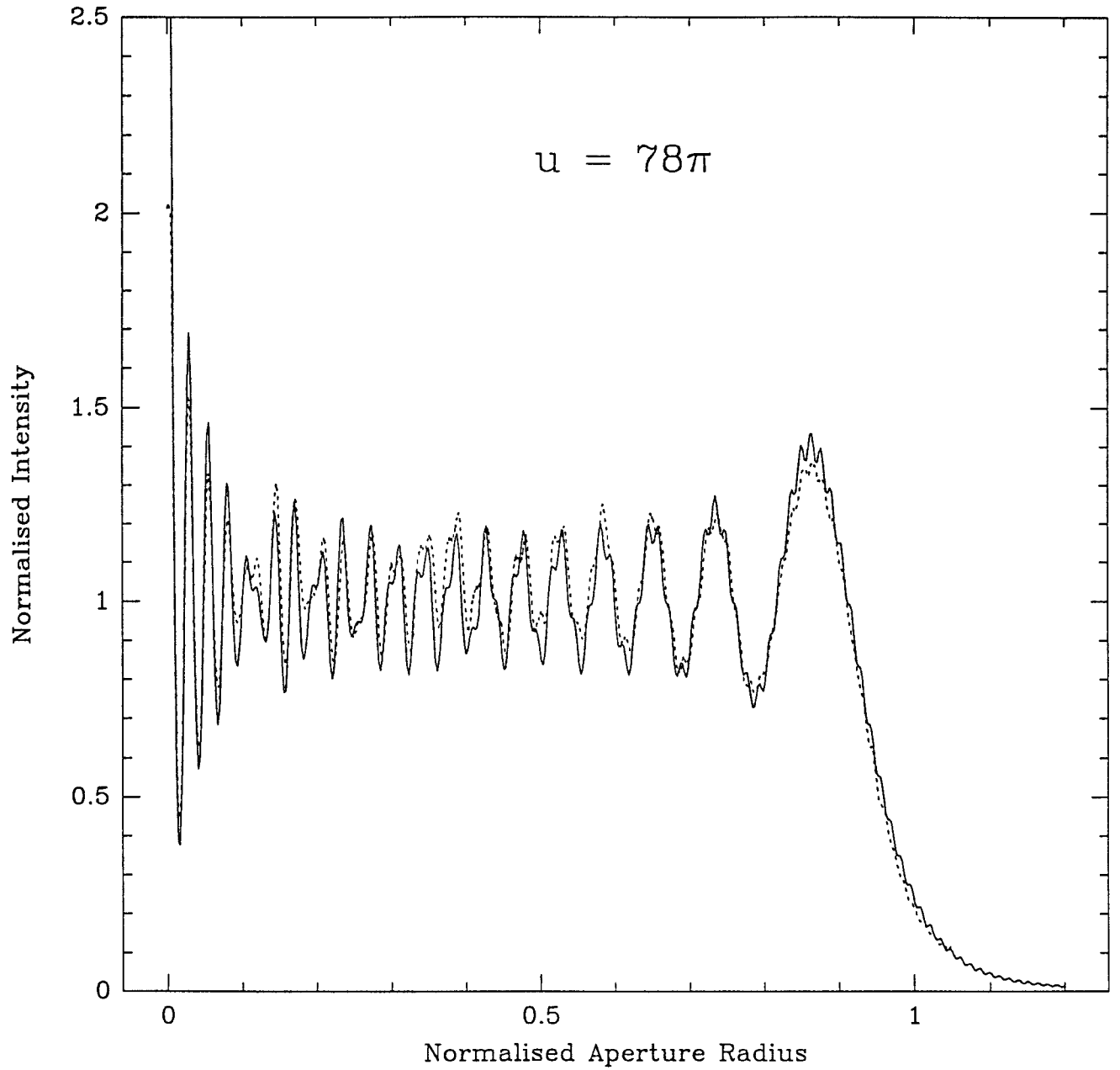
$$U \approx u \left(1 - \frac{a^2 (1 + 3\rho^2 + \rho^4)}{4z^2} + \dots \right)$$

