



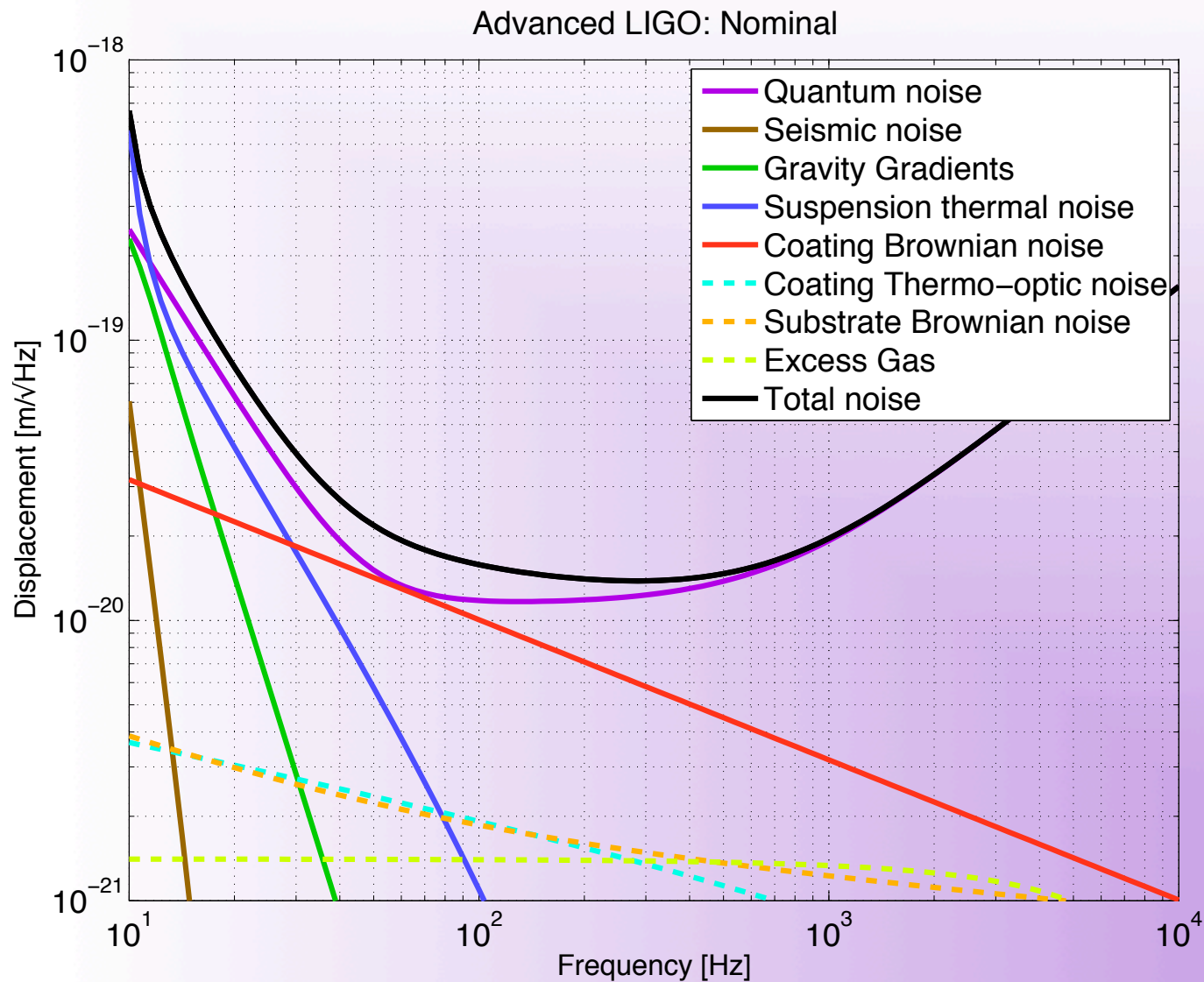
# Beyond Advanced Gravitational Wave Detectors: beating the quantum limit with squeezed states of light

Lisa Barsotti  
(LIGO-MIT)

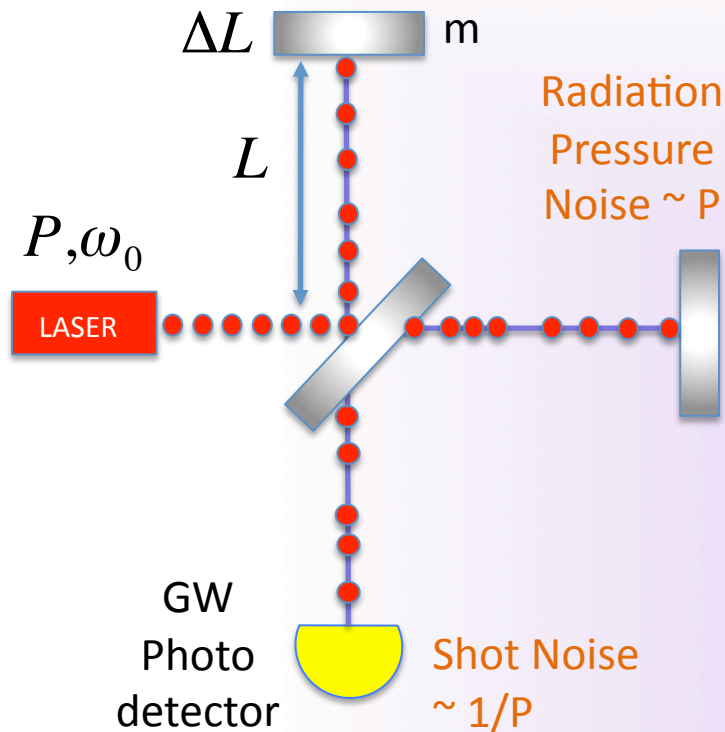


# Why we talk about quantum noise

$$h = \frac{\Delta L}{L}$$



# Where quantum noise comes from



- ✧ **SHOT NOISE**: Photon counting noise produced by fluctuations of the number of photon detected at the interferometer output  
 → Limitation of the precision you can measure arm displacement
- ✧ **RADIATION PRESSURE NOISE**: Back-action noise caused by random motion of the mirrors due to fluctuations of the number of photons impinging on the mirrors  
 → Additional displacement noise

$$\Delta L_{rad} = \frac{1}{cm\Omega^2} \sqrt{8\hbar P \omega_0} \quad \Delta L_{shot} = c \sqrt{\frac{\hbar}{2P\omega_0}}$$

$$\Delta L_{Quantum} = \sqrt{\Delta L_{rad}^2 + \Delta L_{shot}^2}$$

# “Standard Quantum Limit”

$$\Delta L_{\text{Quantum}} = \sqrt{\frac{4\hbar}{m\Omega^2}} \sqrt{\frac{1}{2} \left( \mathbf{K} + \frac{1}{\mathbf{K}} \right)}, \quad \mathbf{K} = \frac{4P\omega_0}{c^2 m\Omega^2}$$

It doesn't depend on the interferometer, just on the quantum mechanics of a harmonic oscillator mass

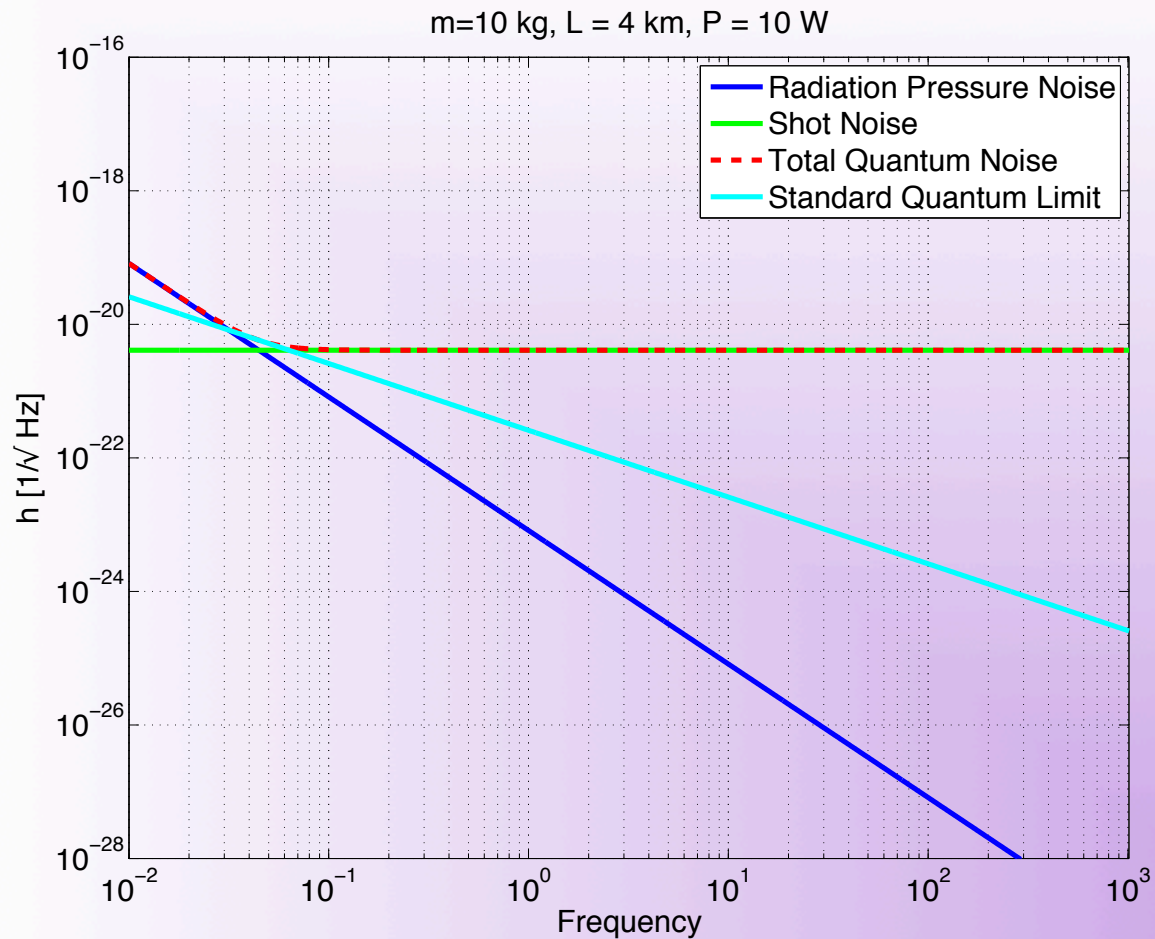
$$\Delta L_{\text{SQL}} \sim \sqrt{\frac{\hbar}{m\Omega^2}}$$

$$h_{\text{Quantum}} = \frac{\Delta L_{\text{Quantum}}}{L} = \sqrt{\frac{4\hbar}{m\Omega^2 L^2}} \sqrt{\frac{1}{2} \left( \mathbf{K} + \frac{1}{\mathbf{K}} \right)}$$

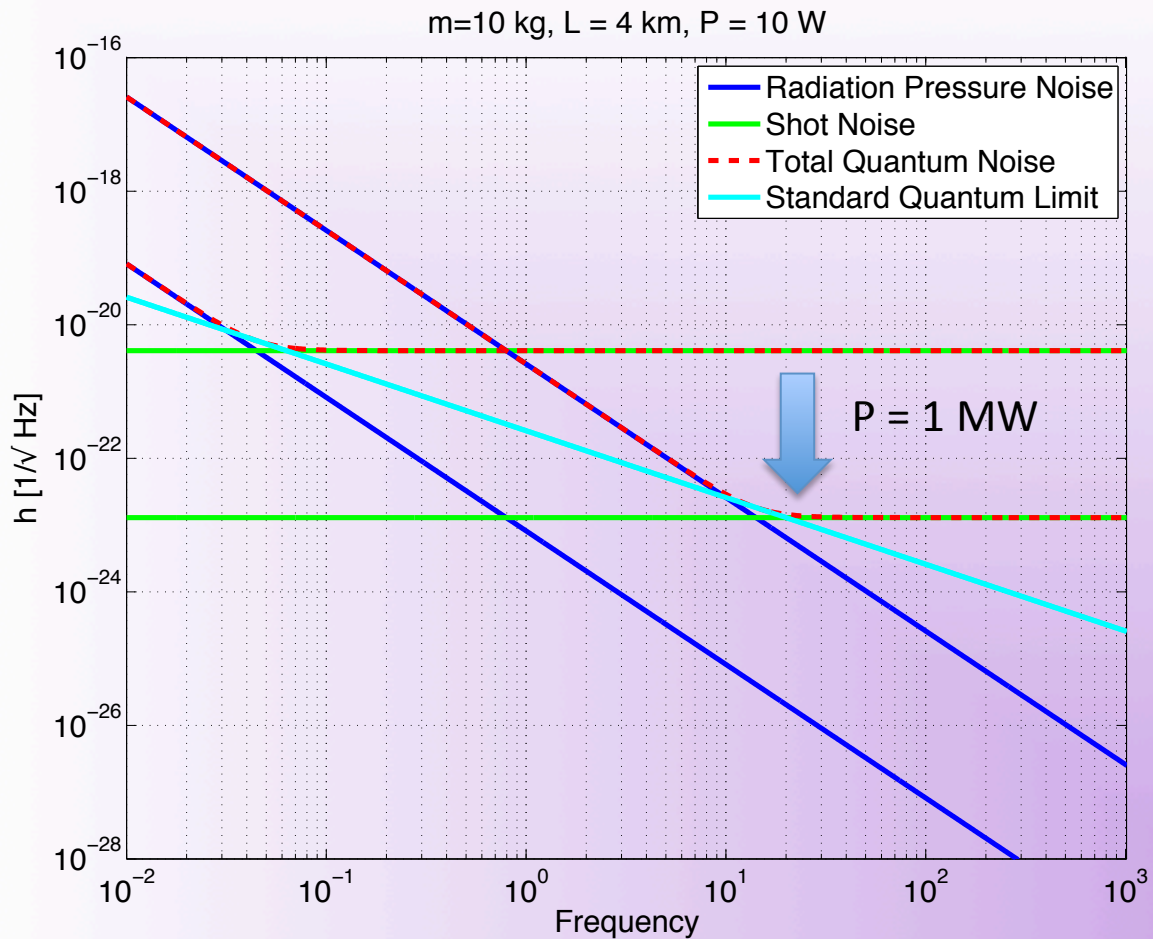
$$\mathbf{K} = 1 \Rightarrow P = \frac{c^2 m\Omega^2}{4\omega_0} \quad h_{\text{SQL}} = \sqrt{\frac{4\hbar}{m\Omega^2 L^2}}$$



# Simple Michelson, $P = 10 \text{ W}$



# Simple Michelson, $P = 1 \text{ MW}$



## “Easy” ways of reducing quantum noise

$$h_{\text{Quantum}} = \sqrt{\frac{1}{2}} \sqrt{\frac{4\hbar}{m\Omega^2 L^2}} \sqrt{\left(\mathbf{K} + \frac{1}{\mathbf{K}}\right)}, \quad \mathbf{K} = \frac{4P\omega_0}{c^2 m\Omega^2}$$

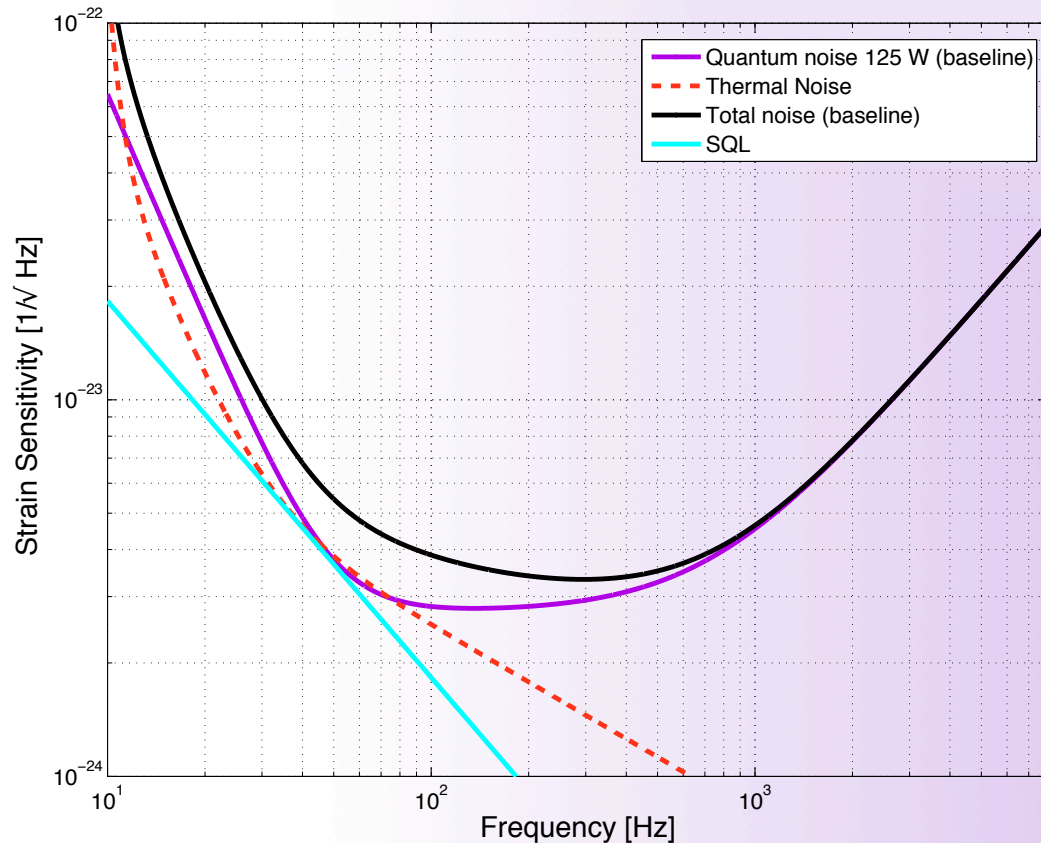
- ✧ Just make your interferometer longer
- ✧ More power to improve shot noise + heavier test masses to compensate for radiation pressure noise

$$h_{\text{Quantum}} \sim \sqrt{\left(\frac{4P\omega_0}{c^2 m^2 \Omega^2} + \frac{c^2 \Omega^2}{4P\omega_0}\right)}$$

