

Introduction to radiation pressure noise squeezing and optomechanical interaction

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Introduction: Why Optomechanics is Amazing

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Outline

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- **Optomechanics**
 - a simple model for GW detector
- **Standard Quantum Limit (SQL)**
- **Approaches for surpassing SQL**
 - Building up quantum correlation
- **Experimental evidence for quantum radiation pressure**

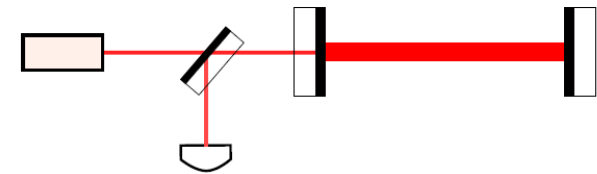
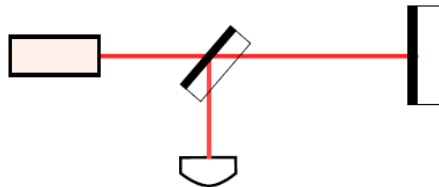
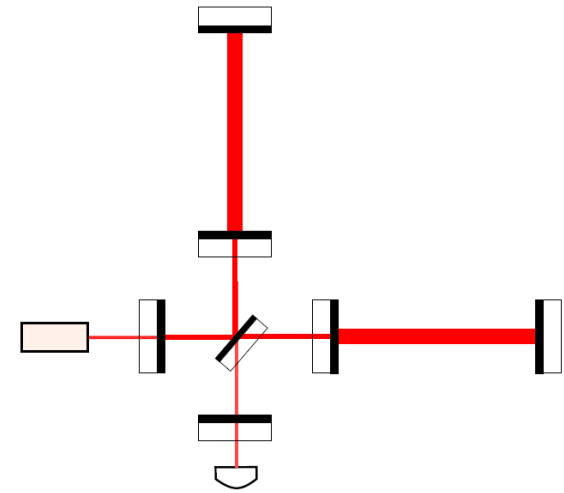
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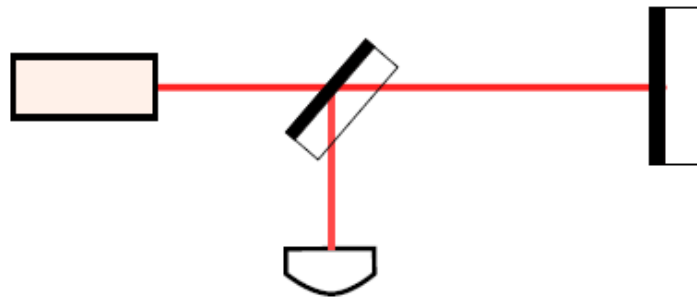
A simple² model for GW detector

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Optomechanics

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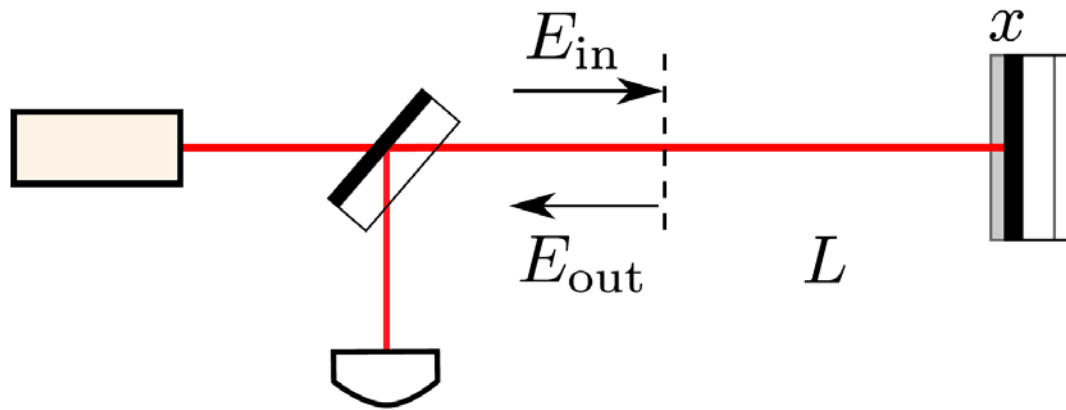
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“White-board” section

Deriving SQL



$$E_{\text{out}} = E_{\text{in}}\left(t - \frac{2L}{c} - \frac{2x}{c}\right)$$

$$m \ddot{x}(t) = \frac{\hbar\omega_0}{c} E_{\text{in}}^2(t) + \frac{m}{2} L \ddot{h}(t)$$