$$V \approx v \left(1 - \frac{a^2 (1 + \rho^2)}{2z^2} + \dots \right)$$

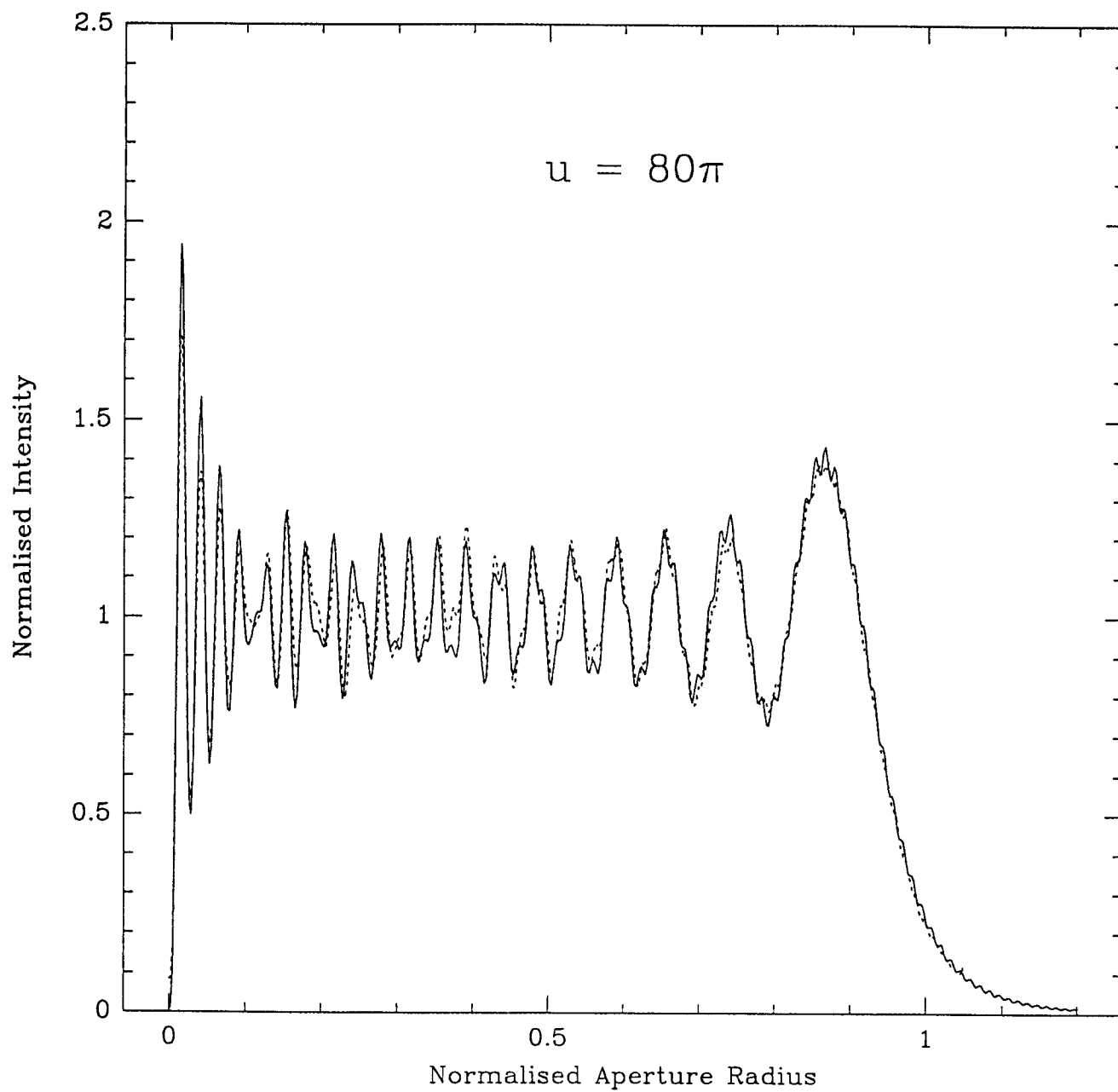




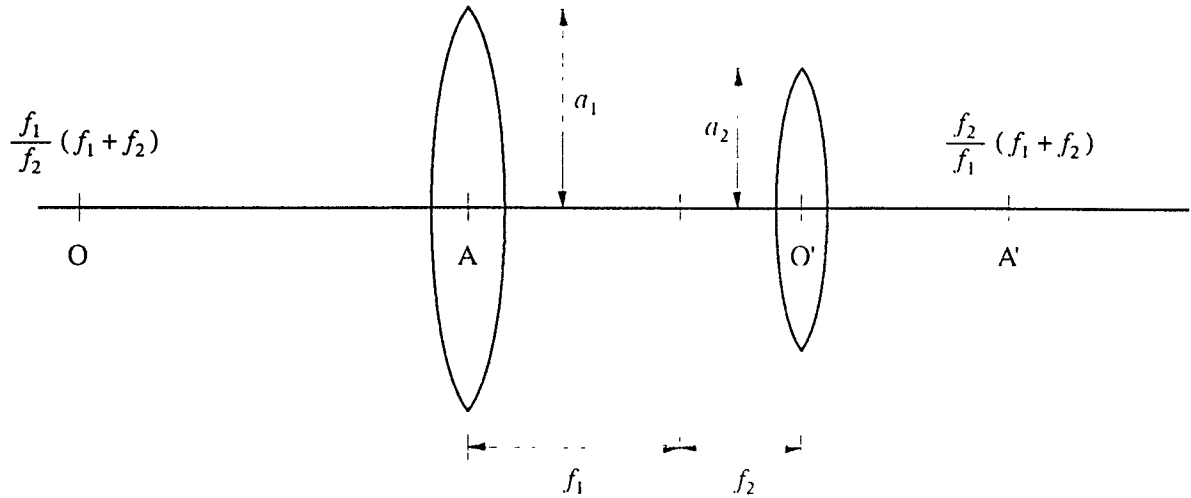