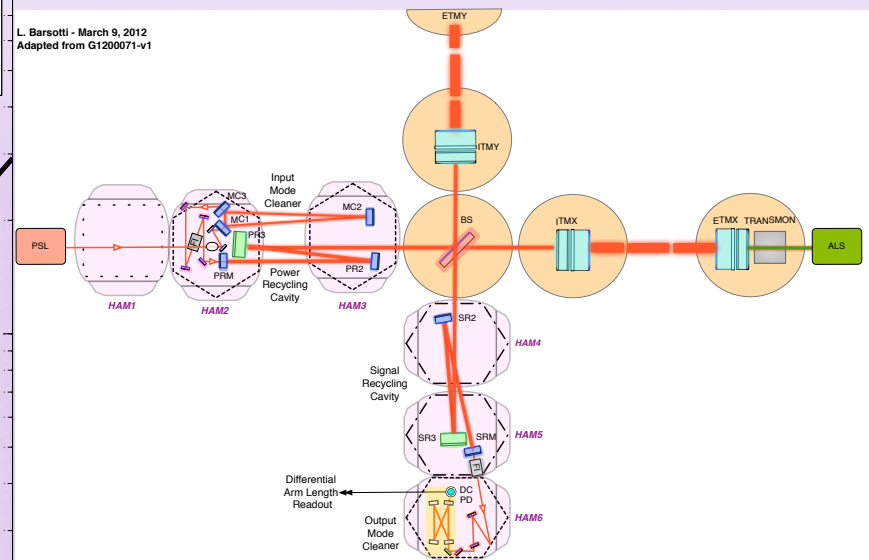
# More Clever: Quantum Noise in aLIGO

$$h_{Quantum} = \sqrt{\frac{1}{2}} \sqrt{\frac{8\hbar}{m\Omega^2 L^2}} \sqrt{\left( K_{SR} + \frac{1}{K_{SR}} \right)},$$

$$K_{SR} \sim \frac{8P_{Arm} \omega_0}{c^2 m \Omega^2} \frac{G_{sig}}{\left( 1 + \frac{\Omega^2}{\gamma_{src}^2} \right)}$$



L. Barsotti - March 9, 2012  
Adapted from G1200071-v1



- ✧ Arm cavities, power and signal recycling cavity
- ✧ Up to ~800 kW of light stored in the arms

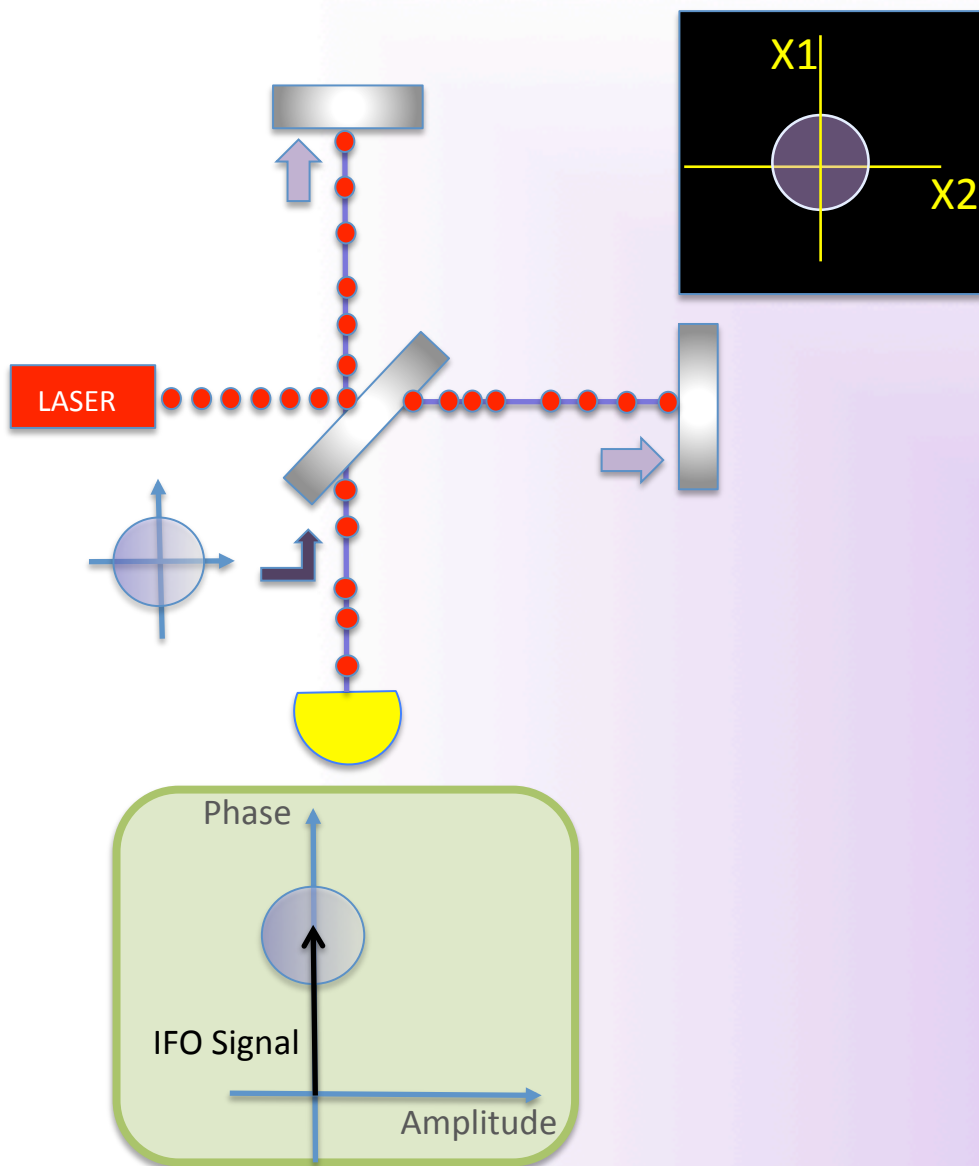
# How we go beyond aLIGO

- ✧ Again...make your interferometer longer!
- ✧ More power + heavier test masses
  - ✧ Already  $\sim 1\text{MW}$  in the arm cavities, need to compensate for thermal effects and instabilities
- ✧ (Even) more complex optical configuration which shapes the interferometer optical response

D. E. McClelland, N. Mavalvala, Y. Chen, and R. Schnabel, "Advanced interferometry, quantum optics and optomechanics in gravitational wave detectors", *Laser and Photonics Rev.*5, 677-696 (2011)

✧ Injection of squeezed states of vacuum

# Quantum Noise and Vacuum



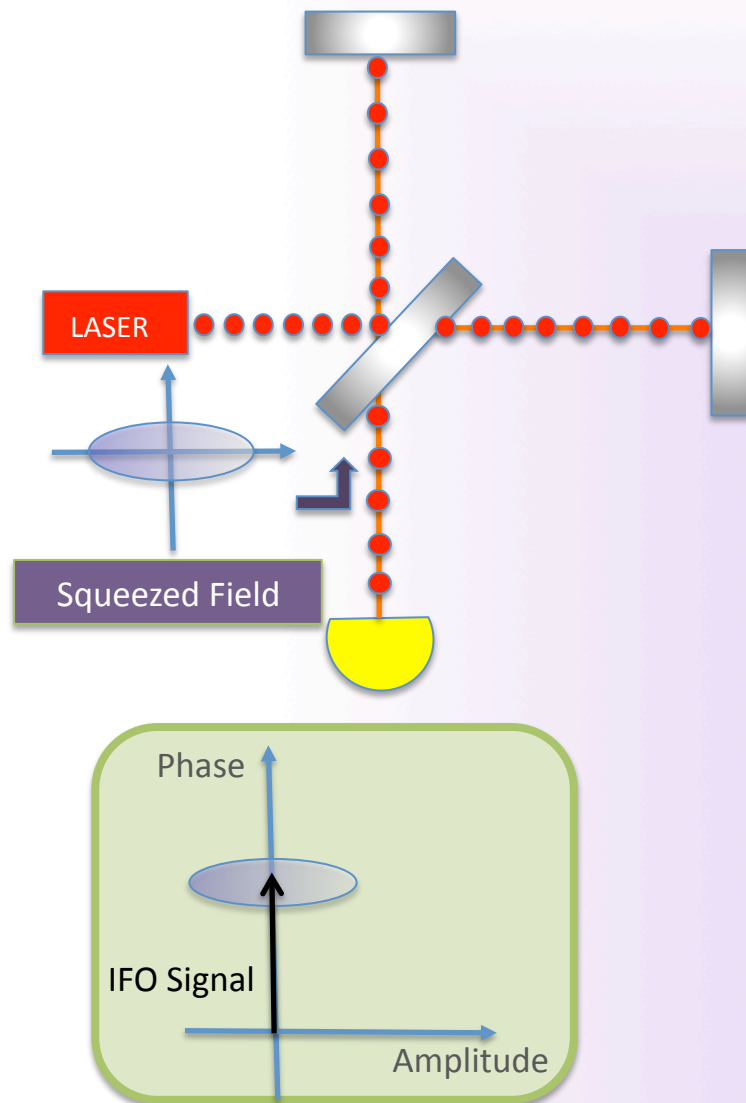
- ✧ Quantization of the electro-magnetic field
- ✧ When average amplitude is zero, the variance remains
- ✧ Heisenberg uncertainty principle:

$$\Delta X_1 \Delta X_2 \geq 1$$

- ✧ Vacuum fluctuations are everywhere that classically there is no field....
- ✧ ...like at the output port of your interferometer!

- ✧ Quantum noise is produced by vacuum fluctuations entering the open ports
- ✧ Vacuum fluctuations have equal uncertainty in phase and amplitude:
  - ❖ **Phase: Shot-Noise**  
(photon counting noise)
  - ❖ **Amplitude: Radiation Pressure Noise**  
(back-action)

# Vacuum Getting Squeezed



- ✧ Reduce quantum noise by injecting **squeezed vacuum**: less uncertainty in one of the two quadratures
- ✧ **Heisenberg uncertainty principle**: if the noise gets smaller in one quadrature, it gets bigger in the other one
- ✧ One can choose the relative orientation between the squeezed vacuum and the interferometer signal (**squeeze angle**)

C. M. Caves, Phys. Rev. Lett. 45, 75 (1980).

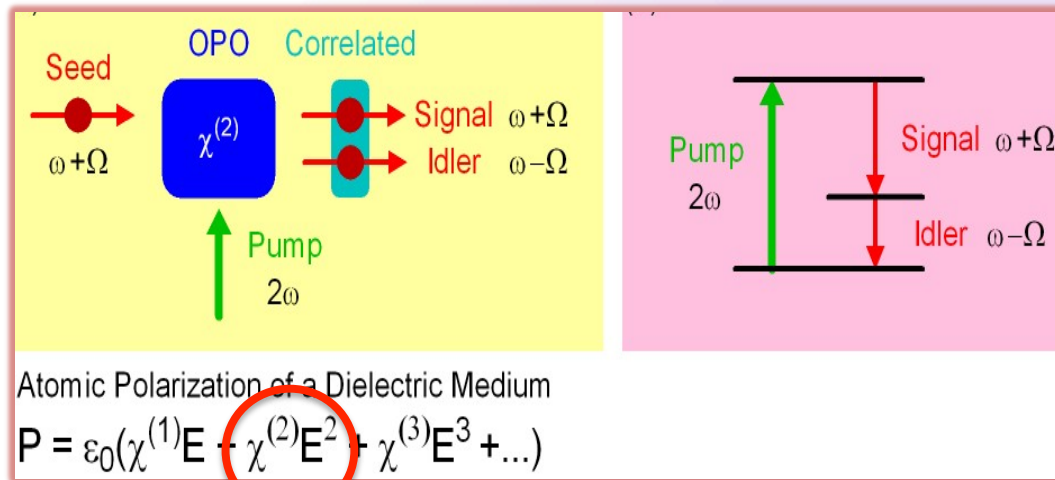
C. M. Caves, Quantum-mechanical noise in an interferometer. Phys. Rev. D 23, p. 1693 (1981).



# How to make squeezed fields..

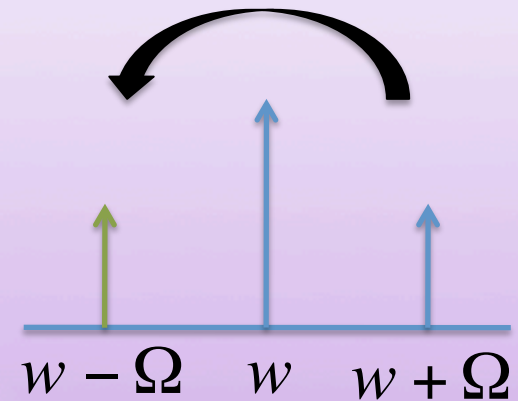
.... in theory

- ✧ Non linear medium with a strong second order polarization component
- ✧ Correlation of upper and lower quantum sidebands



$$P \propto (Ee^{-i2\omega t} + Ee^{-i(\omega+\Omega)t})^2$$

$$\Rightarrow Ee^{-i(\omega-\Omega)t}$$



The OPO makes a “copy” of the quantum sideband, and it correlates the sidebands

# How to make squeezed fields..

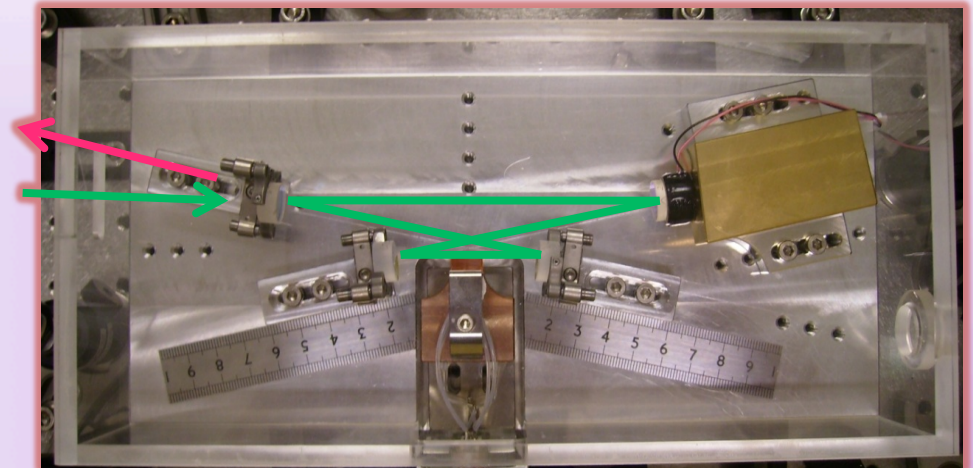
.... in practice

World-wide effort in the last 10 years to make squeezing in the audio-frequency band

✧ Lasers, mirrors, control loops,..