Radiation pressure

GW tidal force

$$(a_1 \ll \bar{E}_0, \quad a_2 \ll \bar{E}_0)$$

$$E_{\text{in}}(t) = [\bar{E}_0 + a_1(t)] \cos \omega_0 t + a_2(t) \sin \omega_0 t$$

$$\approx \bar{E}_0 \left[1 + \frac{a_1(t)}{\bar{E}_0} \right] \cos \left[\omega_0 t - \frac{a_2(t)}{\bar{E}_0} \right]$$

Amplitude Modulation

Phase Modulation

$$E_{\text{out}}(t) = [\bar{E}'_0 + b_1(t)] \cos \omega_0 t + b_2(t) \sin \omega_0 t$$

Linearization: $x \ll \lambda$

Input-output relation:

$$b_1(t) = a_1(t)$$

$$b_2(t) = a_2(t) + \frac{\alpha}{\hbar} x(t) \quad (\alpha \equiv \hbar\omega_0 \bar{E}_0/c)$$

$$m\ddot{x}(t) = \alpha a_1(t) + \frac{1}{2}mL\ddot{h}(t)$$

Radiation pressure

GW tidal force

“One-step” quantization:

Input-output relation:

$$\hat{b}_1(t) = \hat{a}_1(t)$$

$$\hat{b}_2(t) = \hat{a}_2(t) + \frac{\alpha}{\hbar} \hat{x}(t) \quad (\alpha \equiv \hbar\omega_0 \bar{E}_0/c)$$

$$m\ddot{\hat{x}}(t) = \alpha \hat{a}_1(t) + \frac{1}{2}mL\ddot{\hat{h}}(t)$$

Radiation pressure

GW tidal force

Frequency domain:

$$\hat{b}_1(\Omega) = \hat{a}_1(\Omega)$$

$$\hat{b}_2(\Omega) = \hat{a}_2(\Omega) - \kappa \hat{a}_1(\Omega) + \sqrt{2\kappa} \frac{h(\Omega)}{h_{\text{SQL}}}$$

$$\kappa = \frac{\alpha^2}{\hbar m \Omega^2} \quad h_{\text{SQL}} = \sqrt{\frac{2\hbar}{m \Omega^2 L^2}}$$

Measuring phase quadrature:

$$\hat{b}_2(\Omega) = \hat{a}_2(\Omega) - \kappa \hat{a}_1(\Omega) + \sqrt{2\kappa} \frac{h(\Omega)}{h_{\text{SQL}}}$$

Shot noise Radiation-pressure noise GW signal

Noise spectrum (for vacuum):

$$S_{11}(\Omega) = S_{22}(\Omega) = 1, \quad S_{12}(\Omega) = 0$$

Measuring phase quadrature:

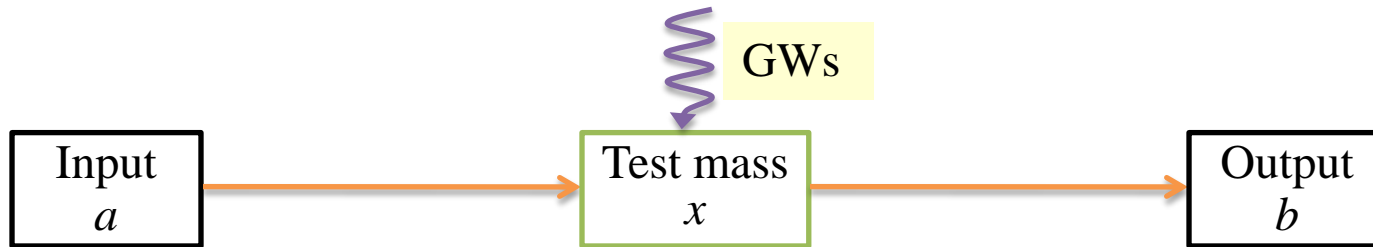
$$\hat{b}_2(\Omega) = \hat{a}_2(\Omega) - \kappa \hat{a}_1(\Omega) + \sqrt{2\kappa} \frac{h(\Omega)}{h_{\text{SQL}}}$$

h -referred noise spectrum:

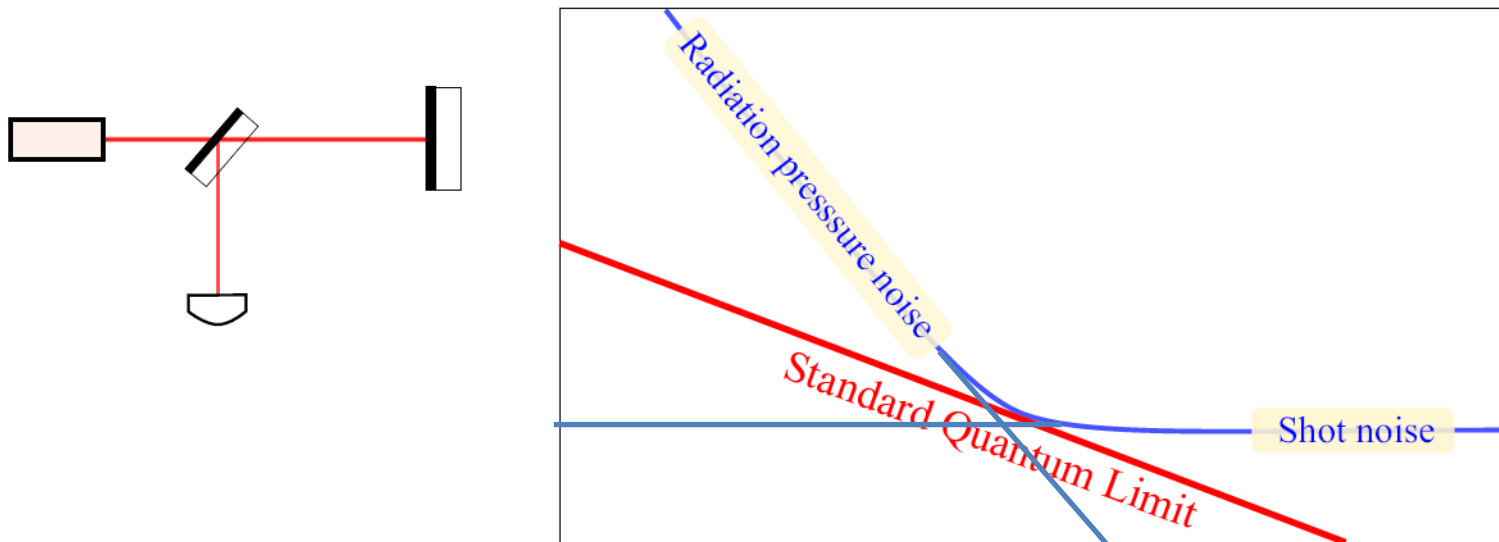
$$\begin{aligned} S_h(\Omega) &= \left(\frac{1}{\kappa} + \kappa \right) \frac{h_{\text{SQL}}^2}{2} \\ &\geq h_{\text{SQL}}^2 = \frac{2\hbar}{m\Omega^2 L^2} \end{aligned}$$

Standard Quantum Limit (SQL)

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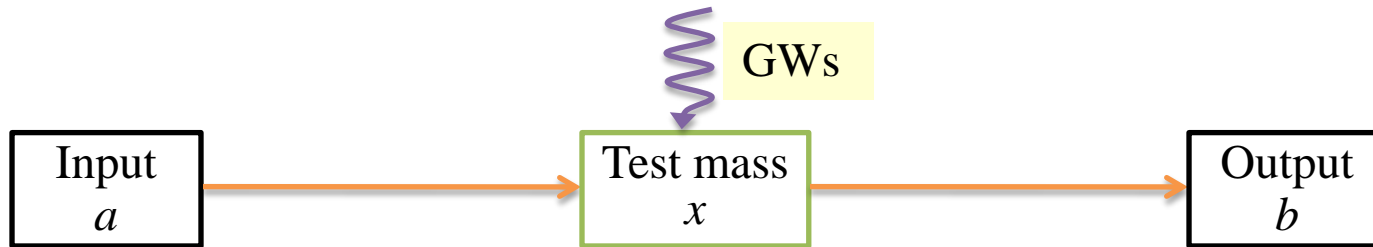


Noise spectrum: $S_h(\Omega) = \left[\frac{1}{\kappa} + \kappa \right] \frac{h_{\text{SQL}}^2}{2} \geq h_{\text{SQL}}^2 = \frac{2\hbar}{m\Omega^2 L^2}$