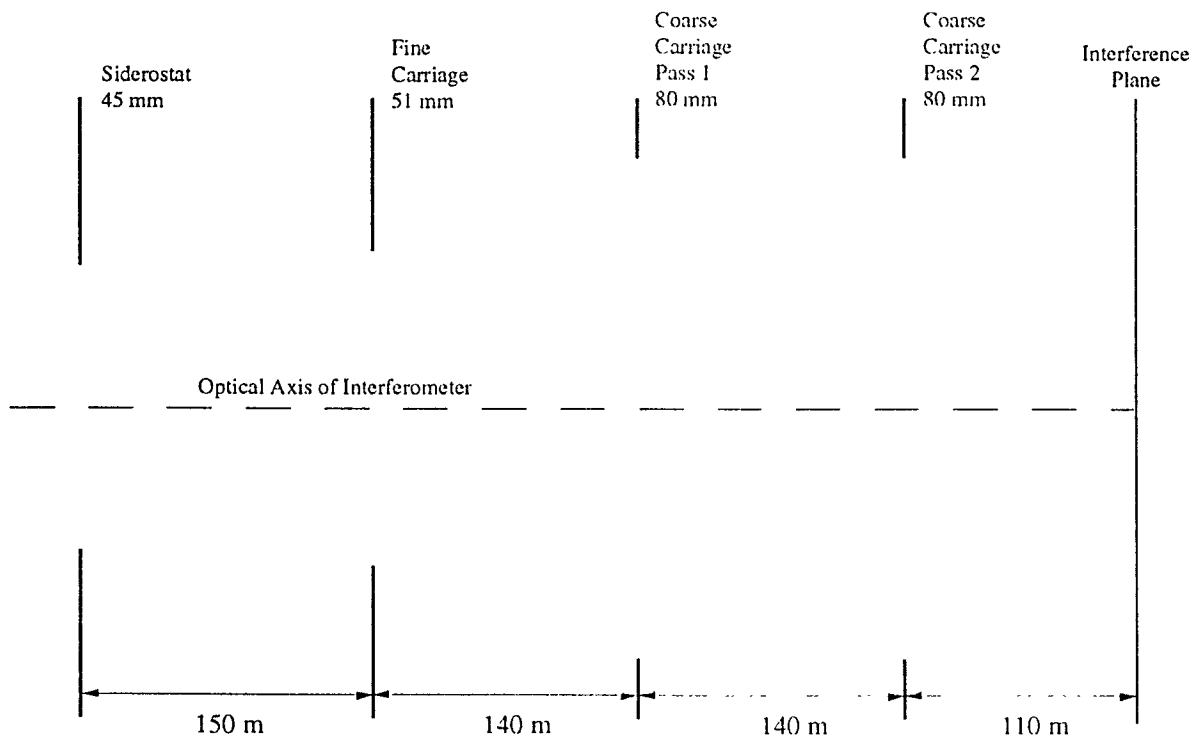
Measurement at $u = 78\pi$. Best fit was found at $u = 244.8509 = 77.938462\pi$