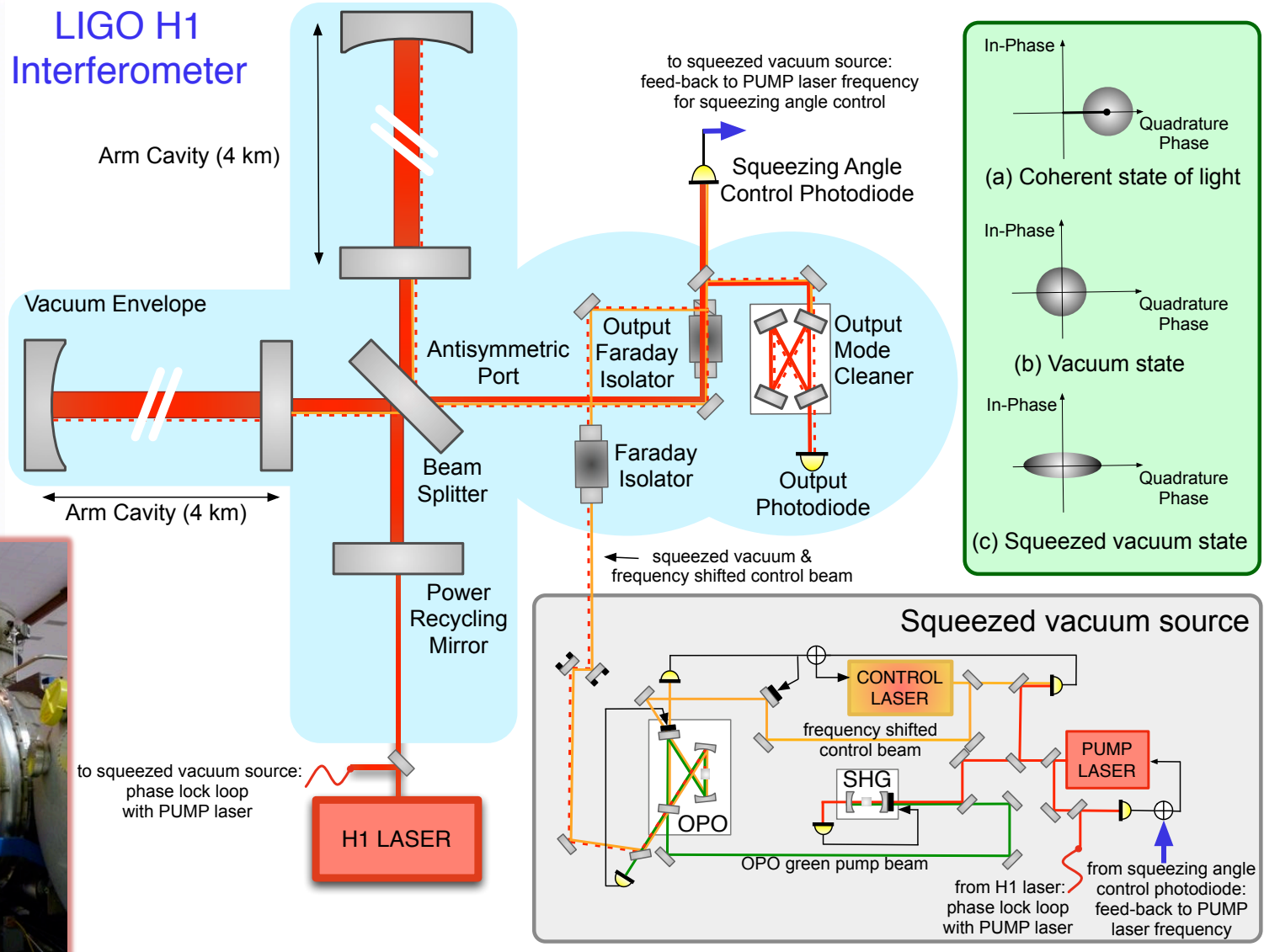


The Squeezer of the GEO600 detector



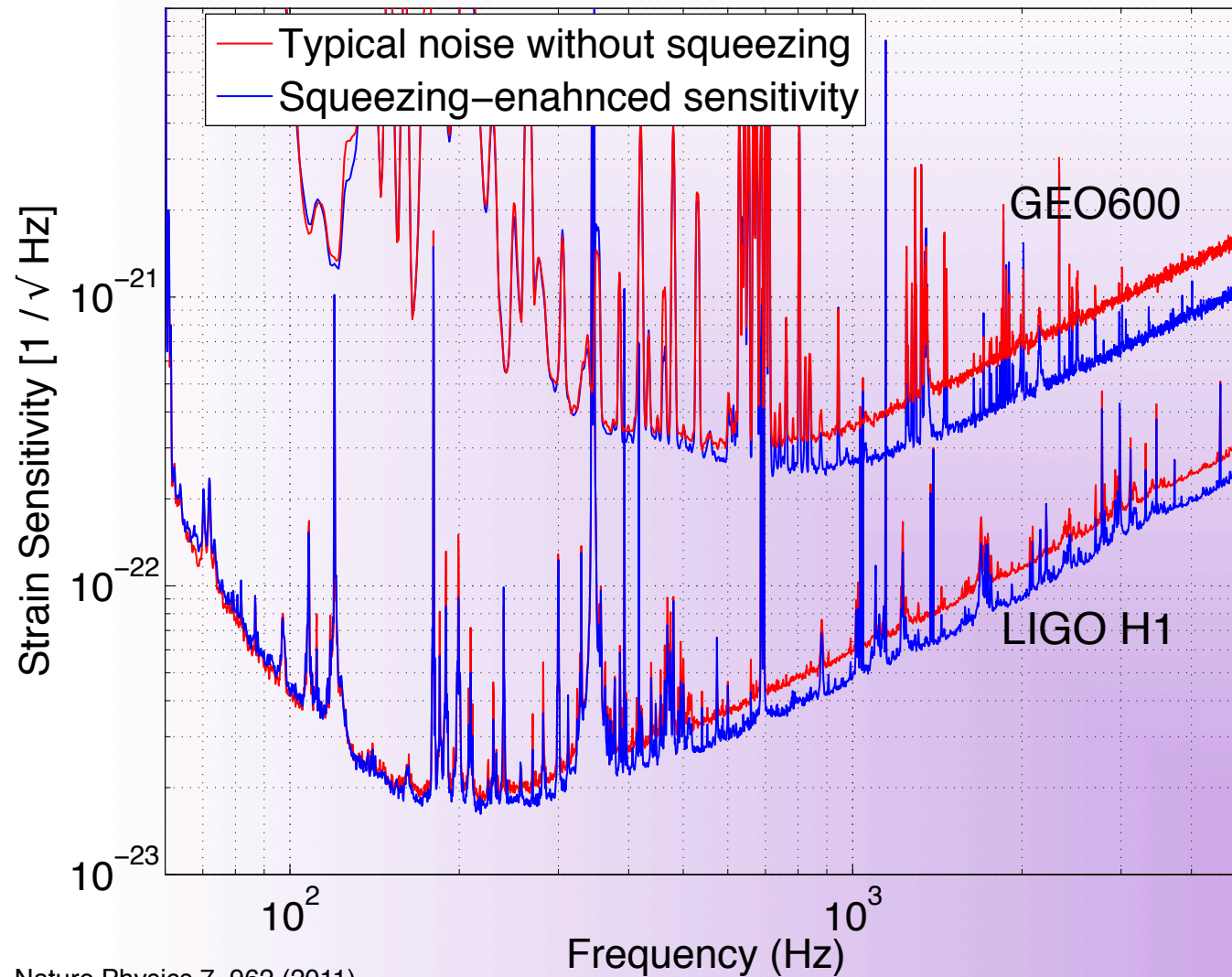
The Optical Parametric Oscillator of the LIGO squeezer (ANU design)

# How to inject squeezed fields





# Squeezing in GEO600 and LIGO H1



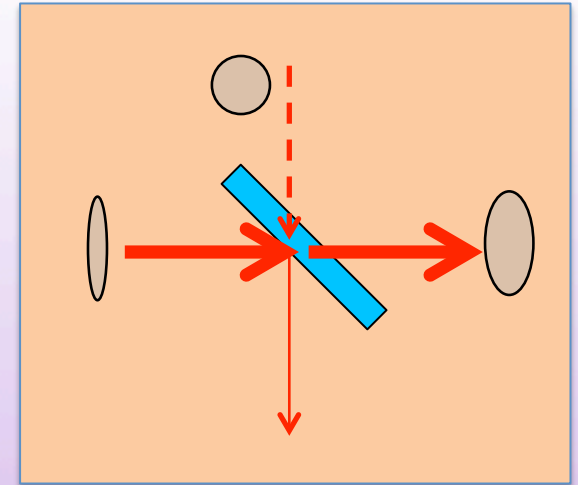
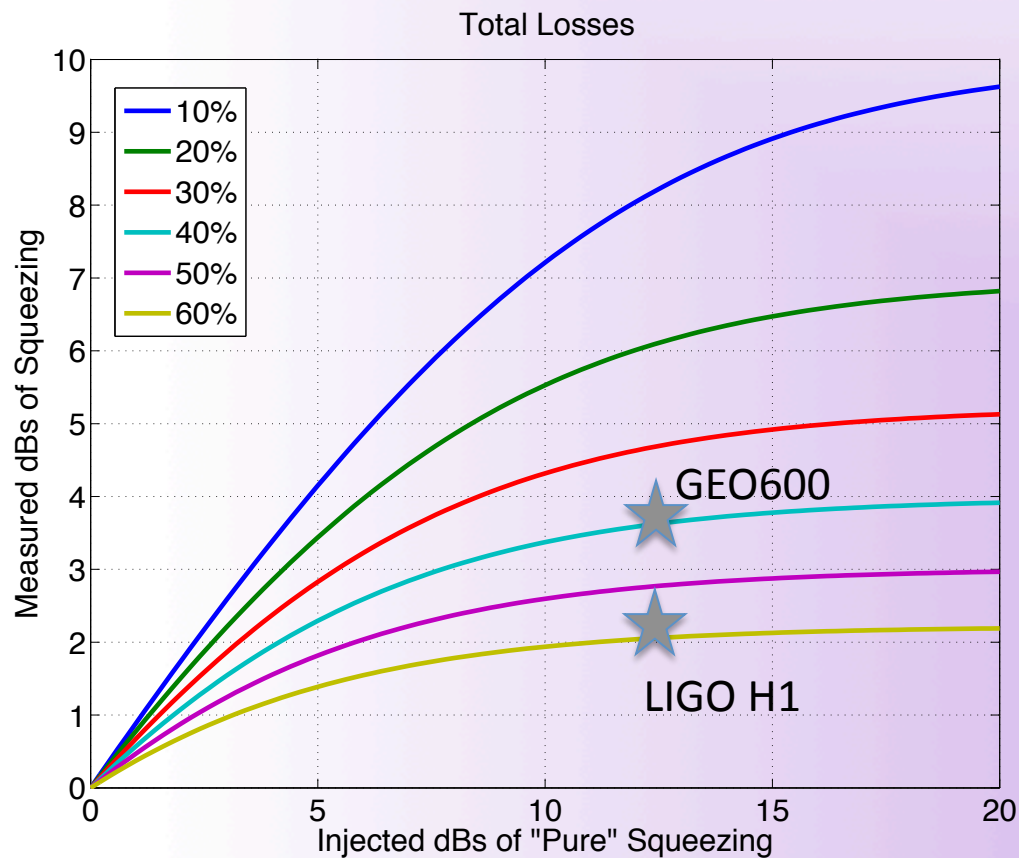
Abadie et al., Nature Physics 7, 962 (2011)  
GEO data are courtesy of H. Grote



# Lessons Learned (I)

## ✧ Losses are very unforgiving

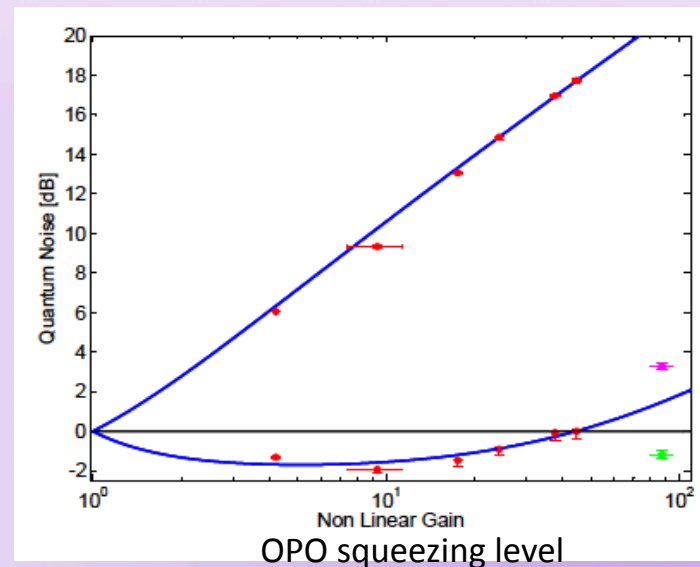
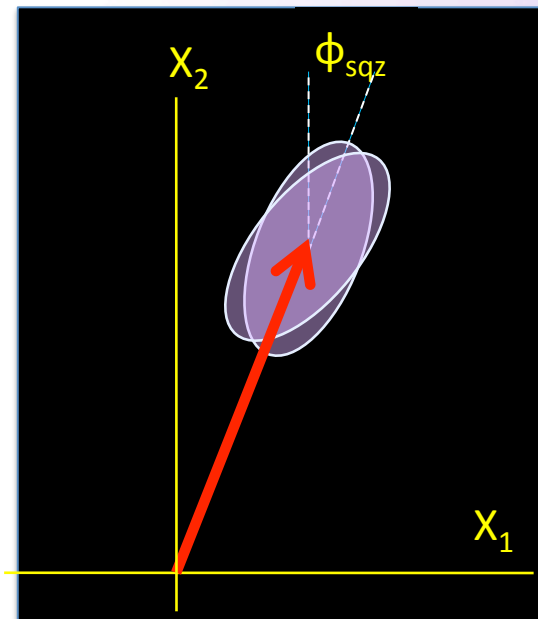
- ✧ GEO aimed for 6 dB → got 3.5 dB
- ✧ LIGO aimed for 3 dB → got 2.15 dB



- ✧ Mode matching
- ✧ Faradays
- ✧ OMC transmission
- ✧ ...

# Lessons Learned (II)

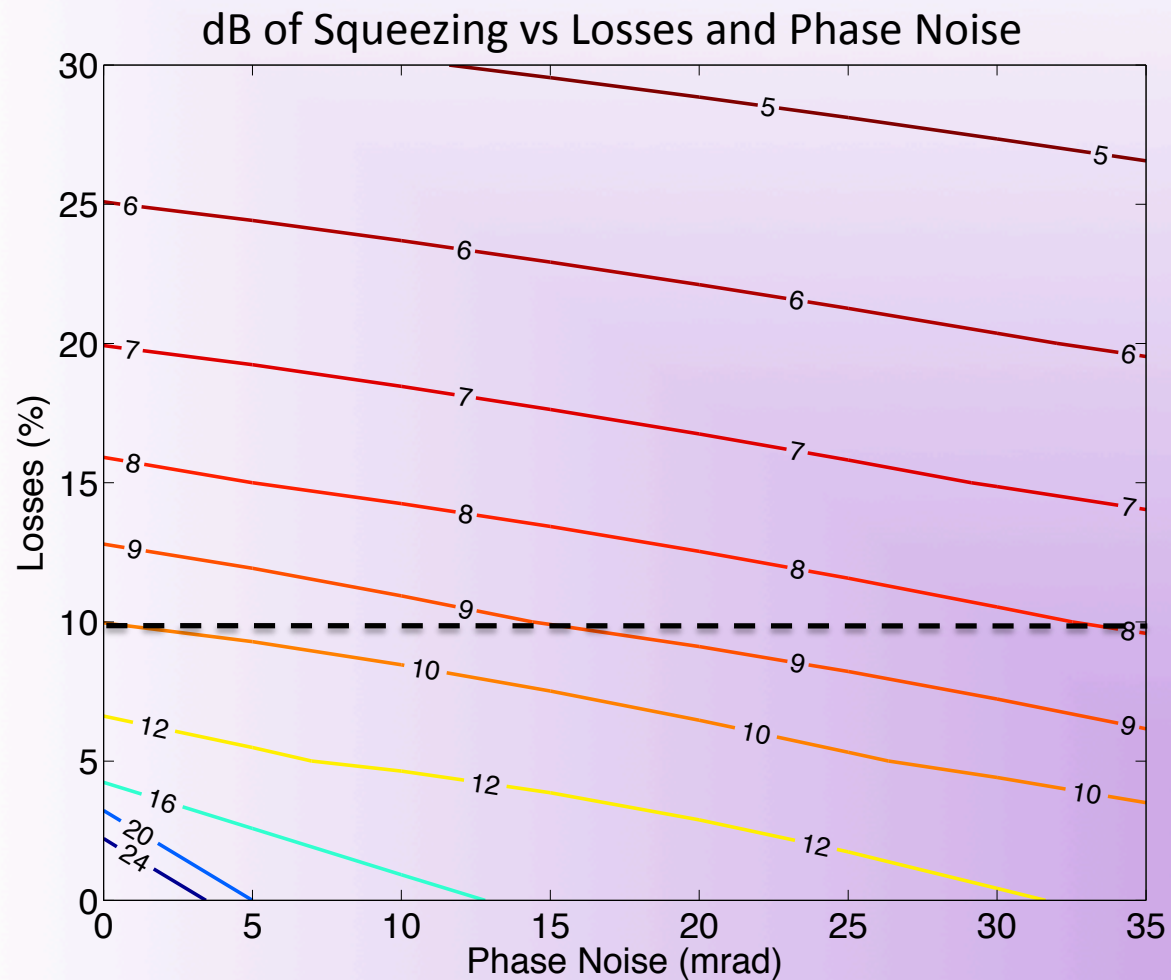
- ✧ Phase noise between squeezed field and interferometer was dependent on interferometer alignment:
  - ✧ Static misalignments will cause a change in the demodulation phase needed to detect the maximum squeezing
  - ✧ Beam jitter will add phase noise when beating against a static misalignment.