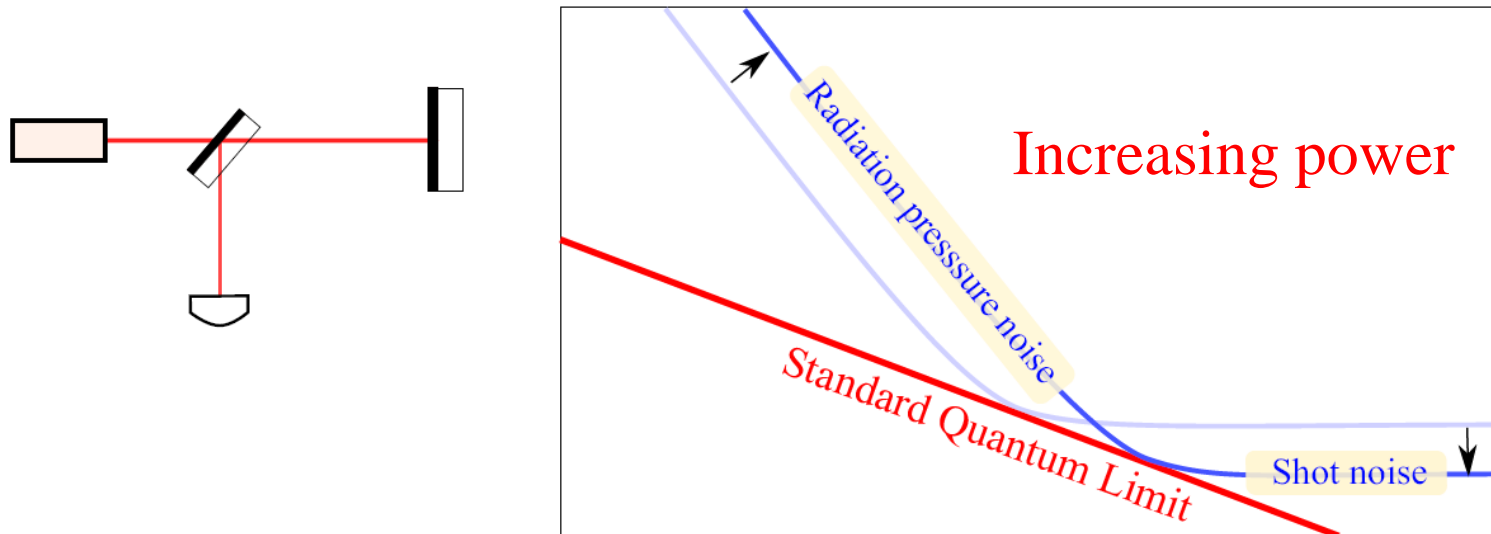


Standard Quantum Limit (SQL)

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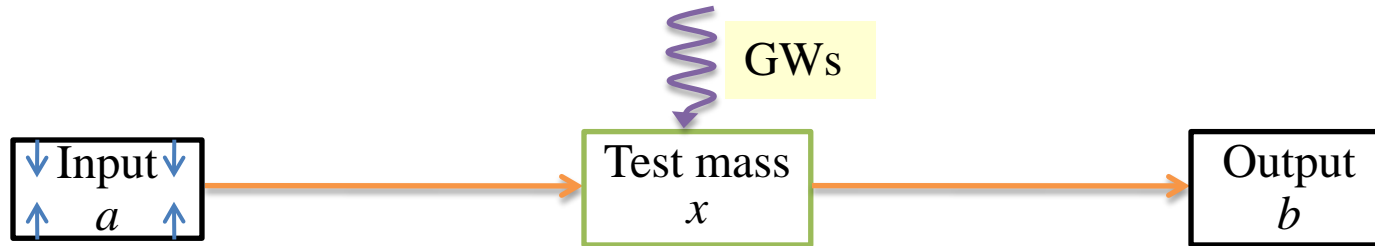


Noise spectrum:
$$S_h(\Omega) = \left[\frac{1}{\kappa} + \kappa \right] \frac{h_{\text{SQL}}^2}{2} \geq h_{\text{SQL}}^2 = \frac{2\hbar}{m\Omega^2 L^2}$$

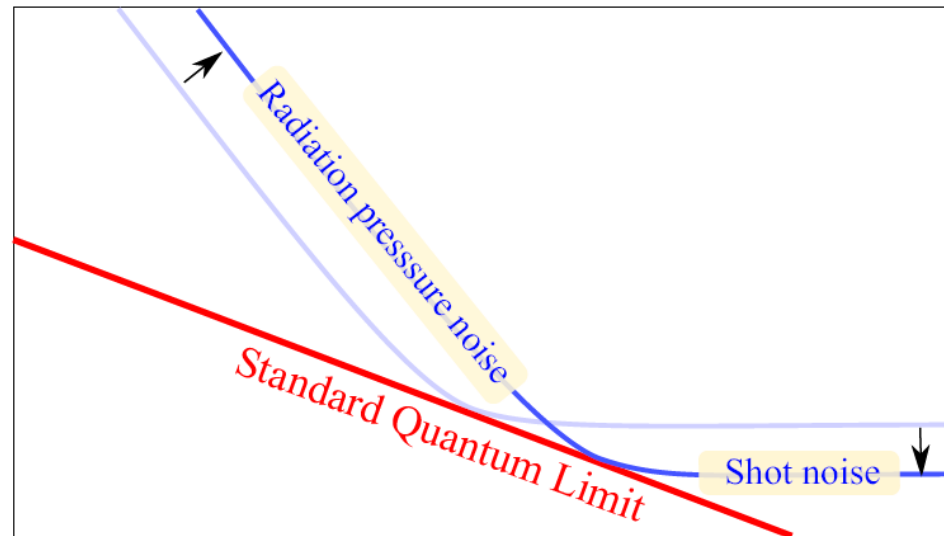
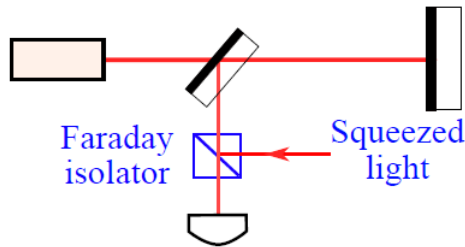