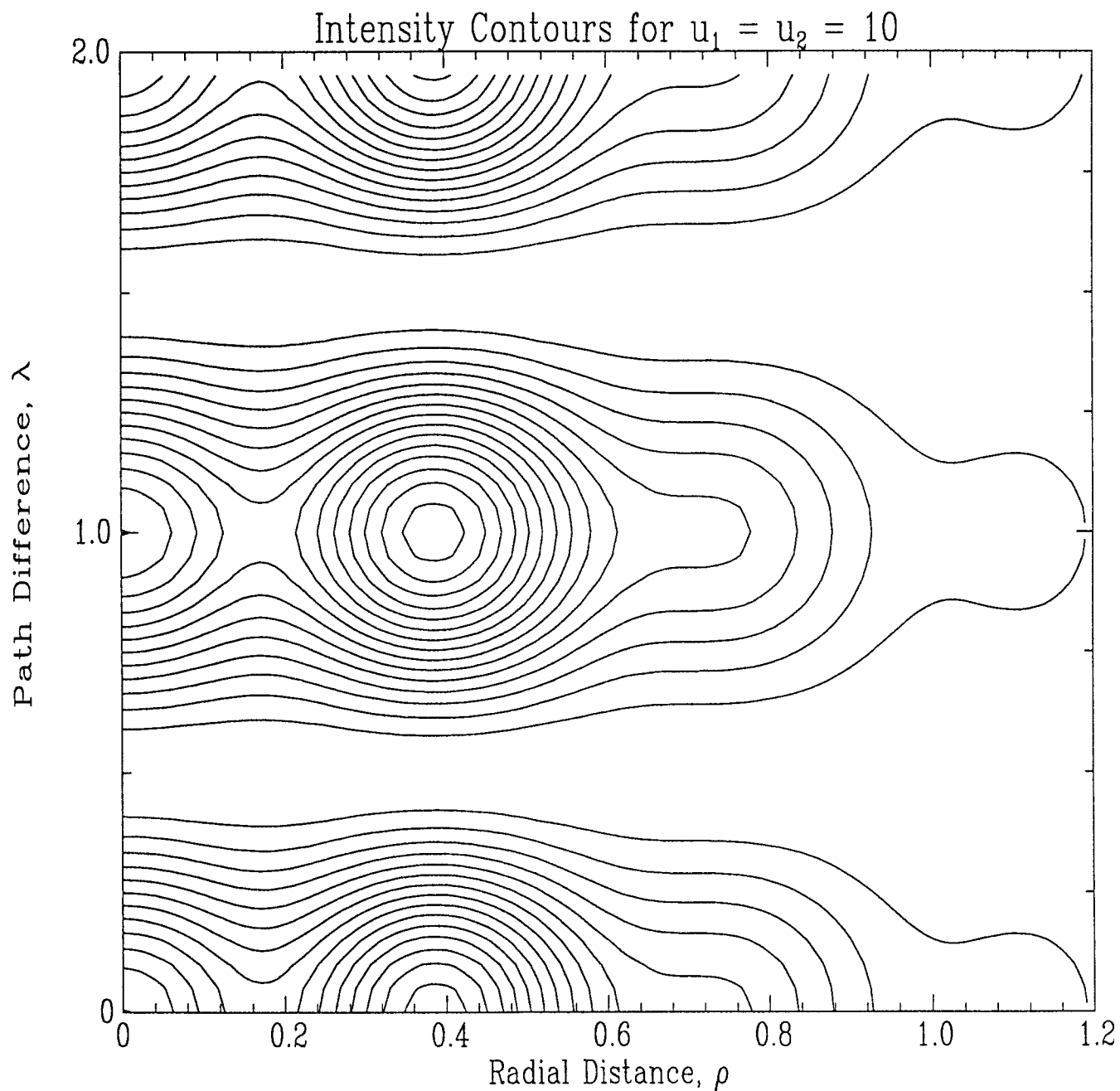


Measurement at $u = 80\pi$. Best fit was found at $u = 251.1794 = 79.952886\pi$

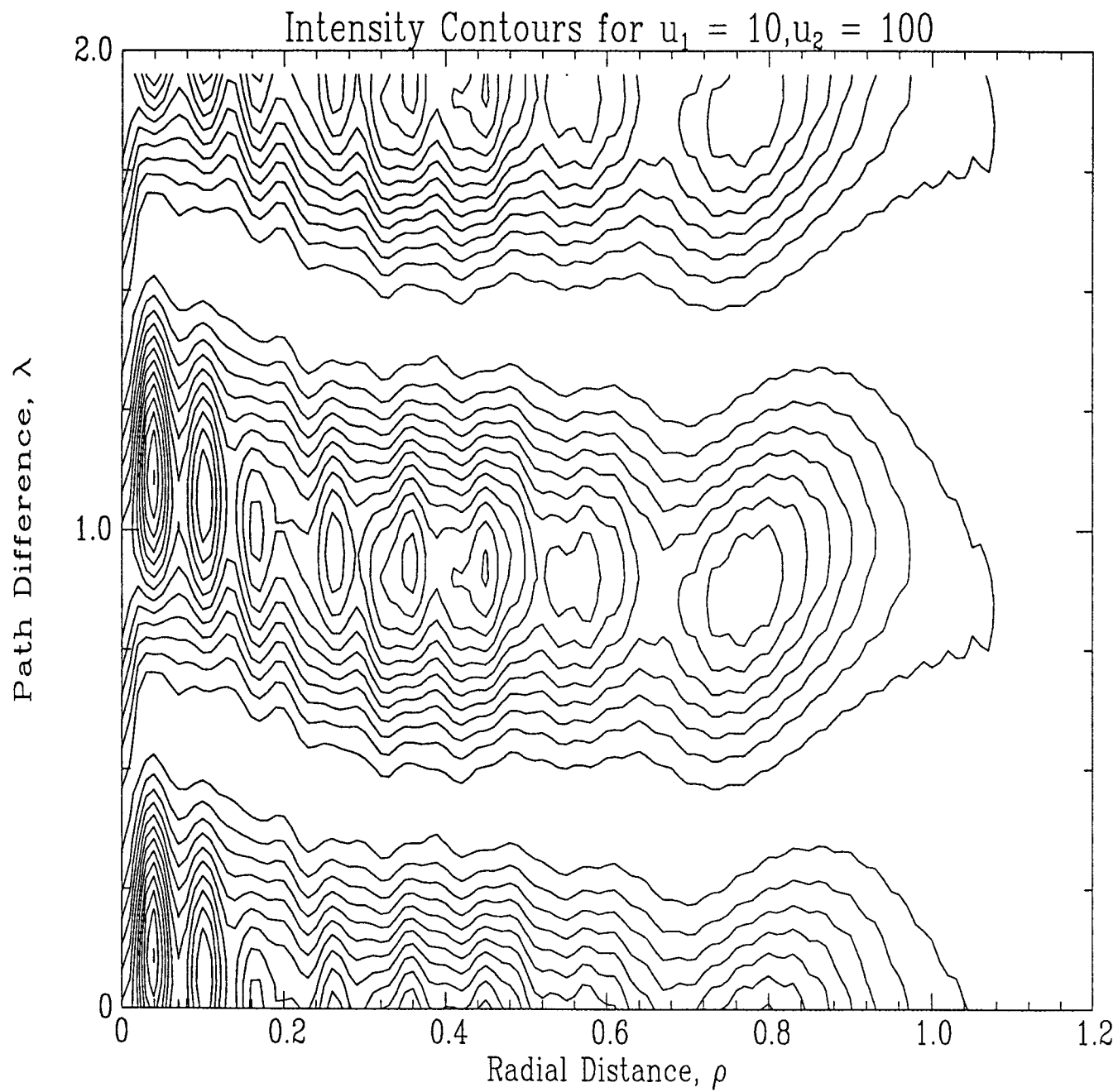


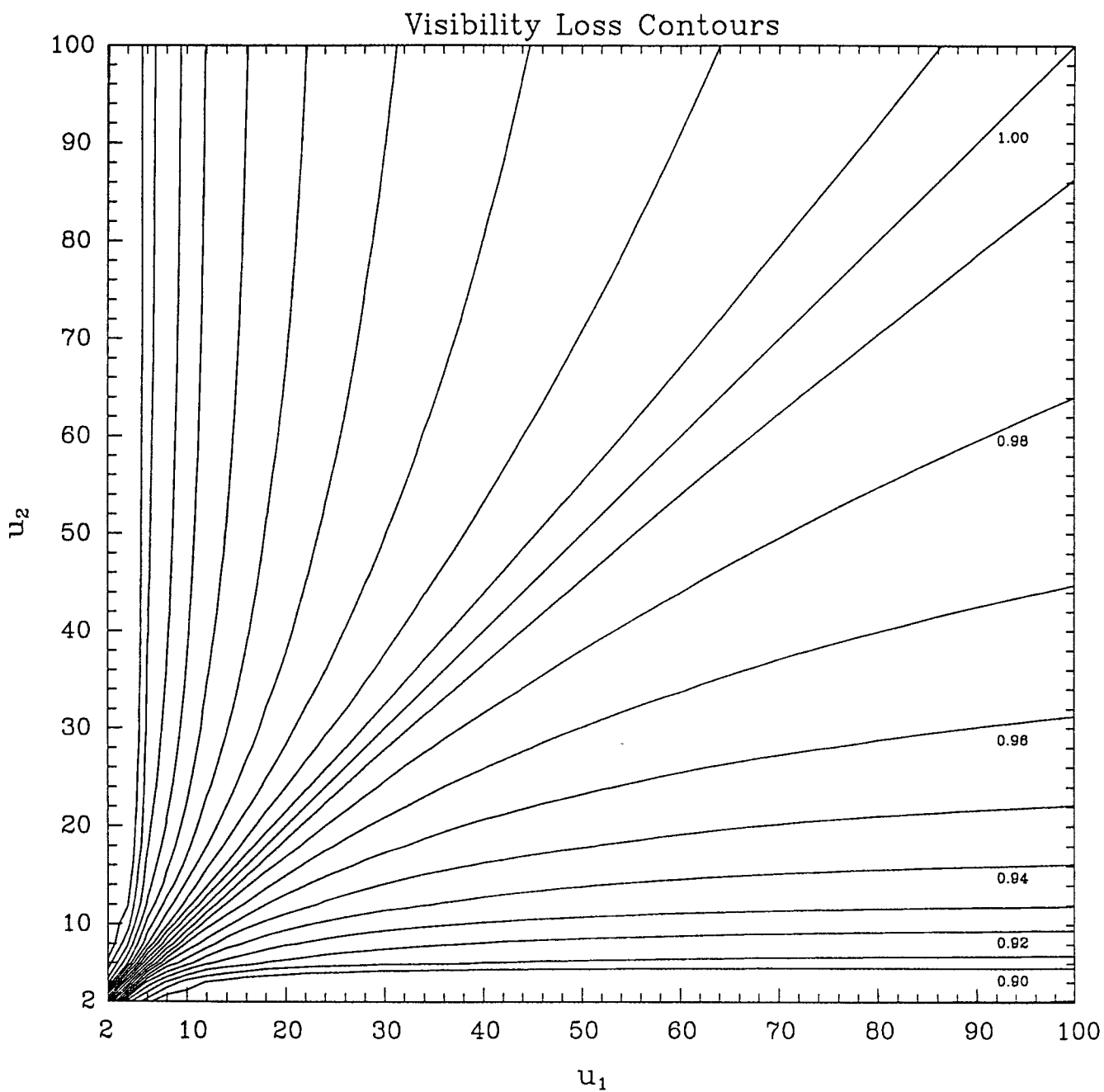
$$z_2 = \left(\frac{f_2}{f_1}\right)^2 z_1 - \frac{f_2}{f_1}(f_1 + f_2)$$