→ at best  $\sim 40$  mrad RMS,  
only  $\sim 25$  mrad from squeezer source

# To keep in mind for the future

With 10% total losses, you can't afford any phase noise at all, if you want to measure 10dB of squeezing

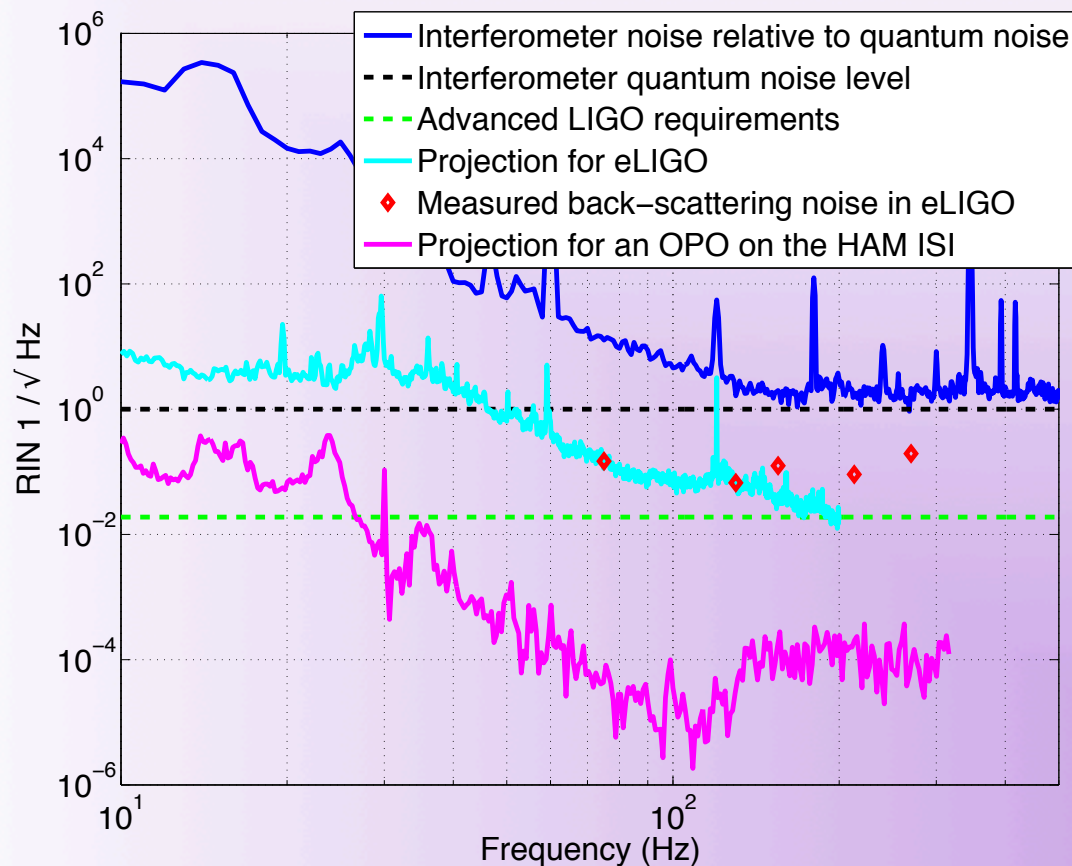




# Lessons Learned (III)

## ✧ Need better isolation from back scattering

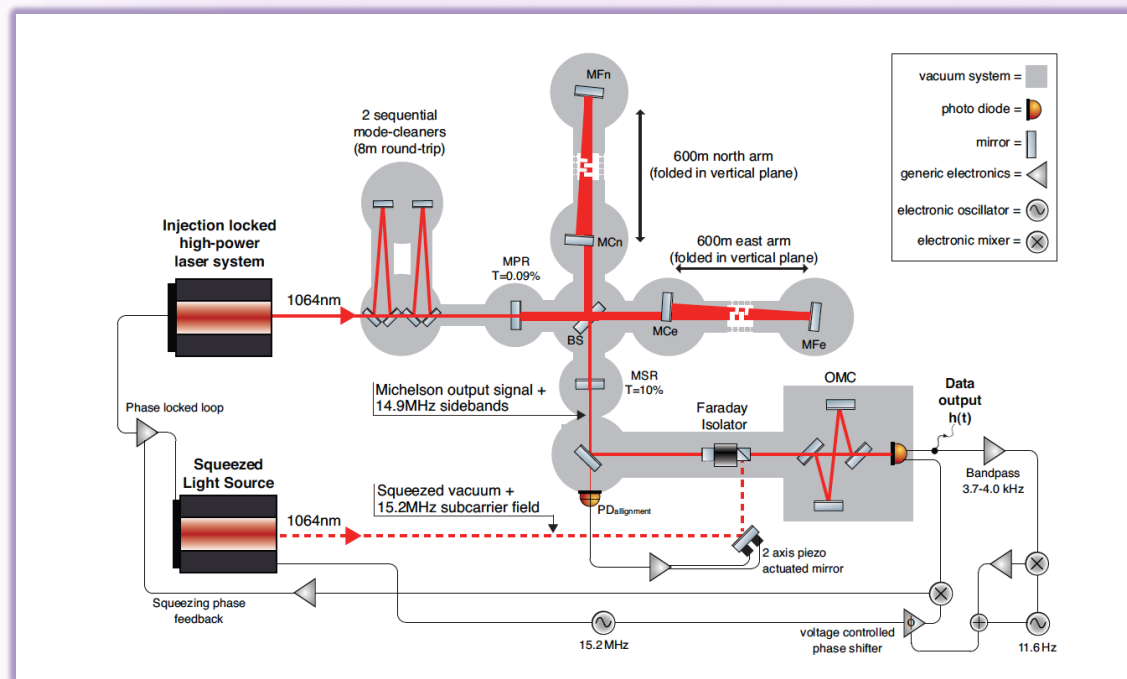
(it was ok for LIGO H1, it won't be enough for aLIGO)



Impact of backscattered-light in a squeezing-enhanced interferometric gravitational-wave detector, S. Chua et al. (in preparation)

# Lessons Learned (VI)

- ✧ From GEO600: Squeezing angle control signals from 1% pick-off are bad
- New “a-la-Hartmut” strategy (use transmission signals from the OMC)

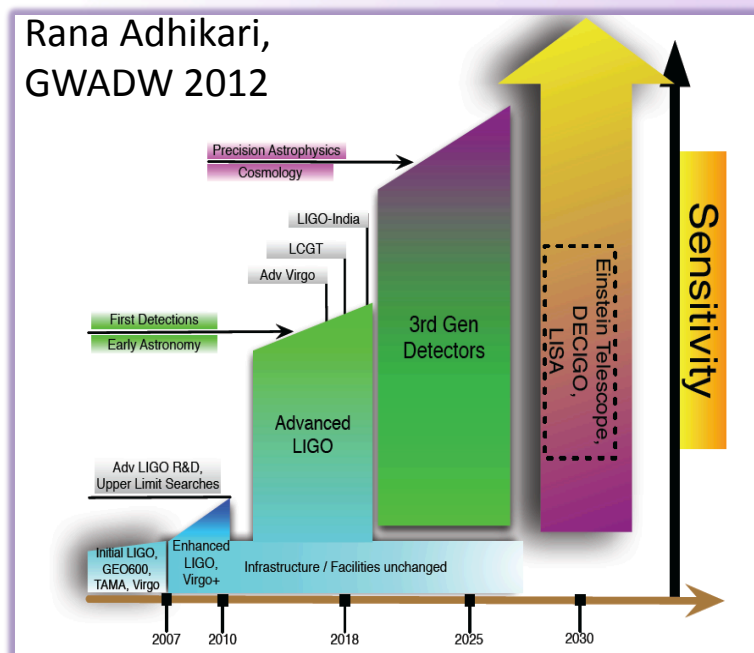


## First Long-Term Application of Squeezed States of Light in a Gravitational-Wave Observatory

H. Grote,<sup>1,\*</sup> K. Danzmann,<sup>1</sup> K.L. Dooley,<sup>1</sup> R. Schnabel,<sup>1</sup> J. Slutsky,<sup>1</sup> and H. Vahlbruch<sup>1</sup>

# How about squeezing in aLIGO (and beyond)?

- ✧ Do we want it?
- ✧ Do we know how to make it? How do we incorporate all the “lessons learned” in the next generation of squeezers?

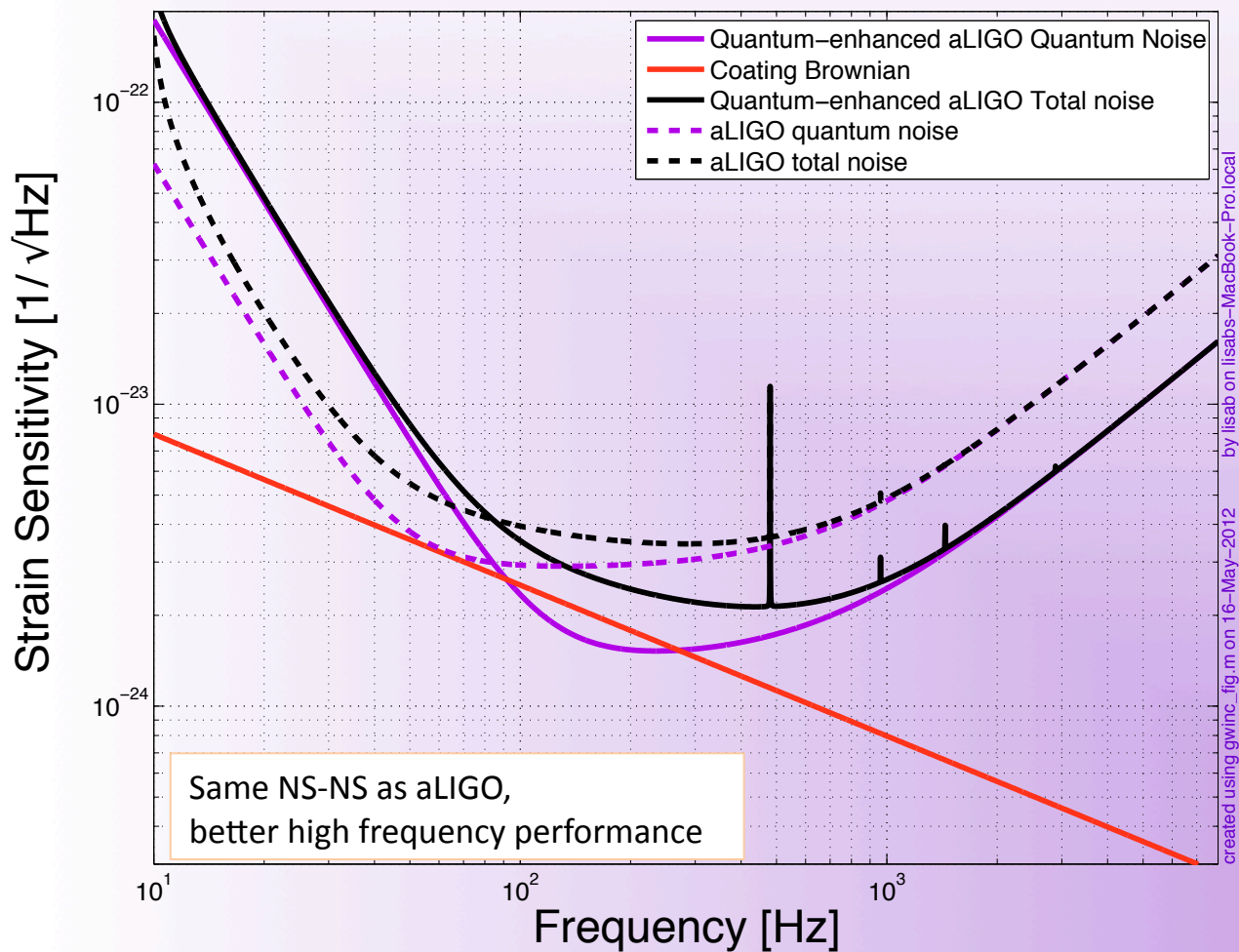


LIGO technical note T1200008-v3  
Comparison of Quantum Noise in 3G  
Interferometer Configurations,  
Haixing Miao et al.

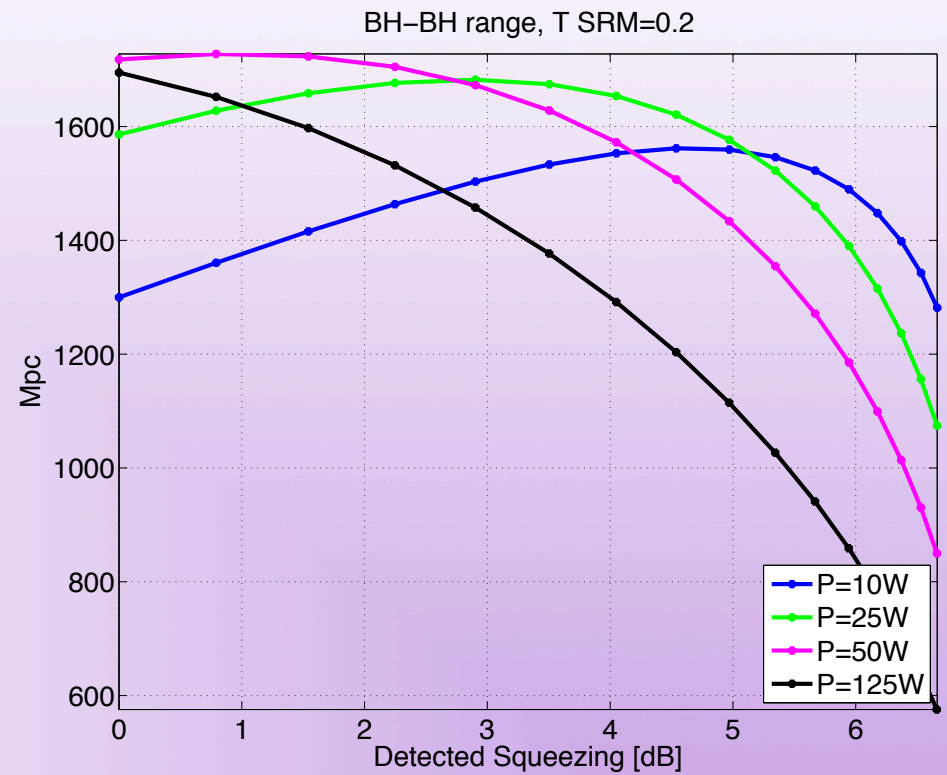
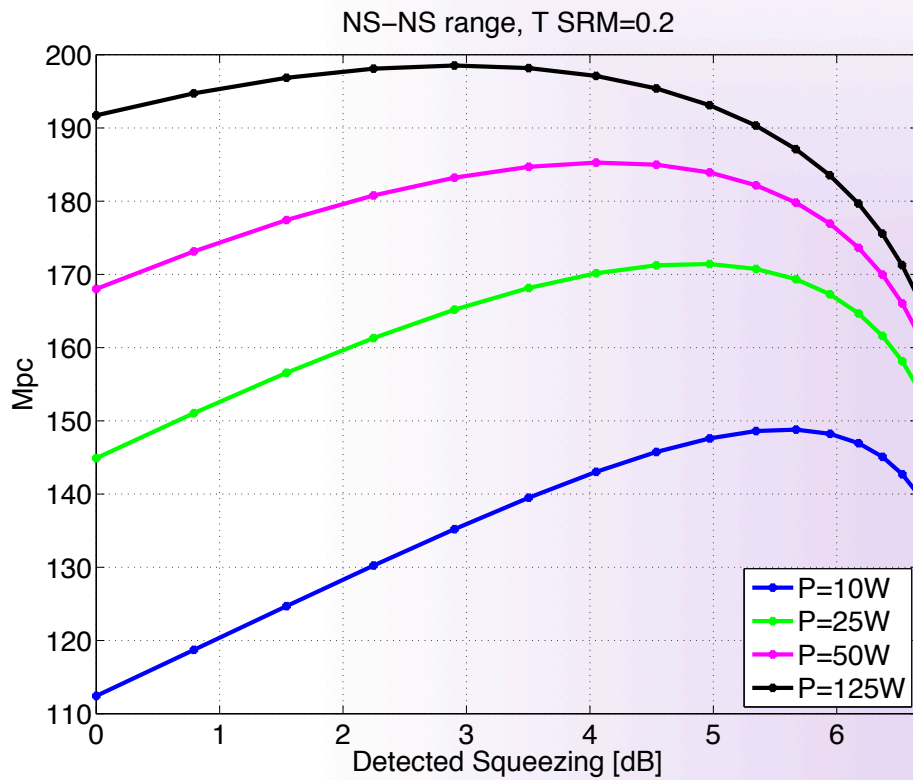
➔ 10 dB of squeezing needed for all future configurations..