Standard Quantum Limit (SQL)

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Squeezed light input:



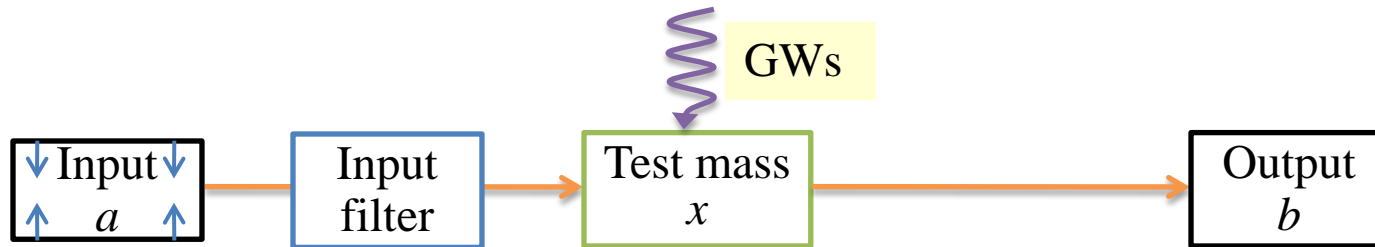
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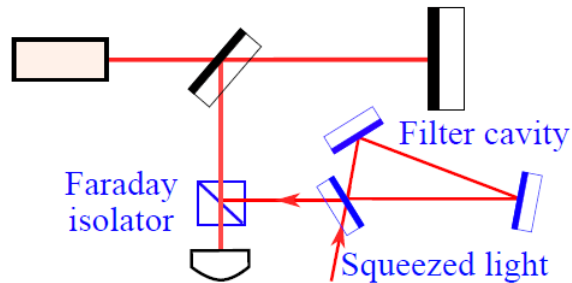
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Surpassing SQL: Freq. dep. squeezing

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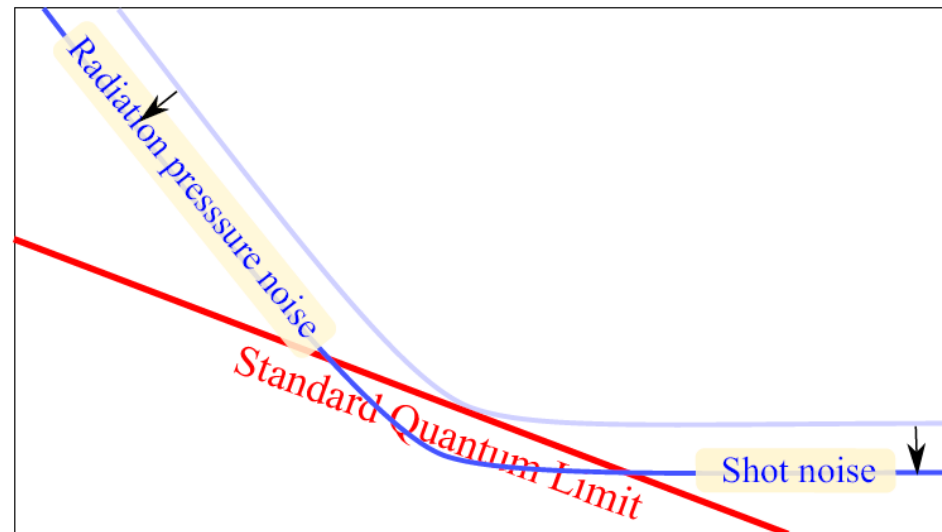


Frequency-dependent squeezing: input filtering



Filter cavity:

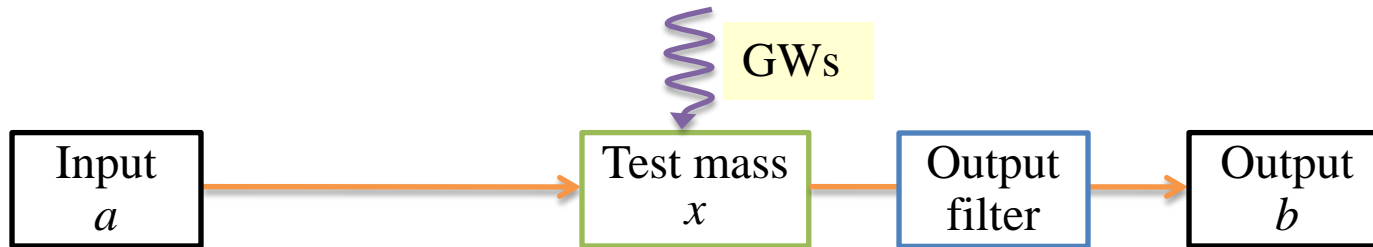
1. Passive optical cavity
2. Optomechanical cavity



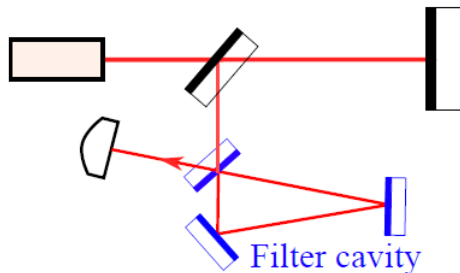
Reference: H. J. Kimble *et al.*, Phys. Rev. D **65**, 022002 (2001).

Surpassing SQL: Variational readout

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Variational readout: Output filtering

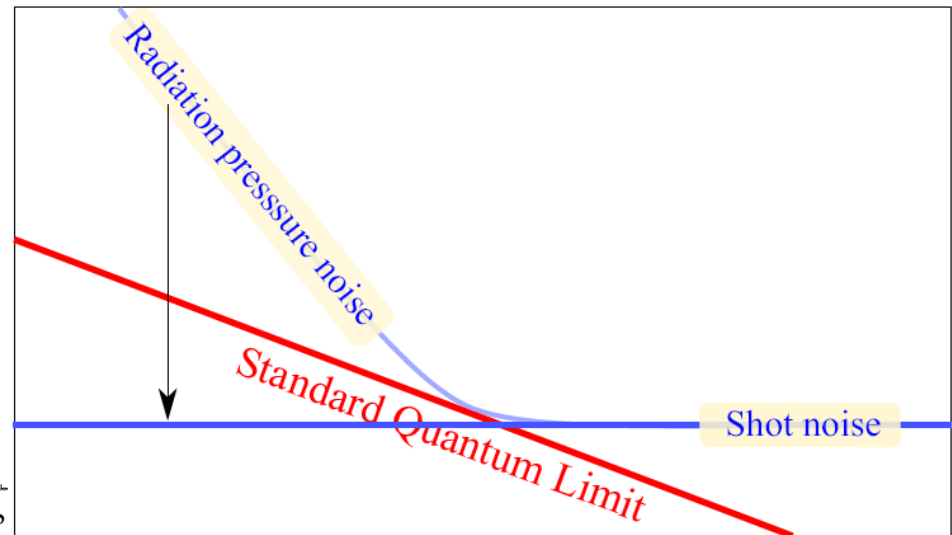


$$b_1 = a_1$$

$$b_2 = a_2 - \kappa a_1 + \text{signal}$$

Measure: $b_\zeta = b_1 \sin \zeta + b_2 \cos \zeta$

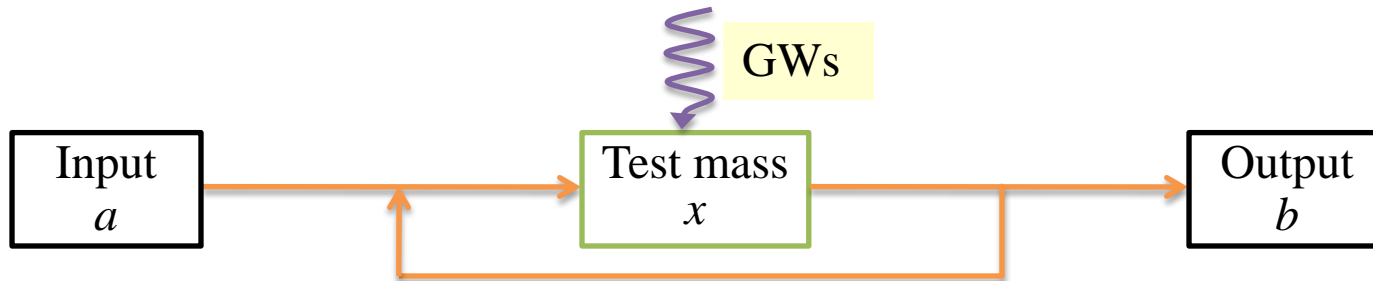
$\tan \zeta = \kappa$ \Rightarrow Rad. noise cancelled



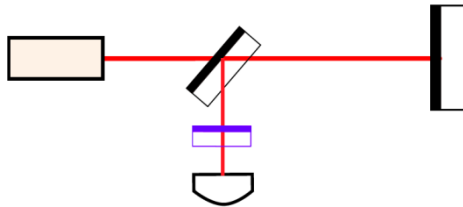
Reference: H. J. Kimble *et al.*, Phys. Rev. D **65**, 022002 (2001).