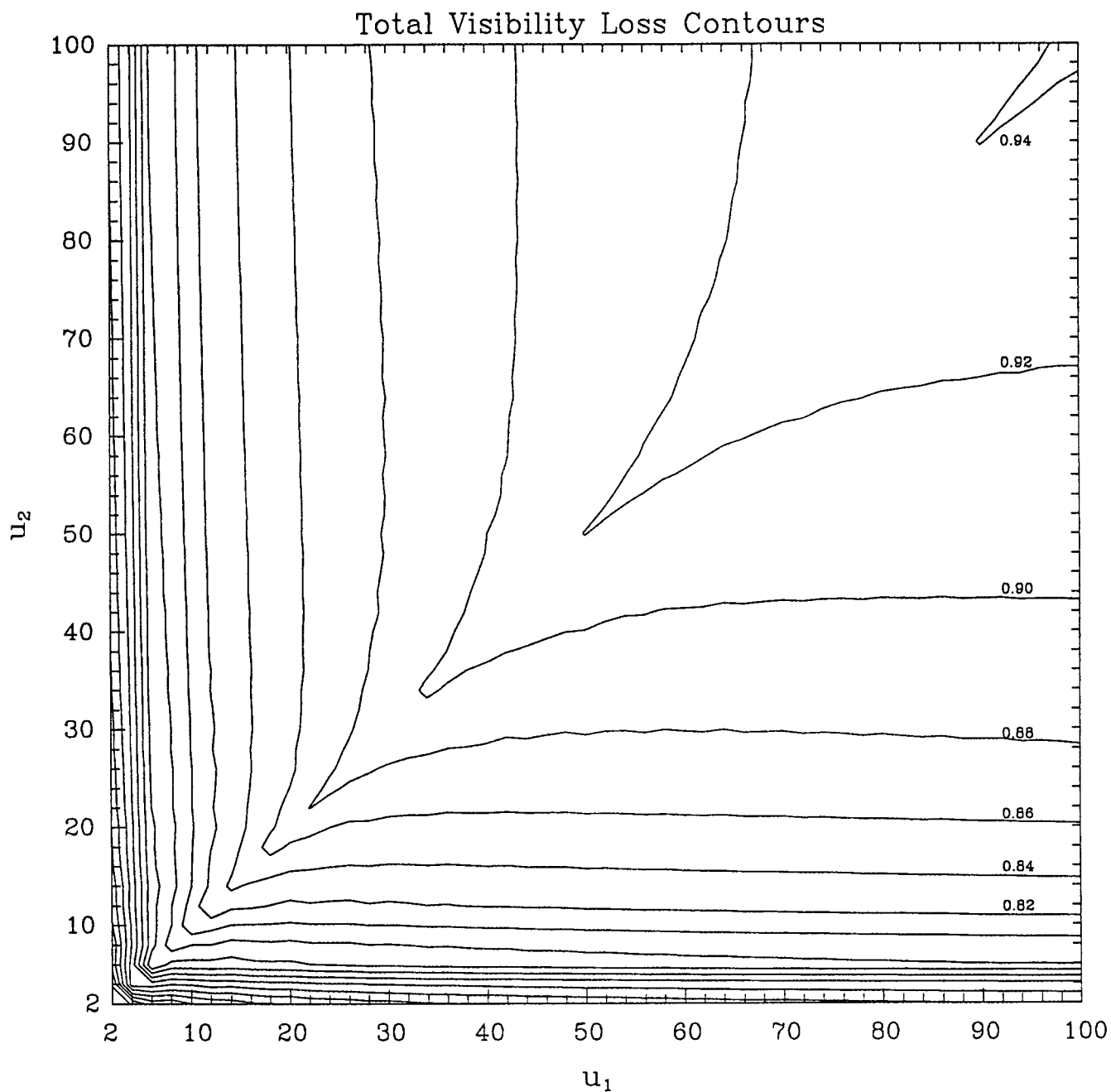


$$I(u_1, u_2, \rho, \delta) = I(u_1, \rho) + I(u_2, \rho) + 2\sqrt{I(u_1, \rho)I(u_2, \rho)} \cos\left(\phi(u_1, \rho) - \phi(u_2, \rho) + \frac{2\pi}{\lambda}\delta\right)$$





$$\nu = \frac{2 \int_0^{\rho_0} \sqrt{I(u_1, \rho) I(u_2, \rho)} \cos(\phi(u_1, \rho) - \phi(u_2, \rho) + \frac{2\pi}{\lambda} \delta_{max}) \rho d\rho}{\int_0^{\rho_0} (I(u_1, \rho) + I(u_2, \rho)) \rho d\rho}$$



$$\frac{\int_0^{\rho_0} (I(u_1, \rho) + I(u_2, \rho)) \rho d\rho}{\int_0^{\infty} (I(u_1, \rho) + I(u_2, \rho)) \rho d\rho}$$

$$\mathcal{V} = 2 \int_0^{\rho_0} \sqrt{I(u_1, \rho) I(u_2, \rho)} \cos(\phi(u_1, \rho) - \phi(u_2, \rho) + \frac{2\pi}{\lambda} \delta_{max}) \rho d\rho$$

Preferred Calculation Scheme

$$U(P) = \frac{k}{\lambda} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} A(\alpha, \beta, z) e^{-ik(\alpha x + \beta y)} d\alpha d\beta$$

$$\alpha = \sin \theta$$

$$\beta = \cos \psi$$

θ and ψ are angles with aperture normal

$$A(\alpha, \beta, z) = A_0(\alpha, \beta) e^{ikz\sqrt{1-(\alpha^2+\beta^2)}}$$

$$A_0(\alpha, \beta) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} U(x, y, 0) e^{ik(\alpha x + \beta y)} dx dy$$