# How squeezing in aLIGO would look

Projections for a “Quantum-Enhanced Advanced LIGO”

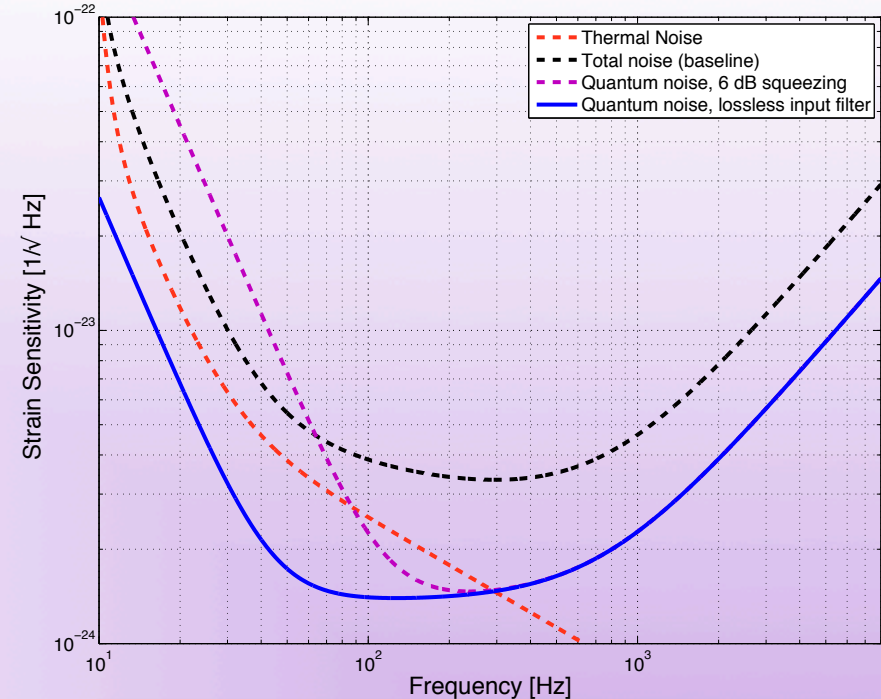
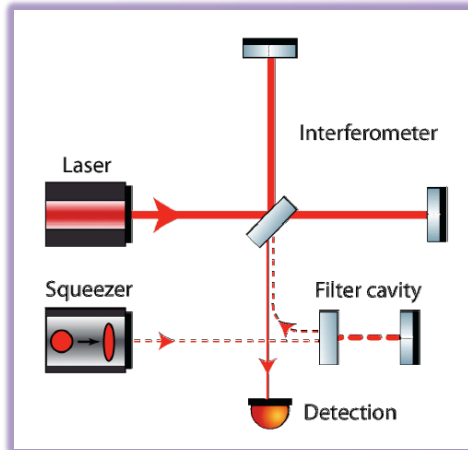


# aLIGO + Squeezing: NS-NS and BH-BH Ranges



# What we really want: Frequency Dependent Squeezing

High finesse detuned cavity which rotates the squeezing angle as function of frequency



PHYSICAL REVIEW D, VOLUME 65, 022002

**Conversion of conventional gravitational-wave interferometers into quantum nondemolition interferometers by modifying their input and/or output optics**

H. J. Kimble,<sup>1</sup> Yuri Levin,<sup>2,\*</sup> Andrey B. Matsko,<sup>3</sup> Kip S. Thorne,<sup>2</sup> and Sergey P. Vyatchanin<sup>4</sup>

PHYSICAL REVIEW D 68, 042001 (2003)

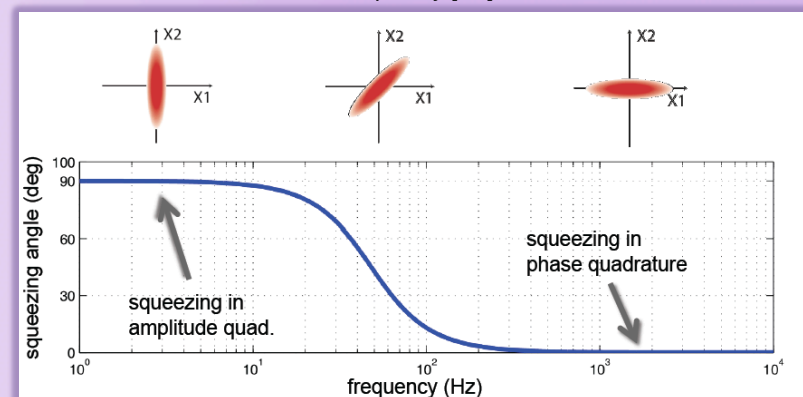
**Squeezed-input, optical-spring, signal-recycled gravitational-wave detectors**

Jan Harms,<sup>1</sup> Yanbei Chen,<sup>2</sup> Simon Chelkowski,<sup>1</sup> Alexander Franzen,<sup>1</sup> Henning Vahlbruch,<sup>1</sup> Karsten Danzmann,<sup>1</sup> and Roman Schnabel<sup>1</sup>

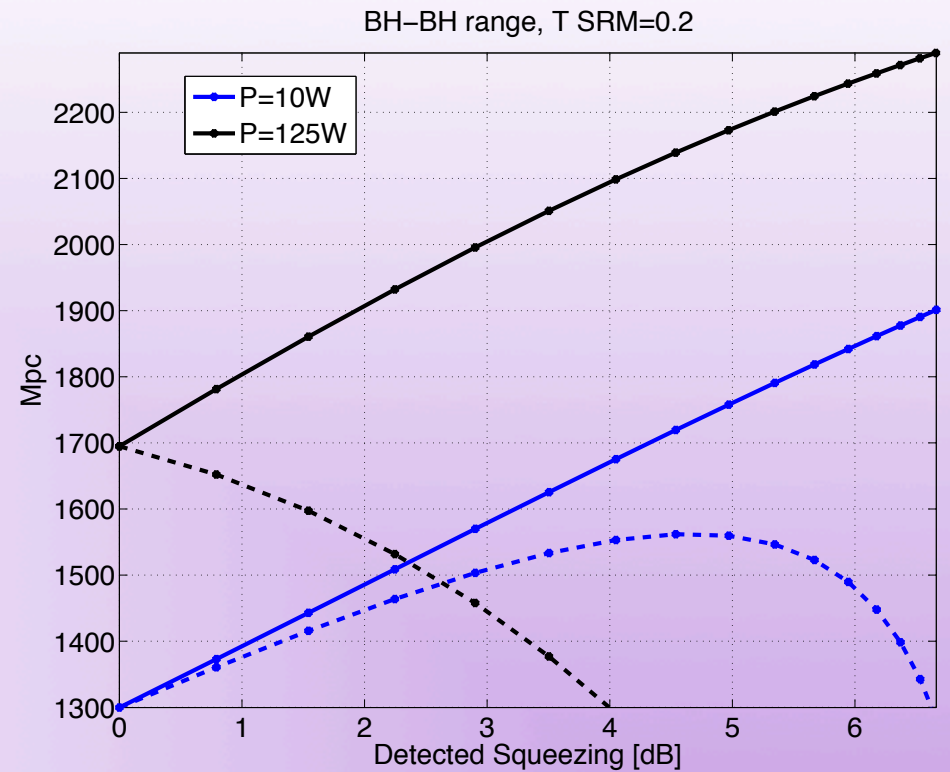
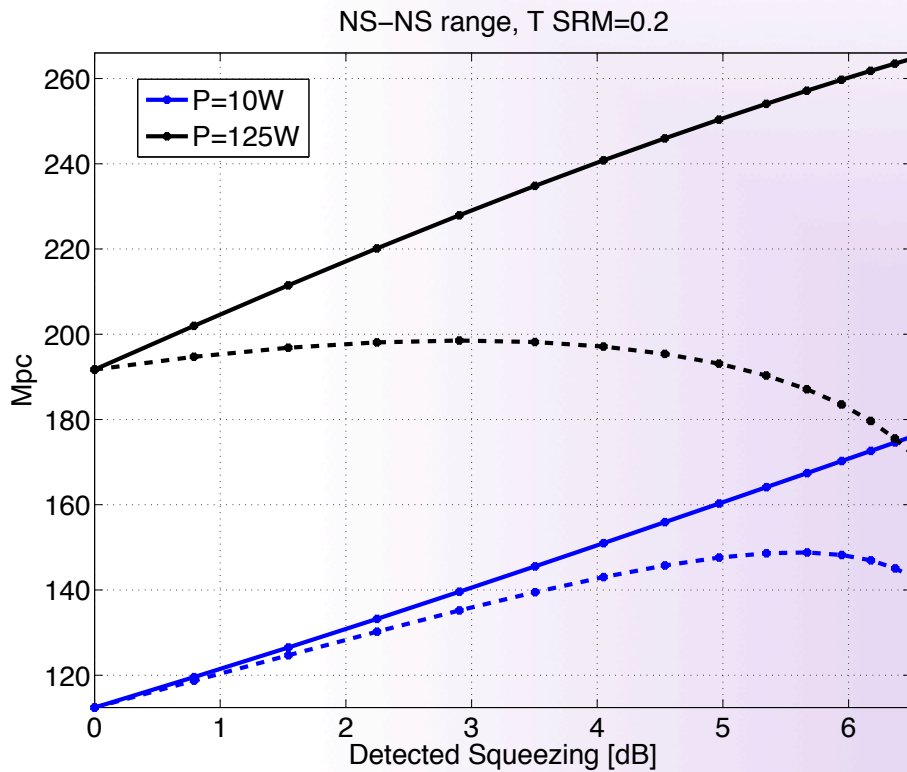
PHYSICAL REVIEW A 71, 013806 (2005)

**Experimental characterization of frequency-dependent squeezed light**

Simon Chelkowski, Henning Vahlbruch, Boris Hage, Alexander Franzen, Nico Lastzka, Karsten Danzmann, and Roman Schnabel



# aLIGO + Frequency Dependent Squeezing: NS-NS and BH-BH Ranges



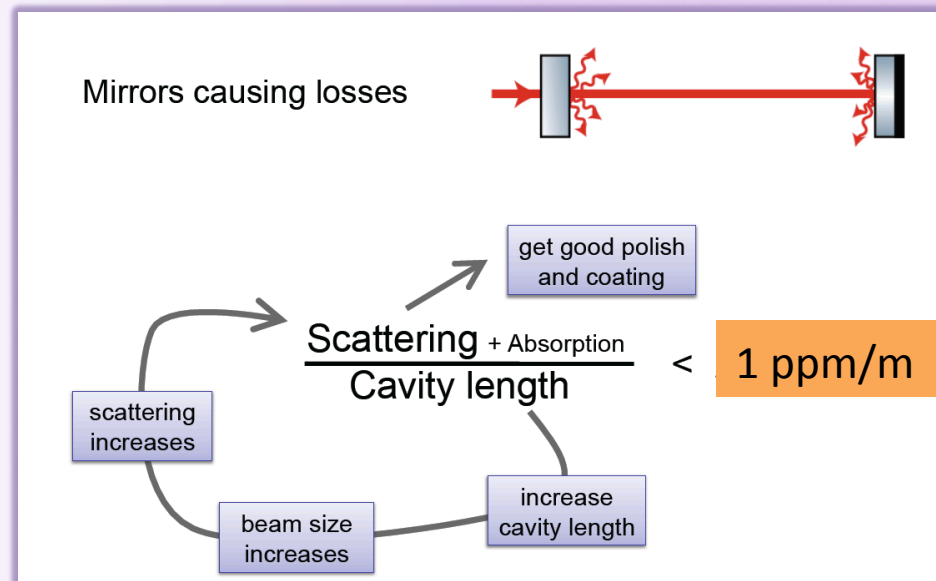


# Nothing comes cheap: losses again..

Losses in a filter cavity deteriorate, if too high, make the filter cavity useless...

$$\text{Total Loss } E = \frac{4\varepsilon}{T} = \frac{\varepsilon}{L} \frac{c}{\gamma_{filter}}, \quad \gamma_{filter} = \frac{Tc}{4L}$$

- ✧ Per-round-trip loss depends on the beam spot size  
(big beam size  $\rightarrow$  higher scatter losses), which depends on L



# Squeezing @ MIT

## FILTER CAVITY EXPERIMENT

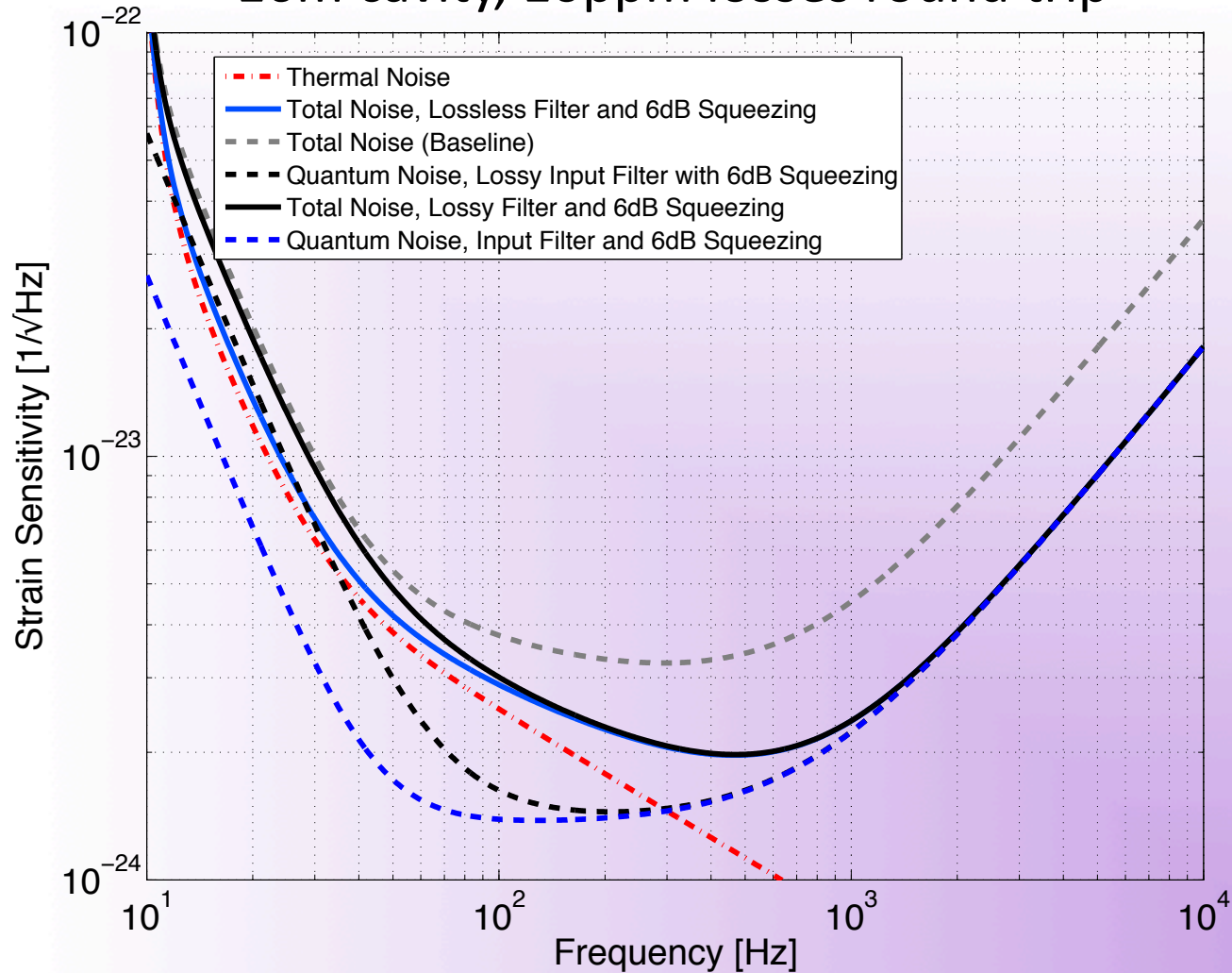
- ✧ Measuring optical losses to determine Advanced LIGO filter cavity design
- ✧ Implementing practical filter cavity control scheme
- ✧ Characterizing technical noises
- ✧ Preparing for demonstration of audio-band frequency dependent squeezing

NEW SQUEEZER SOURCE compatible with aLIGO requirements



Tomoki Isogai, John Miller, Eric Oelker, (Patrick Kwee)

# For aLIGO, we could afford a “lossy” cavity 16m cavity, 10ppm losses round trip



## Realistic Filter Cavities for Advanced Gravitational Wave Detectors

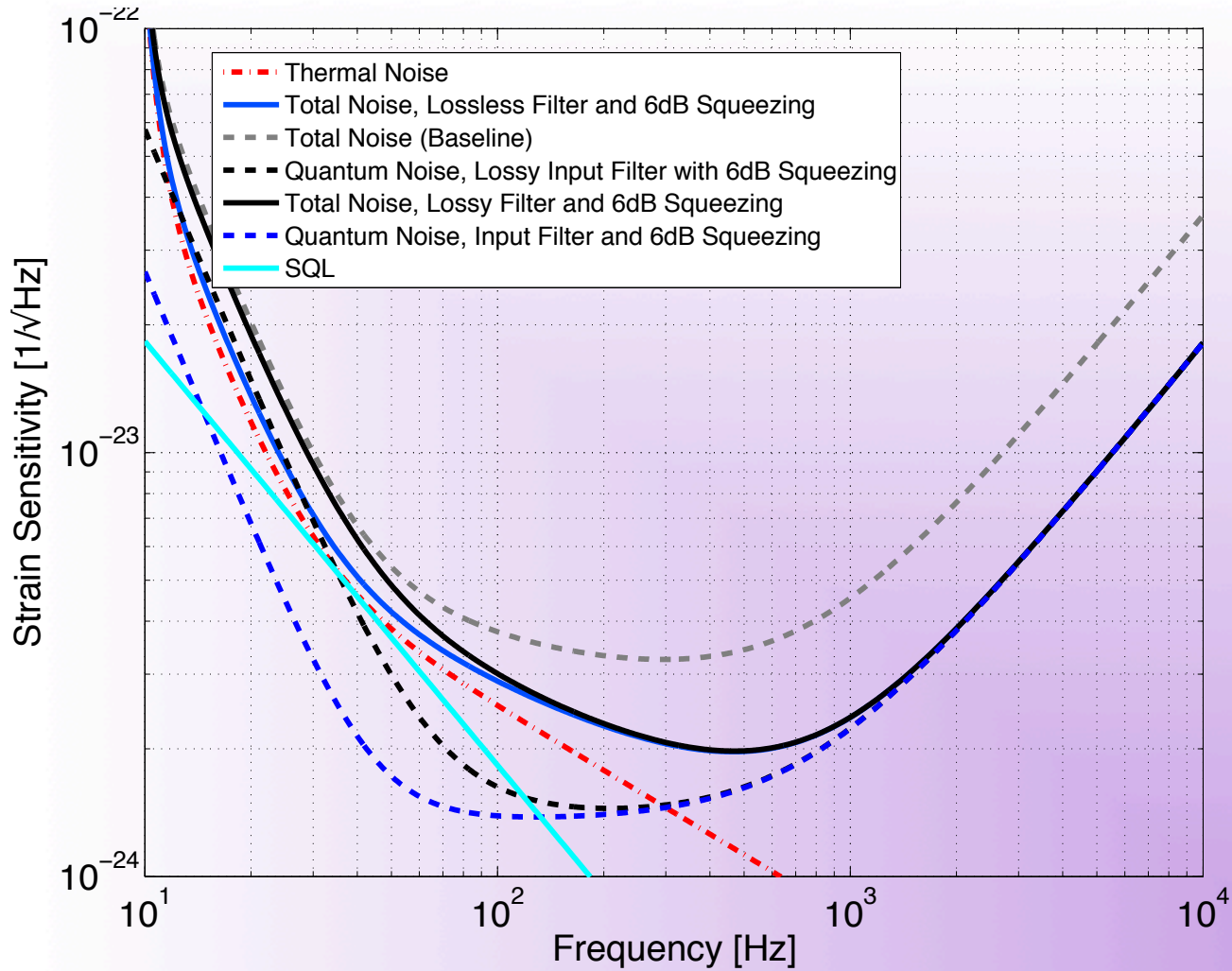
M. Evans,<sup>1</sup> L. Barsotti,<sup>1</sup> J. Harms,<sup>2</sup> P. Kwee,<sup>1</sup> and H. Miao<sup>2</sup>