Surpassing SQL: Coherent feedback

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Coherent feedback:

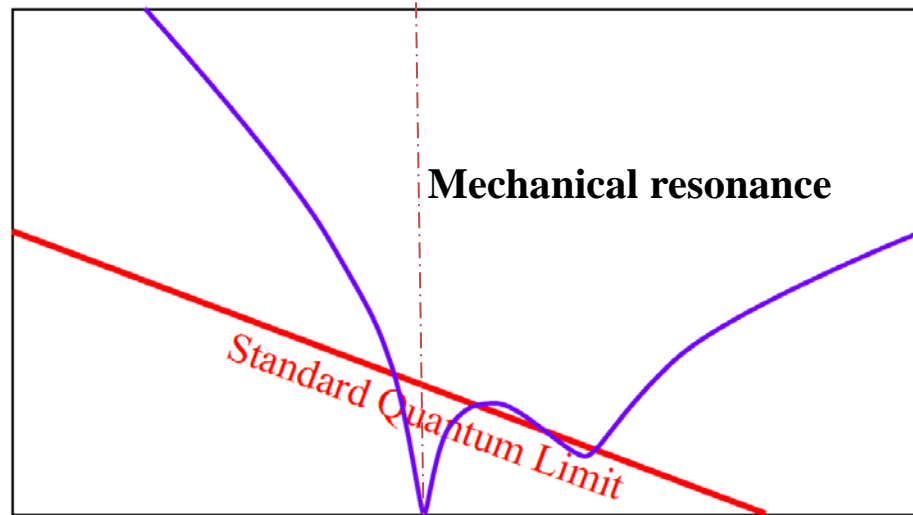


1. Modifying test-mass dynamics:

Optical-spring effect

2. Quantum correlation

between
Shot and rad. pre. noise



Reference: A. Buonanno and Y. Chen, Phys. Rev. D **65**, 042001 (2002).

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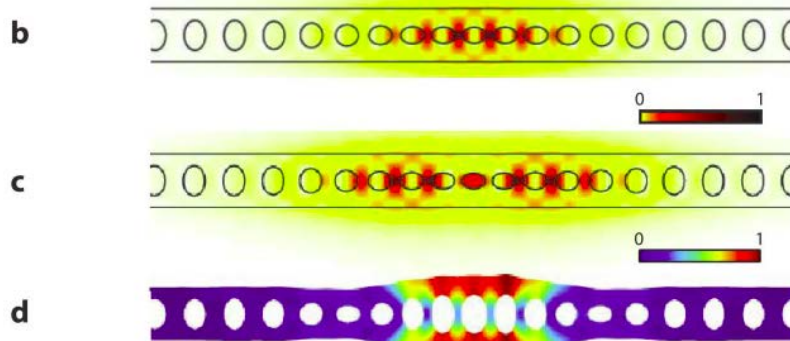
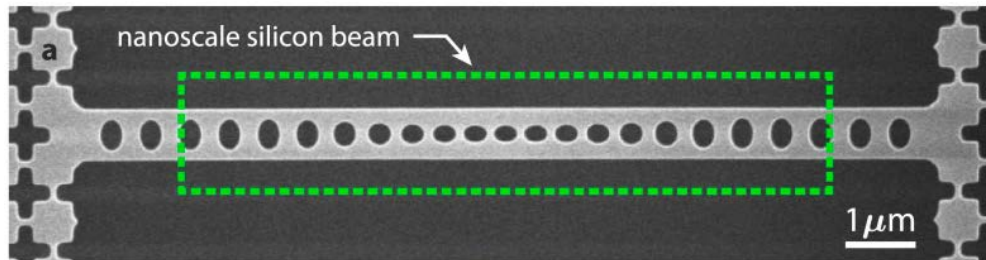
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Experimental evidence finally comes

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Setup:



Cavity parameters

$$\omega_c/2\pi = 205.3 \text{ THz}$$

$$\kappa_c/2\pi = 390 \text{ MHz}$$

$$g_c/2\pi = 960 \text{ KHz}$$

$$\omega_r/2\pi = 194.1 \text{ THz}$$

$$\kappa_r/2\pi = 1 \text{ GHz}$$

$$g_r/2\pi = 430 \text{ KHz}$$

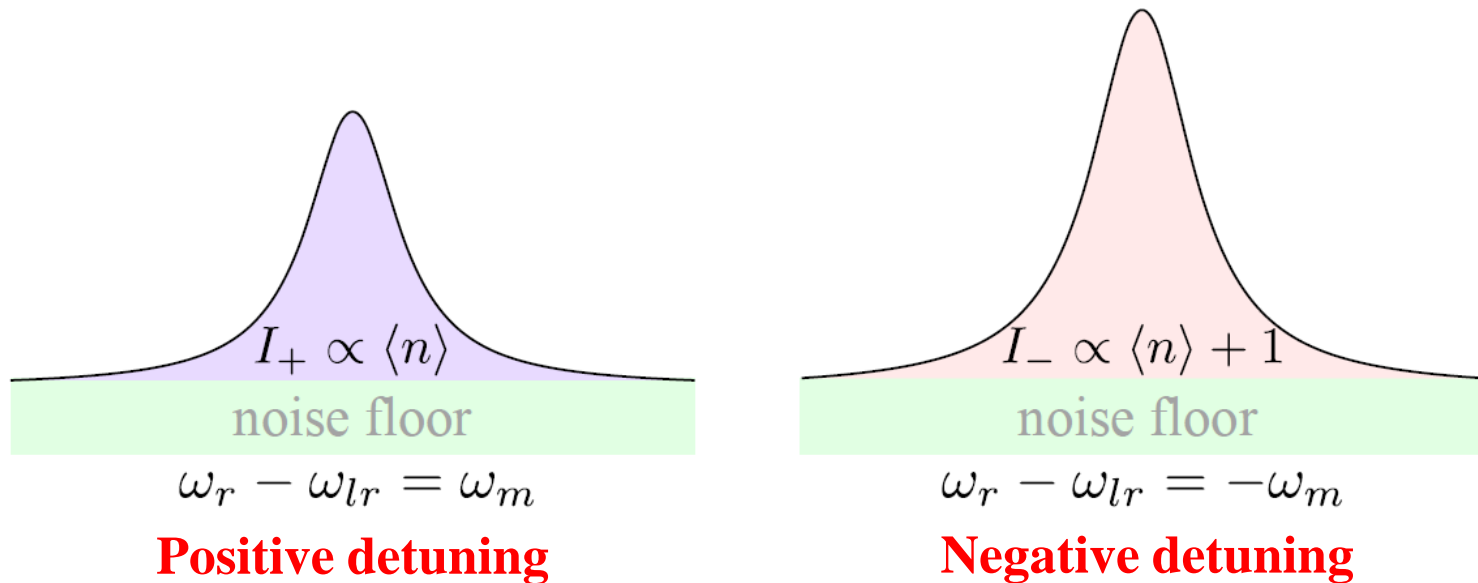
$$\omega_m/2\pi = 3.99 \text{ GHz}$$

- Reference:** [1] A. Safavi-Naeini *et al.*, Phys. Rev. Lett. **108**, 033602 (2012).
[2] F. Khalili *et al.*, In preparation (2012)

Experimental evidence finally comes

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Asymmetric Spectra:



Explanation:

$$\hat{y}_{\text{out}} = \alpha^{-1} \hat{z} + \hat{x} = \alpha^{-1} \hat{z} + \chi[\alpha \hat{F}_{\text{rad}} + F_{\text{th}}]$$

Experimental evidence finally comes

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Output spectrum:

$$S_{yy} = \alpha^{-2} S_{zz} + 2\Re\{\chi^* S_{zF}\} + |\chi|^2 (\alpha^2 S_{FF}^{\text{rad}} + S_{FF}^{\text{th}})$$



Shot noise

Correlation

Rad. Pres. noise

$$S_{zF} \approx -i\hbar \frac{\omega_m}{\Delta} \Big|_{\Delta=\pm\omega_m} = \mp i\hbar$$



Correlation at the quantum level

Experimental evidence finally comes

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Output spectrum:

$$S_{yy} = \alpha^{-2} S_{zz} + 2\Re\{\chi^* S_{zF}\} + |\chi|^2 (\alpha^2 S_{FF}^{\text{rad}} + S_{FF}^{\text{th}})$$

$$S_{zF} \approx -i\hbar \frac{\omega_m}{\Delta} \Big|_{\Delta=\pm\omega_m} = \mp i\hbar$$

Correlation at the quantum level