<sup>1</sup>MIT

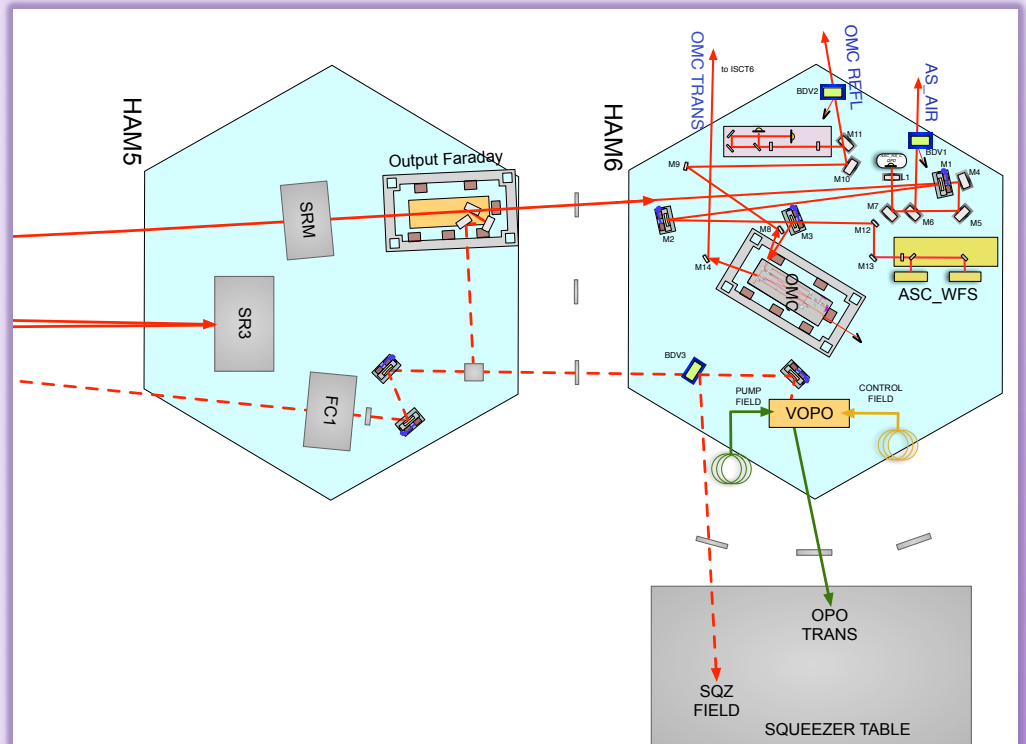
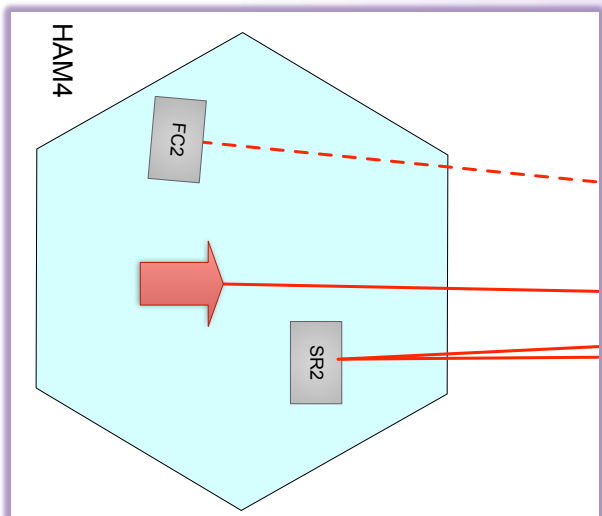
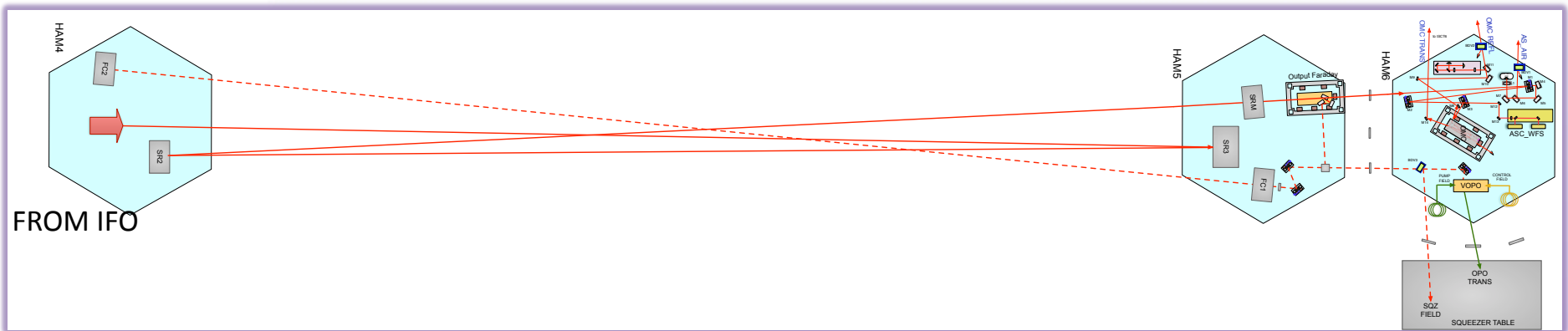
<sup>2</sup>Caltech

In preparation

# Beyond the “Standard Quantum Limit”



# Something like this, maybe....



Just a cartoon!  
Not a conceptual design yet!



# H1 Squeezing Experiment



LHO: Daniel Sigg, Keita Kawabe, Robert Schofield, Cheryl Vorvick, Dick Gustafson (Univ Michigan), Max Factourovich (Columbia), Grant Meadors (Univ Michigan), the LHO staff

MIT: Sheila Dwyer, L. Barsotti, Nergis Mavalvala, Nicolas Smith-Lefebvre, Matt Evans

ANU: Sheon Chua, Michael Stefszky, Conor Mow-Lowry, Ping Koy Lam, Ben Buchler, David McClelland

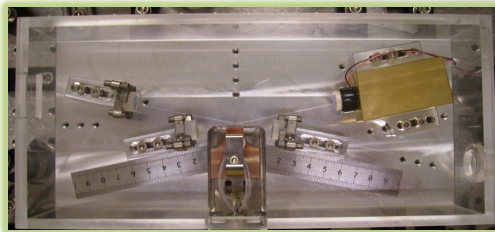
AEI: Alexander Khalaidovski, Roman Schnabel



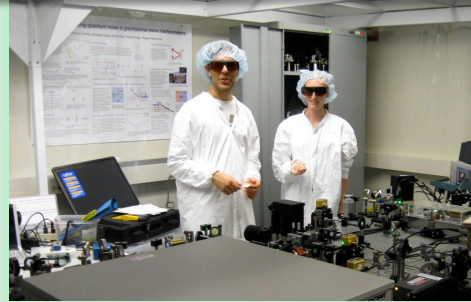
# Proposal for a Squeezed H1 Interferometer

Daniel Sigg, Nergis Mavalvala, David McClelland, Ping Koy Lam, Roman Schnabel, Henning Vahlbruch and Stan Whitcomb

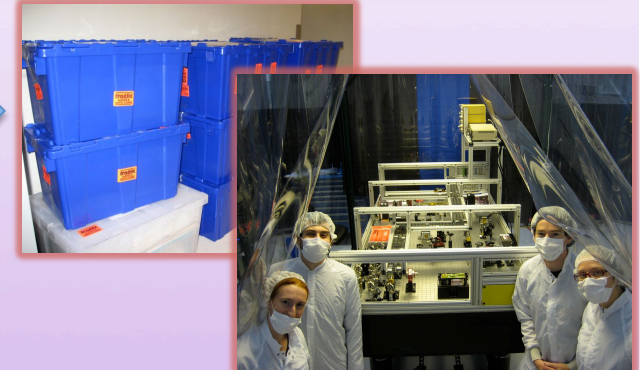
Bow-tie cavity OPO design at ANU (2008)



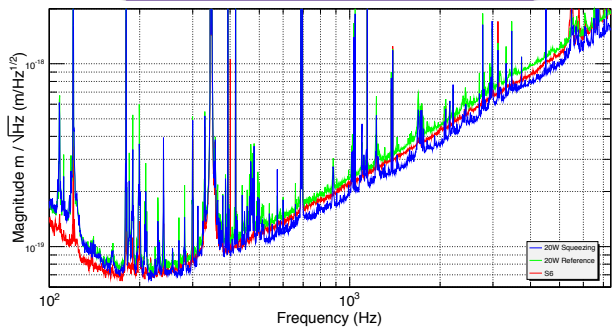
H1 Squeezer assembling at MIT (2009-2010)



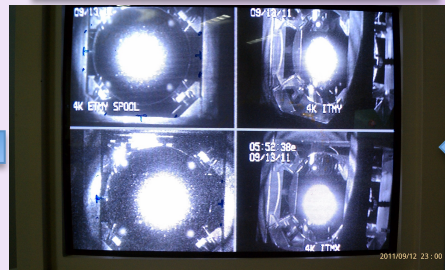
H1 Squeezer parts shipped to LHO (Oct 2010)



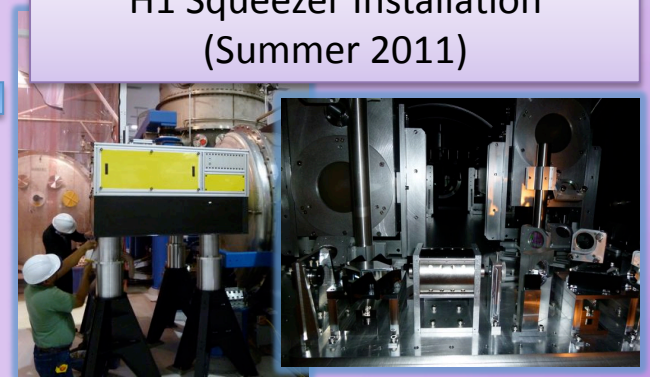
Squeezing in H1 (Oct 3 – Dec 4)



H1 Recovery (Sept 2011)



H1 Squeezer Installation (Summer 2011)

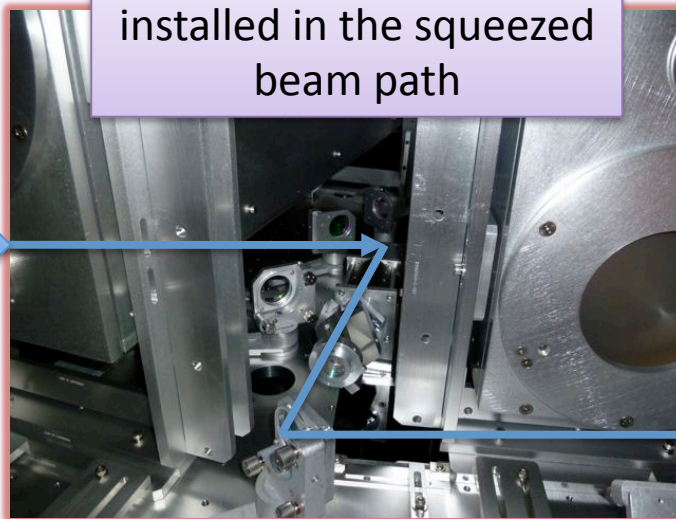




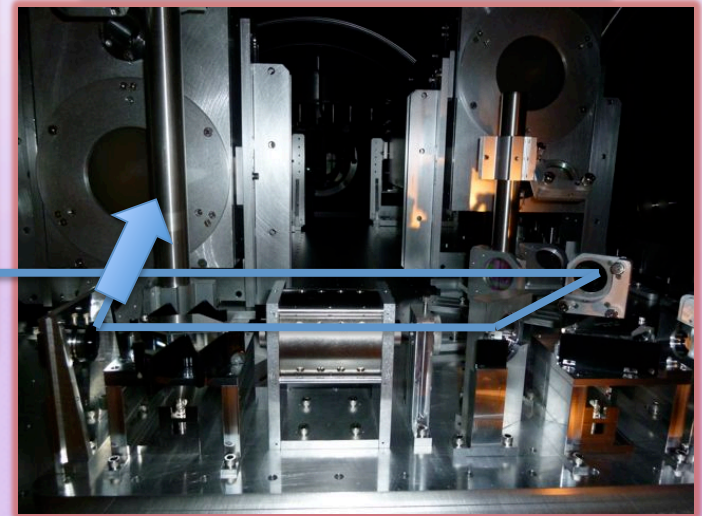
# H1 Squeezing Experiment: Squeezer Installation



Squeezer table  
craned to its final  
location



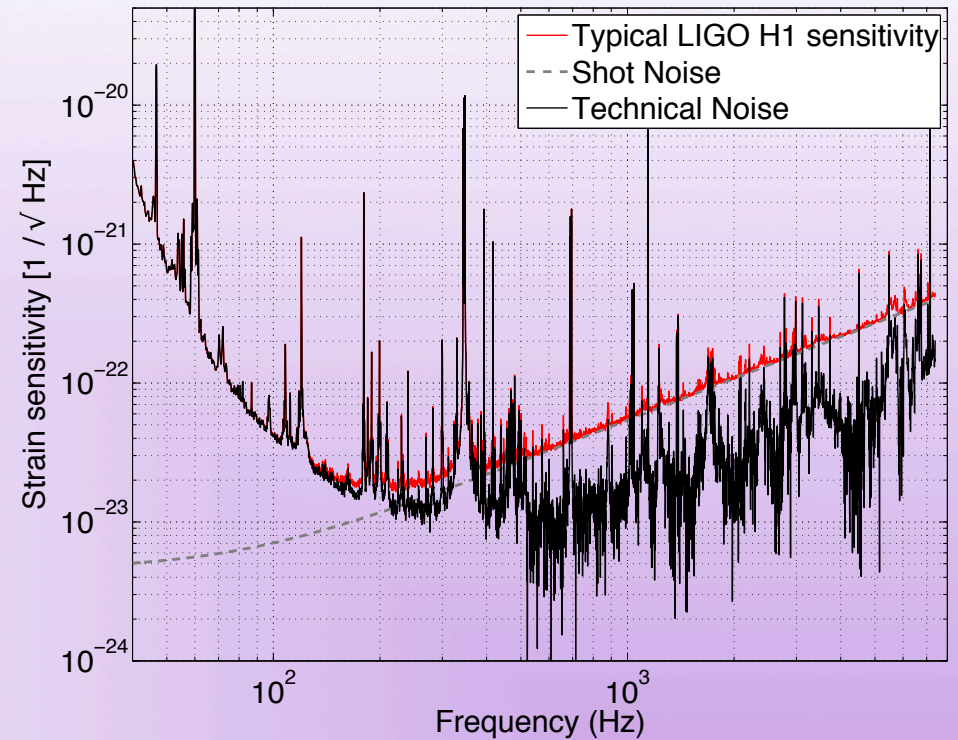
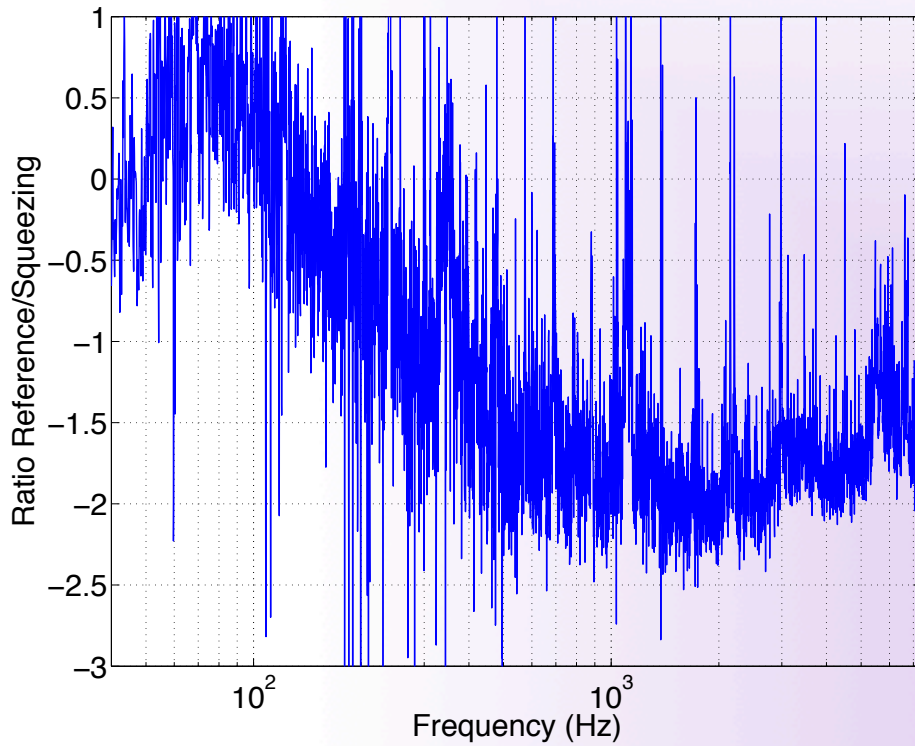
Additional Faraday  
installed in the squeezed  
beam path



New H1 Output Faraday  
(first aLIGO unit)

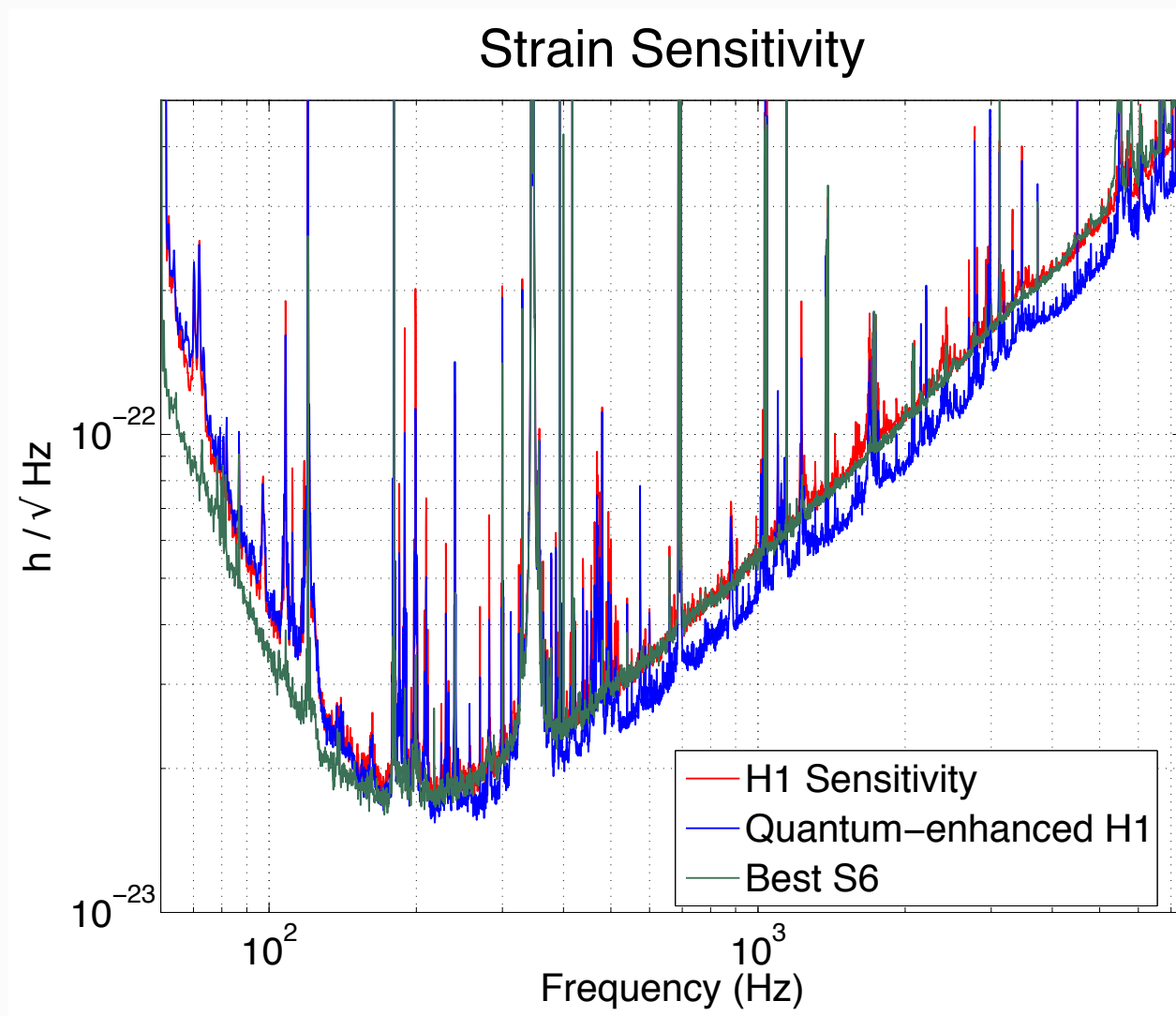


# 2.15 dB (28%) improvement over quantum noise

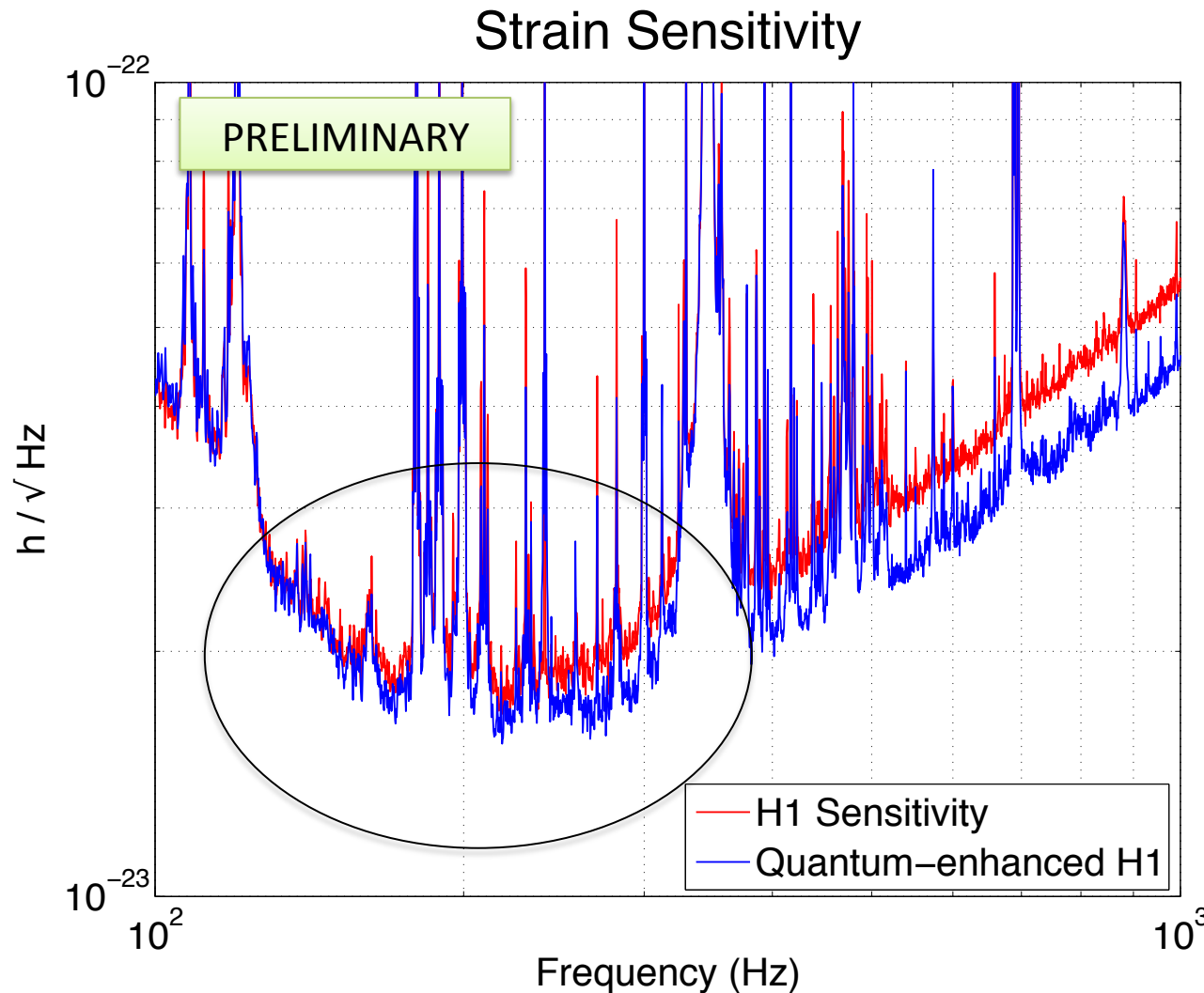


Squeezing improves only quantum noise, not other technical noises

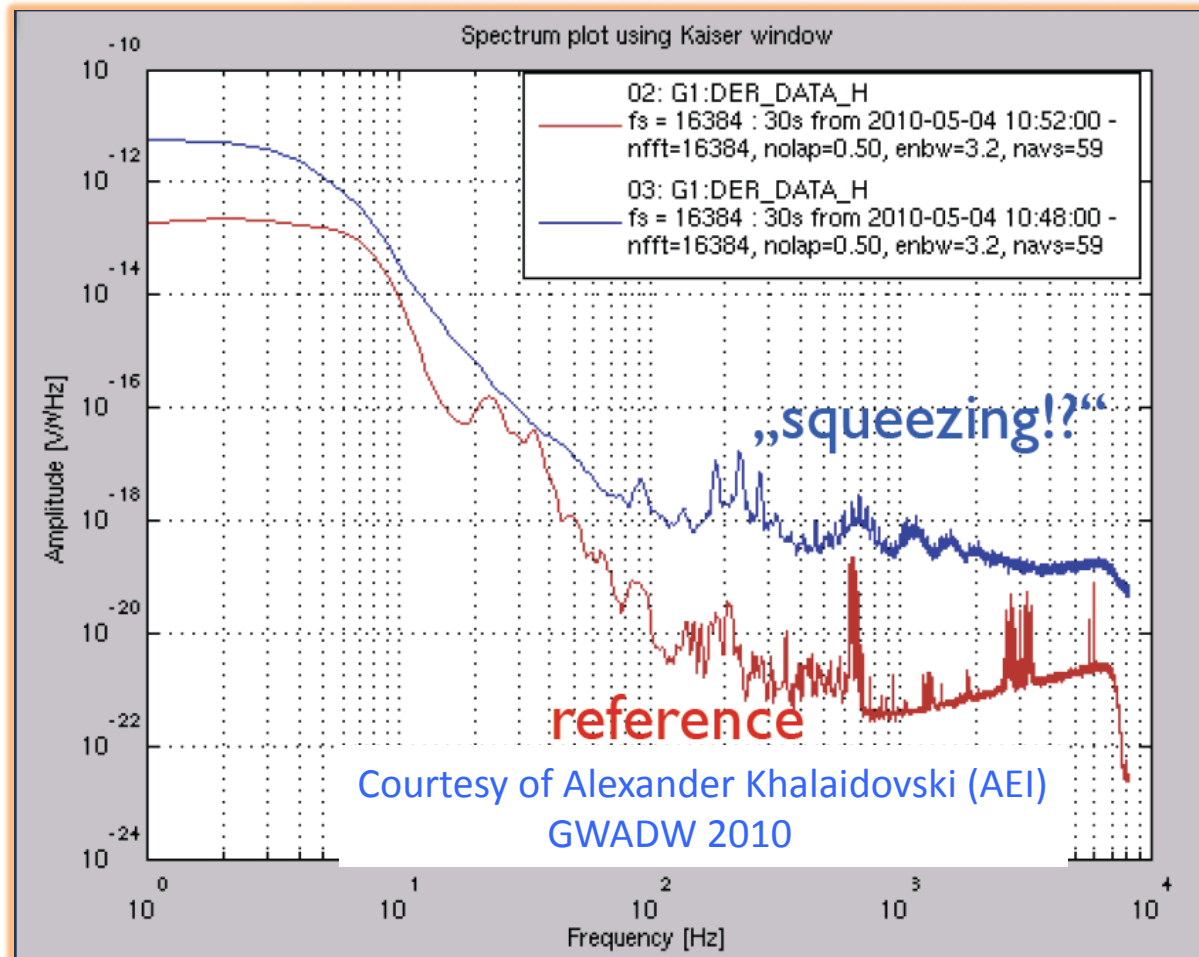
# Best broadband sensitivity ever



Improving H1 by 2 dB (28%) with squeezing  
..without spoiling the sensitivity at 200 Hz

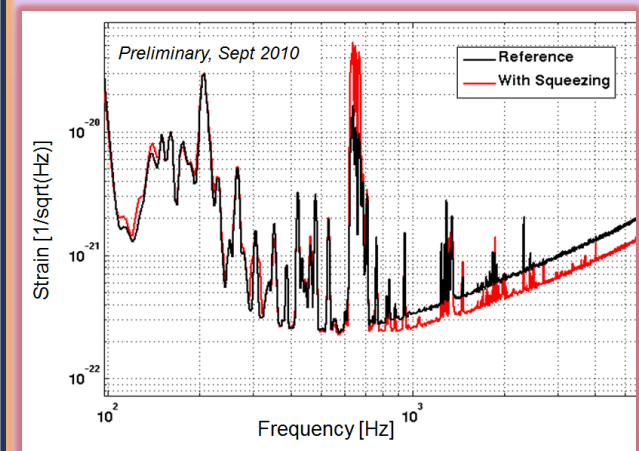


# First try at squeezing in GEO



✧ First squeezing injection: back scattered noise limits the sensitivity

✧ Additional Faraday to reduce back scattering and measure squeezing





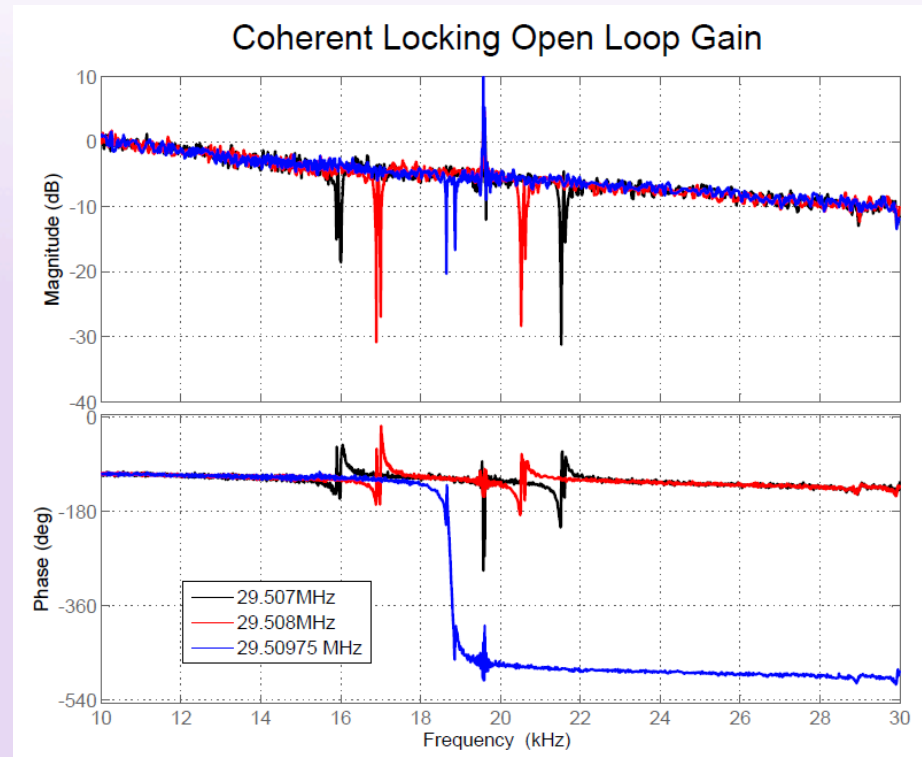
## Where the main losses came from

- ✧ Mode matching (~30% losses)
  - ✧ Faradays (3 passes ~ 20% losses)
  - ✧ OMC transmission (18% losses)
- ➔ “Technical” problems, total losses should be down to 10-15% in aLIGO

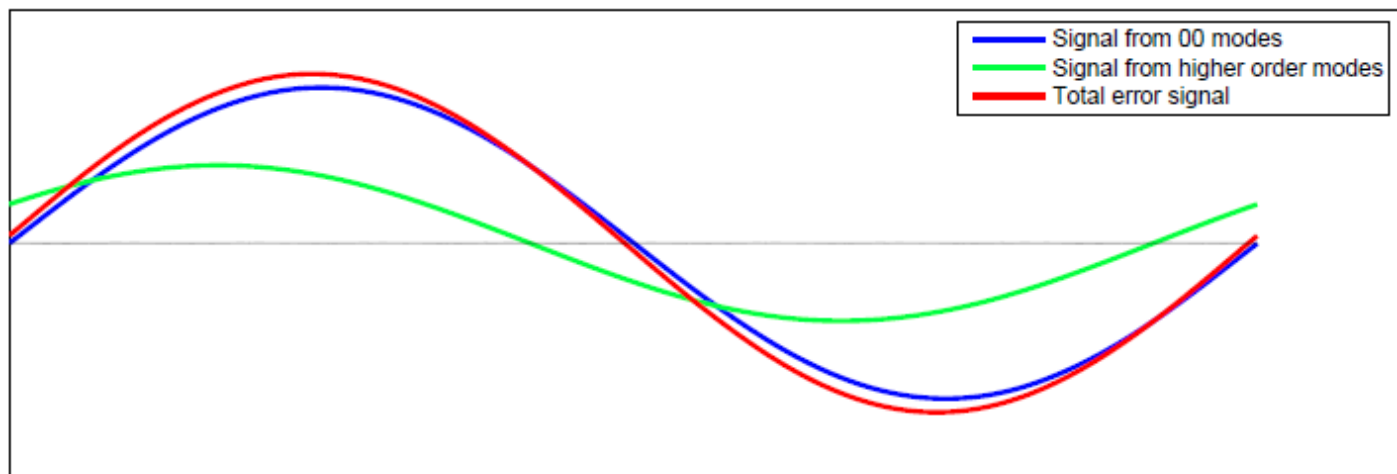
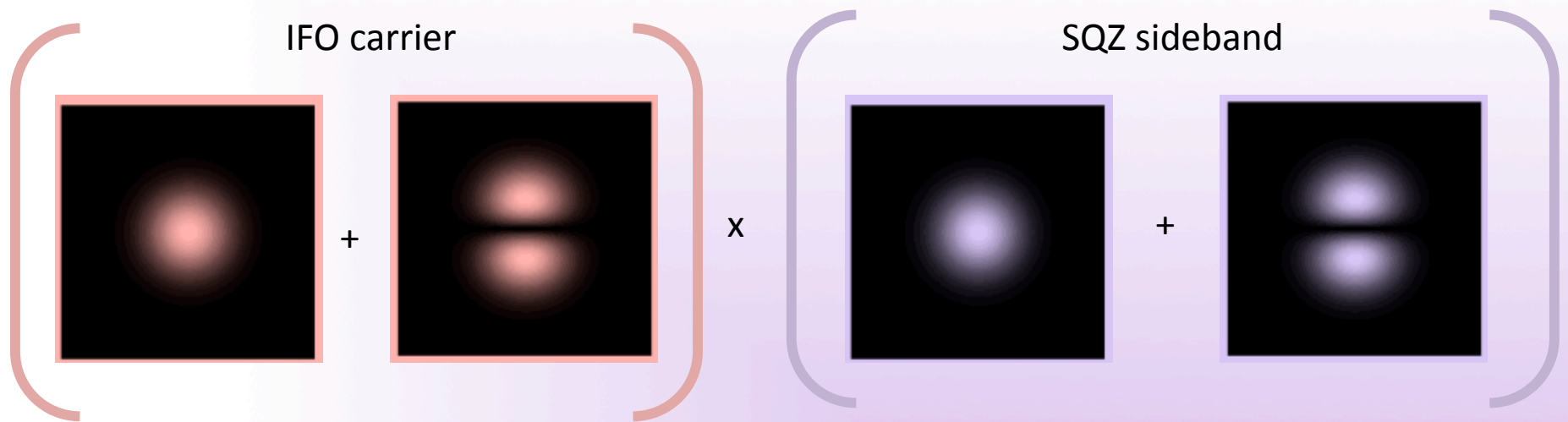


# Phase noise control

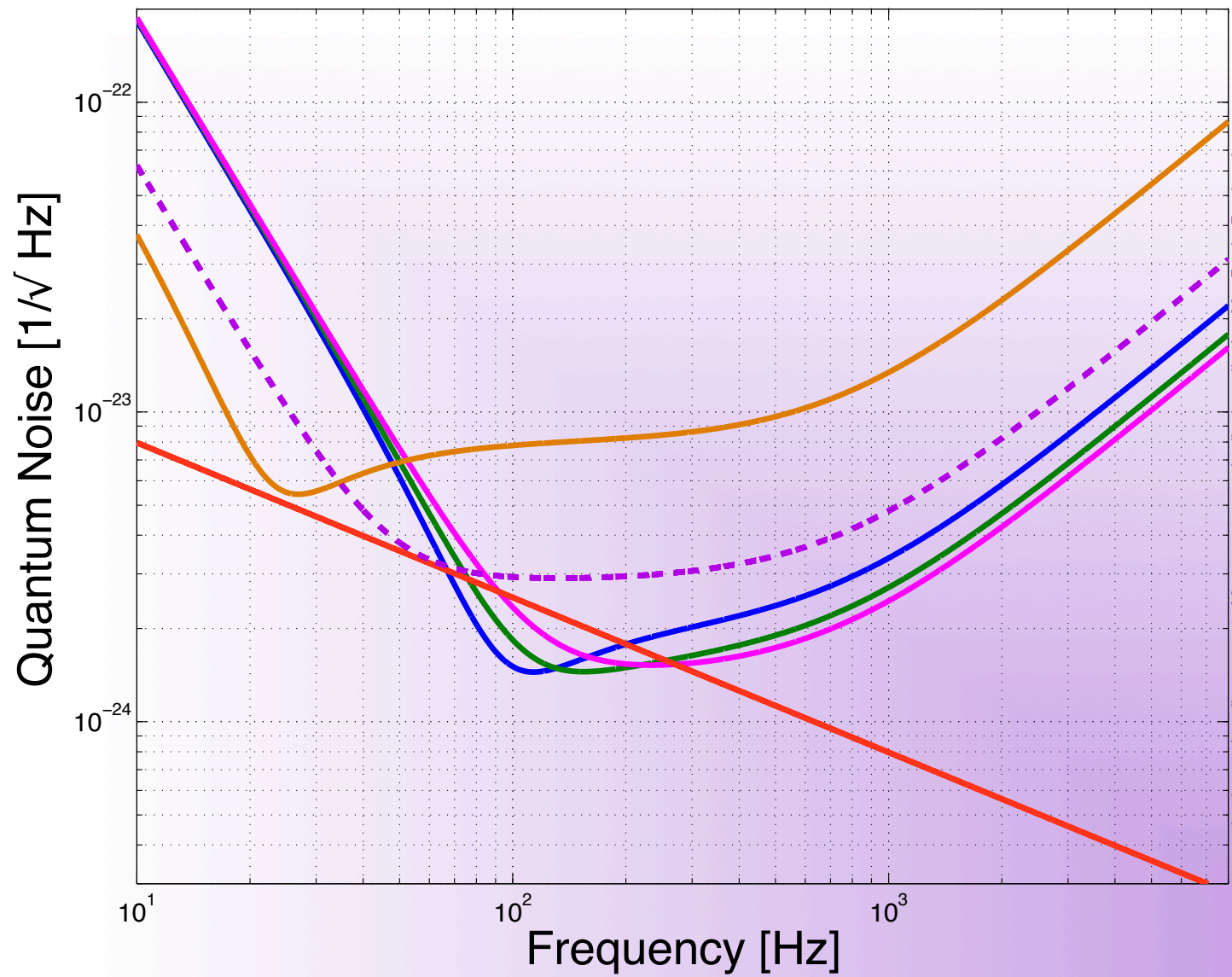
- Bandwidth is limited to 10kHz by arm cavities
- Need to mitigate phase noise at the source
- Changes to control scheme and in vacuum OPO may be necessary for 10-15 dB of squeezing



# Squeezing angle error signal



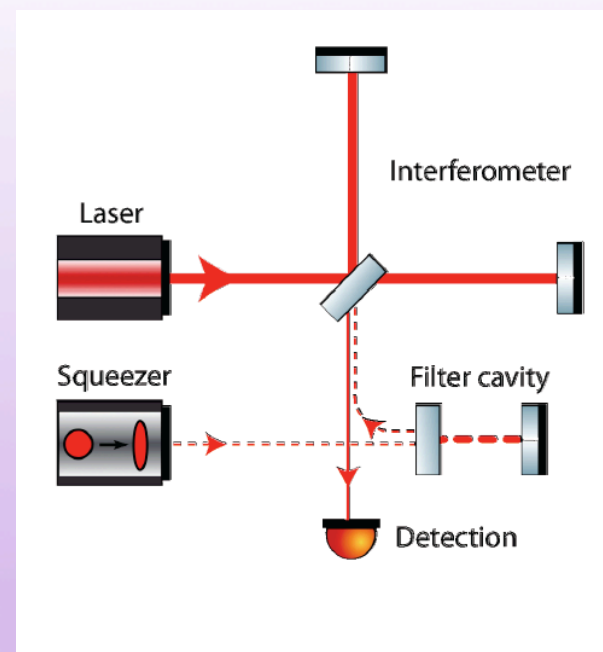
# Quantum noise shaped by squeezed angle



created using gwinc\_fig.m on 16-May-2012 by lisab on lisabs-MacBook-Pro.local

# Frequency Dependent Squeezing

- ✧ High finesse detuned cavity which does the rotation for you
- ✧ Broadband improvement of the quantum noise
- ✧ Theoretically well understood, experimentally challenging
- ✧ Low loss needed:  $F \sim 50,000$  for 100m scale cavities
- ✧ R&D in progress – MIT (P. Kwee and others)  
Caltech (J. Harms and others)

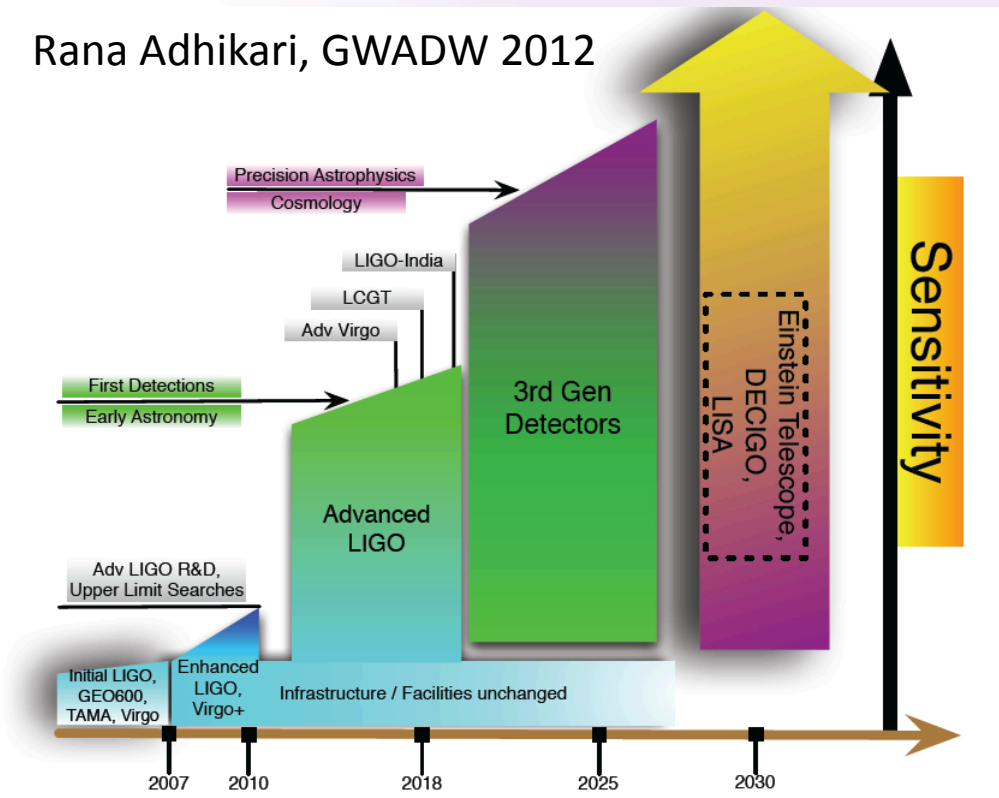


H. J. Kimble, Y. Levin, A. B. Matsko, K. S. Thorne and S. P. Vyatchanin,  
Conversion of conventional gravitational-wave interferometers into quantum  
nondemolition interferometers by modifying their input and/or  
output optics. Phys. Rev. D 65, 022002 (2001).

# Beyond aLIGO: 3<sup>rd</sup> generation

## Can we take another factor of 10 step?

Rana Adhikari, GWADW 2012



✧ Basic idea is to use the same LIGO vacuum envelope

✧ Design study happening now

✧ Still work in progress, one thing already clear:

➔ 10 dB of frequency dependent squeezing needed!

# The Message

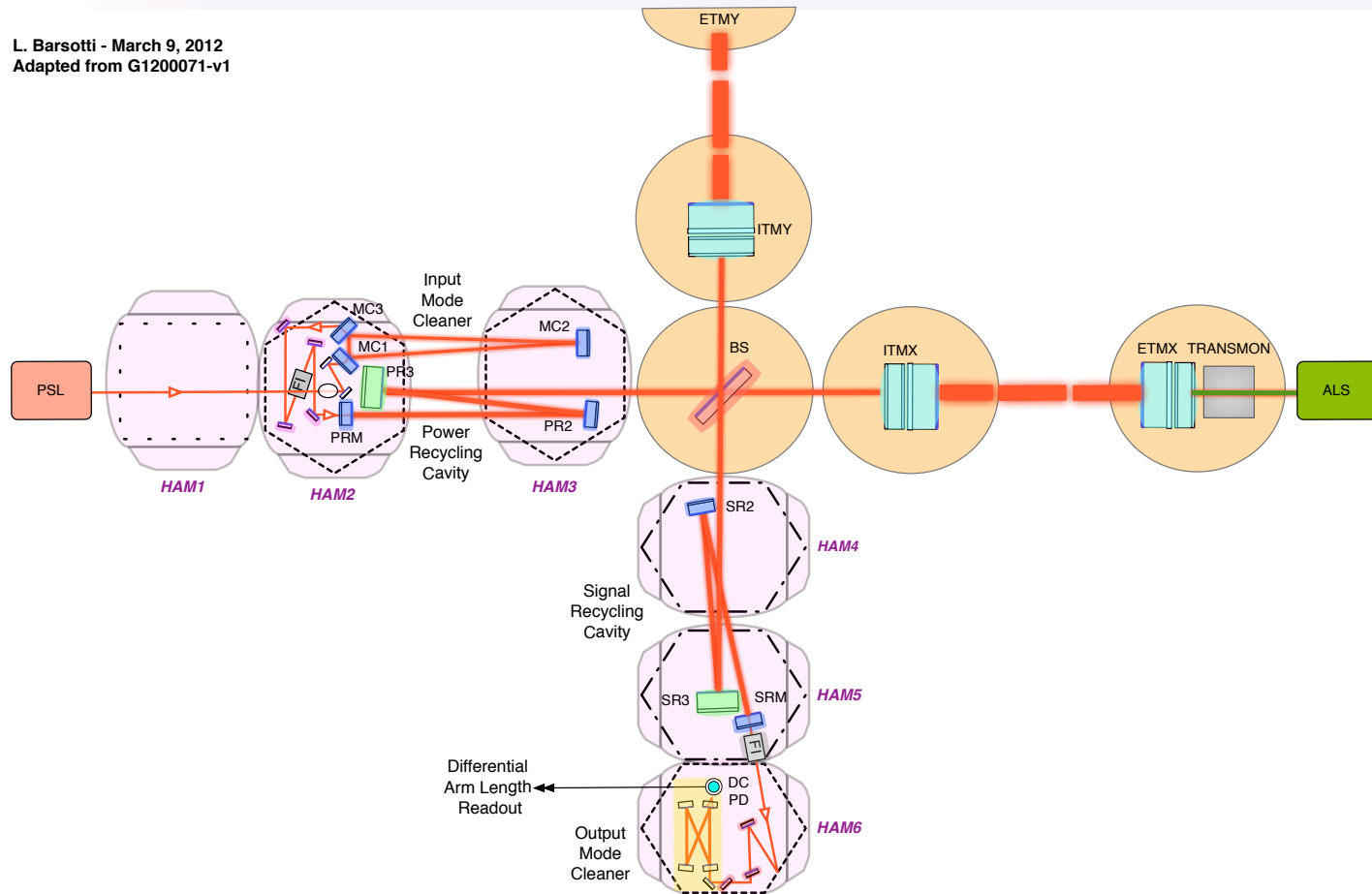
- ✧ Squeezing can reduce quantum noise, and improve the sensitivity of GW detectors
- ✧ Large scale interferometers with squeezing: **DONE!**
  - ✧ Work needed to achieve 24/7 long term stability at maximum squeezing and reduce optical losses
  - ✧ H1 squeezing experiment completed, GEO600 operating with squeezing right now
- ✧ In a good position to make squeezing available for Advanced detectors and beyond





# Advanced LIGO configuration

L. Barsotti - March 9, 2012  
Adapted from G1200071-v1



- ✧ Arm cavities, power and signal recycling cavity
- ✧ Up to ~800 kW of light stored in the arms

# Quantum States

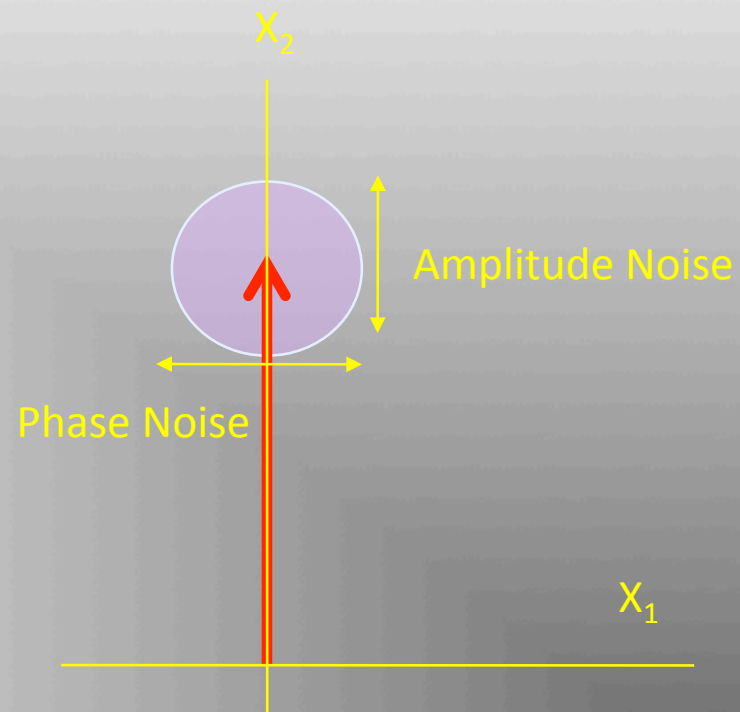
## ✧ Quantization of the electro-magnetic field

Quadrature Field Amplitudes

$$\hat{E} = \hat{X}_1 \cos \omega t + i\hat{X}_2 \sin \omega t$$

Heisenberg uncertainty principle:

$$\Delta X_1 \Delta X_2 \geq 1$$



# Vacuum Fluctuations

- ✧ When average amplitude is zero, the variance remains
- ✧ Vacuum fluctuations are everywhere that classically there is no field....
- ✧ ...like at the output port of your interferometer!

